

Lattice QCD computations of proton decay matrix elements

Yasumichi Aoki

RIKEN Center for Computational Science

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@INT

Proton Decay Matrix Elements project in PACS collaboration

Members:


- **Eigo Shintani** (Tsukuba)
- **Ryutaro Tsuji** (KEK)
- Yoshinobu Kuramashi (Tsukuba)
- YA

using PACS Wilson configurations on physical-mass ud , s quarks

$V=64^4$ $L=5.5$ fm

$1/a=2.3$ GeV

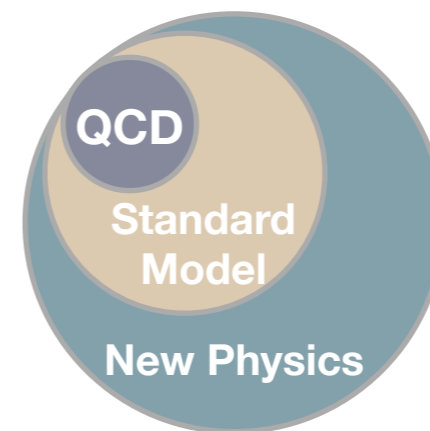
$a = 0.08$ fm



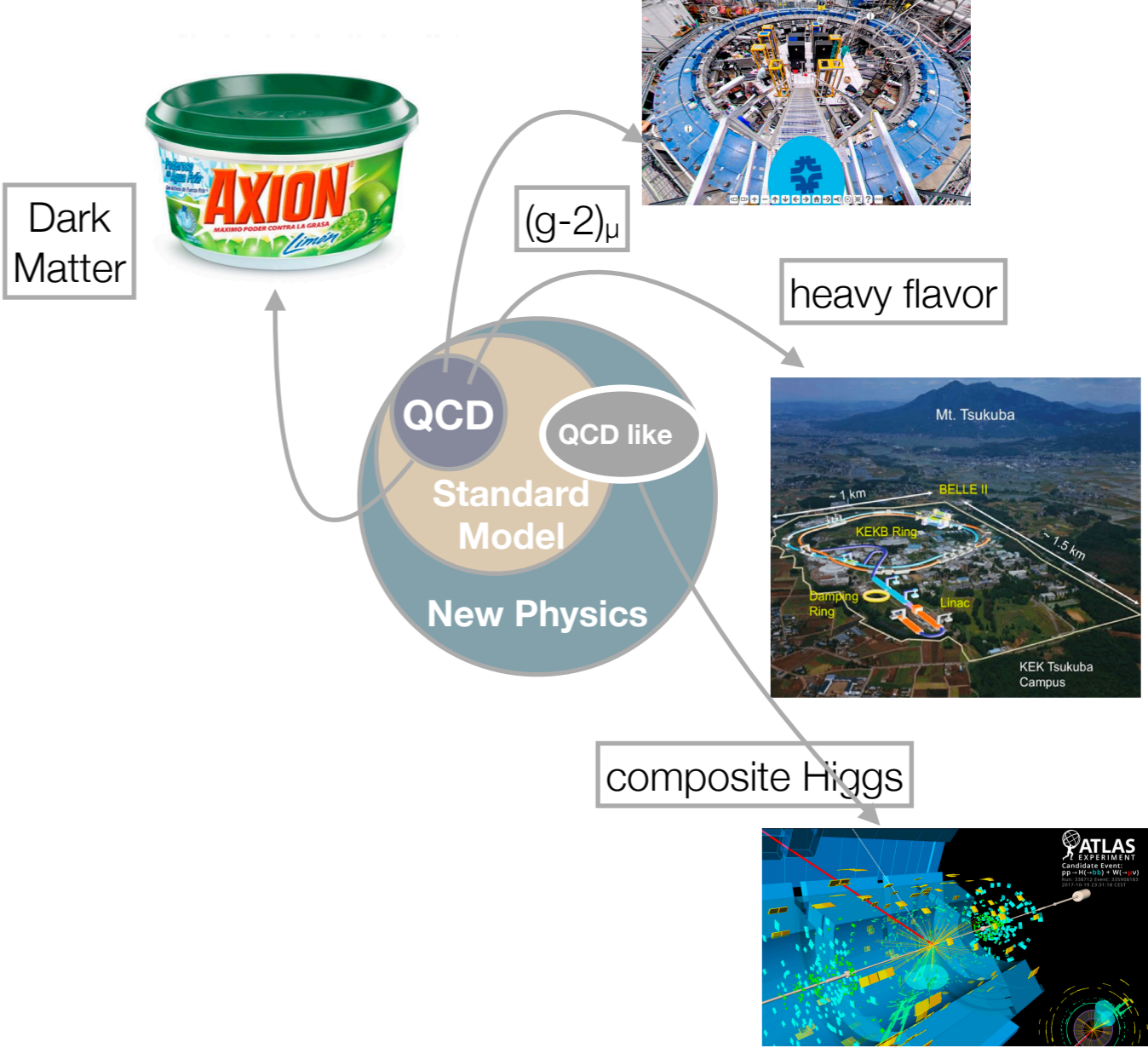
Program for Promoting Researches
on the Supercomputer Fugaku
Large-scale lattice QCD simulation
and development of AI technology

new physics through QCD

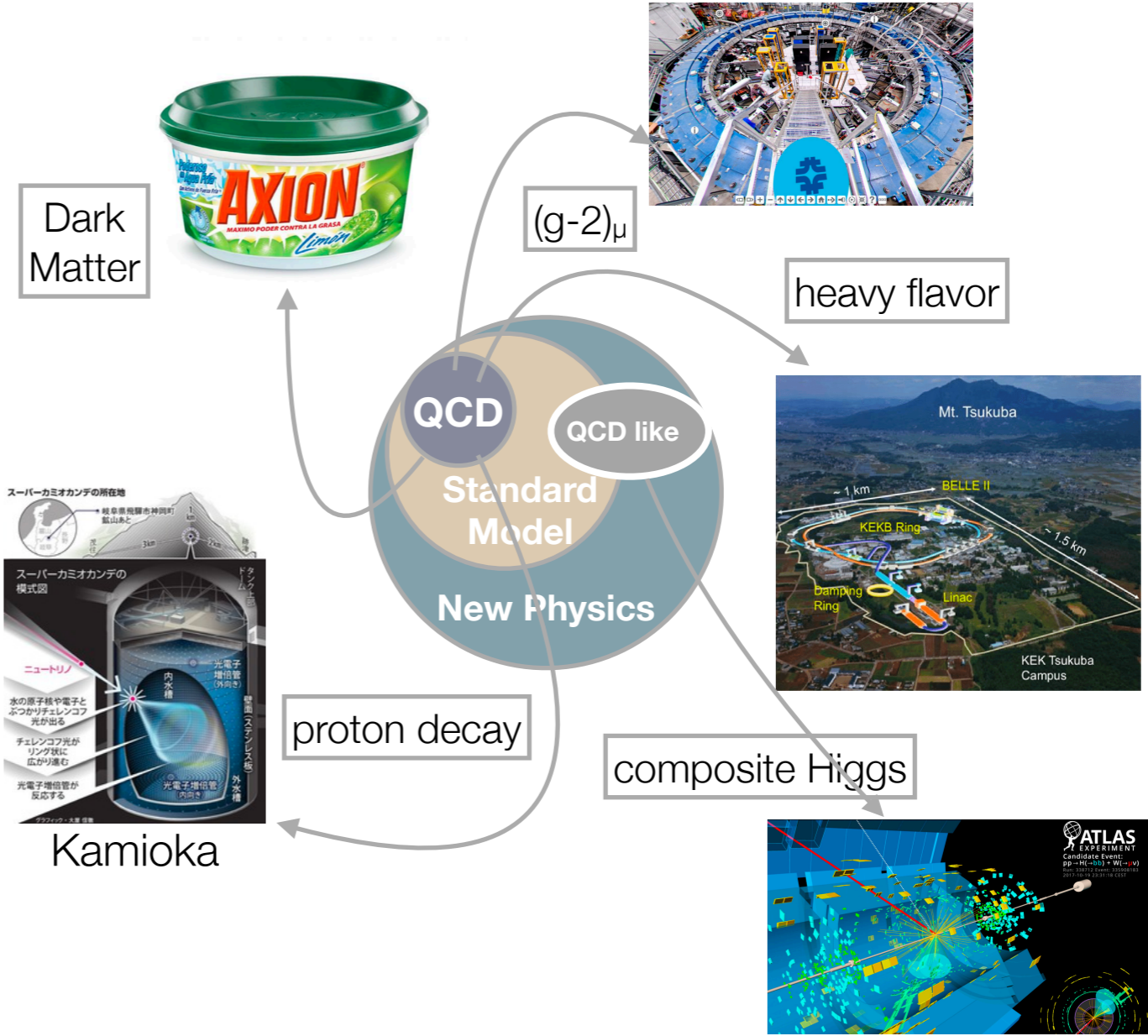
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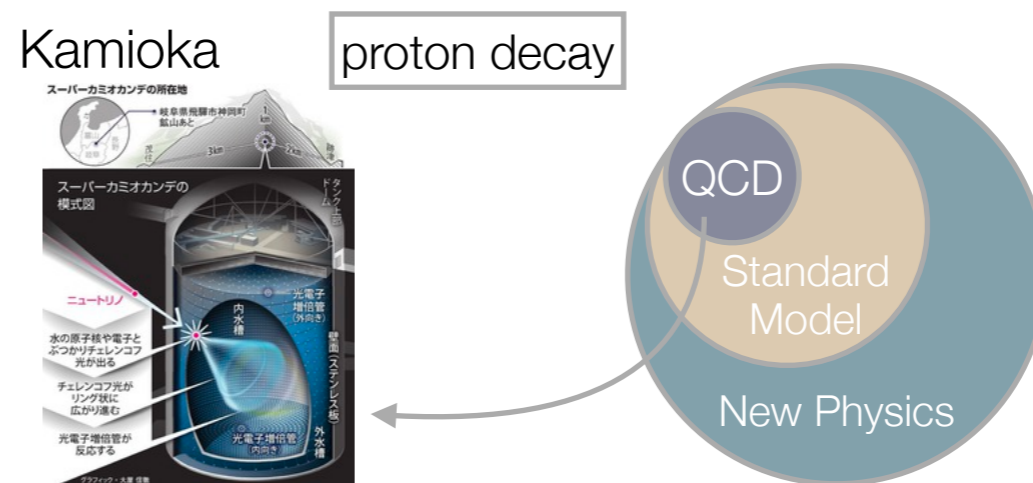
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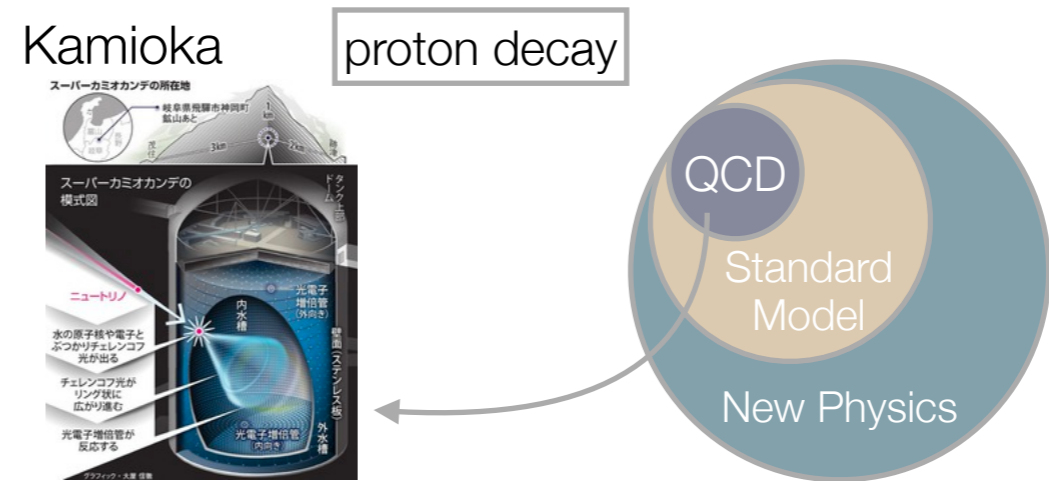
陽子崩壊 - 新実験計画



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- Smoking Gun of New Physics
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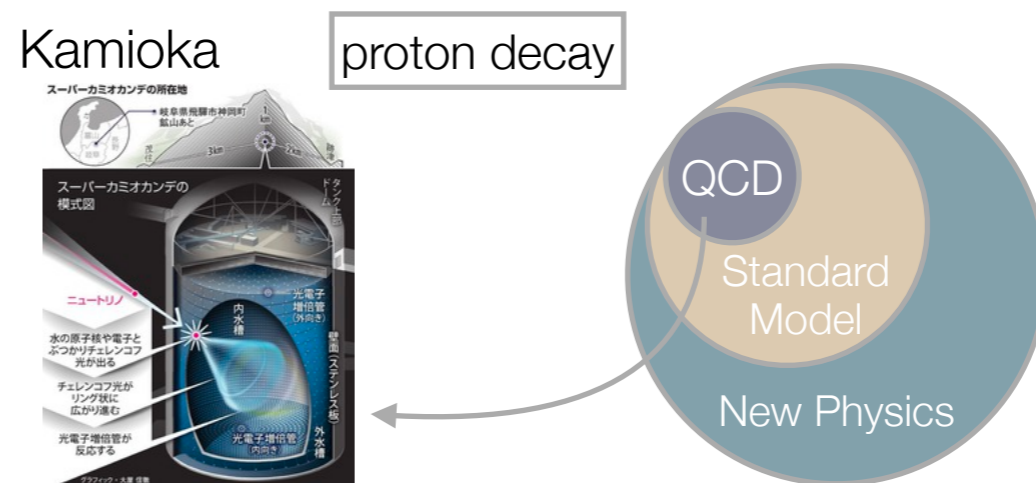
$$1/\tau_p \propto [\text{QCD param.}] * [\text{NewPhys. param}]$$



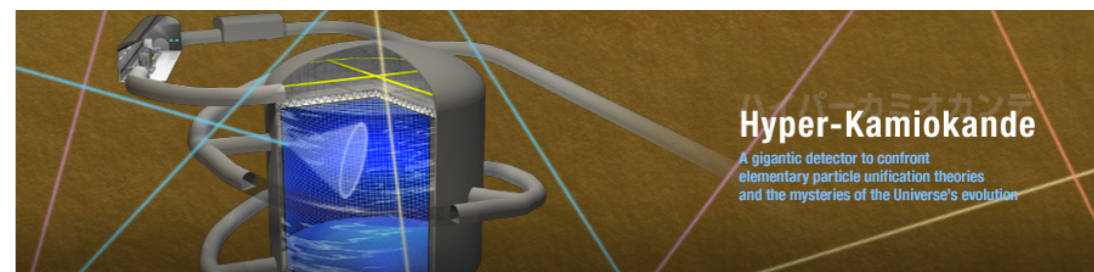
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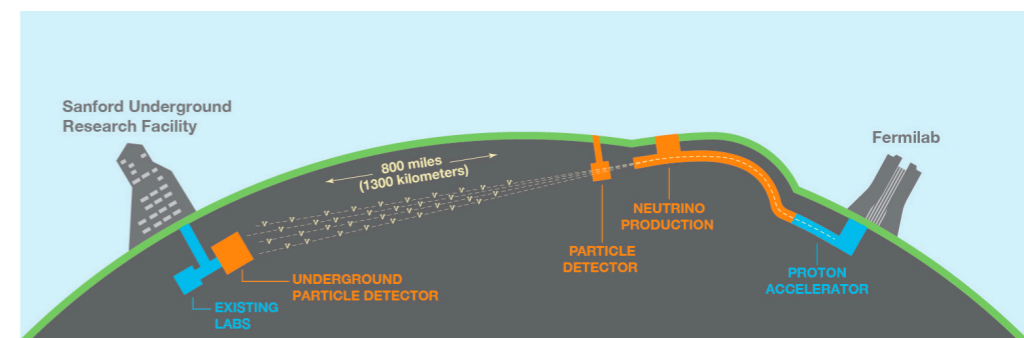
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- New experiments are under preparation
 - HyperKamiokande in Japan
 - ~ x5 V of SuperKamiokande



- DUNE (Deep Underground Neutrino Experiment) in USA
 - Liquid Argon
 - sensitivity to Kaon



relevant form factors

Relevant form factors

- $$\mathcal{L}^{GUT} \rightarrow \mathcal{L}^{SM} + \sum_i C^i(\mu) \cdot O^i(\mu) + \dots$$

$$O^i(\mu) = (qq)_\Gamma (ql)_{\Gamma'}$$

- yields a decay: baryon \rightarrow meson + anti-lepton

$$\langle \pi^0, e^+ | p \rangle_{GUT} = \sum_i C^i(\mu) \cdot \langle \pi^0, e^+ | O^i(\mu) | p \rangle_{SM}$$

$$\langle \pi^0, e^+ | (ud)(eu) | p \rangle^i = \bar{v}_e^c \cdot \langle \pi^0 | (ud)u | p \rangle$$

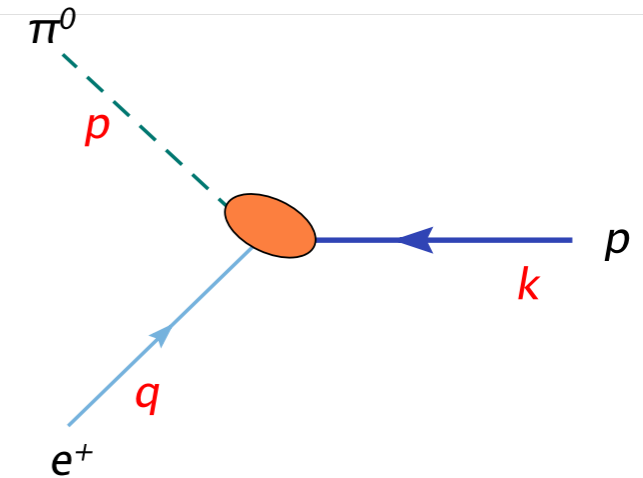
- a convenient parametrization using 2 form factors [JLQCD 2000]

$$\langle \pi^0 | (ud)_\Gamma u_L | p \rangle = P_L [W_0 - \frac{i \not{q}}{m_p} W_1] u_p$$

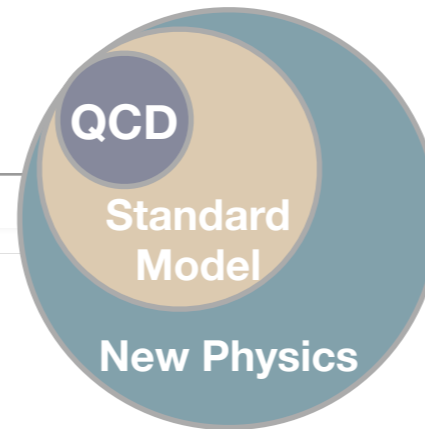
$$\langle \pi^0, e^+ | (ud)_\Gamma (eu)_L | p \rangle = \underbrace{W_0}_{\text{relevant}} \cdot (v_e, u_p)_L + \frac{m_e}{m_p} \underbrace{W_1}_{\text{irrelevant}} \cdot (v_e, u_p)_R$$

- partial width

$$\Gamma(p \rightarrow \pi^0 + e^+) = \frac{m_p}{32\pi} \left[1 - \left(\frac{m_\pi}{m_p} \right)^2 \right]^2 \left| \sum_i C^i W_0^i(p \rightarrow \pi^0) \right|^2$$



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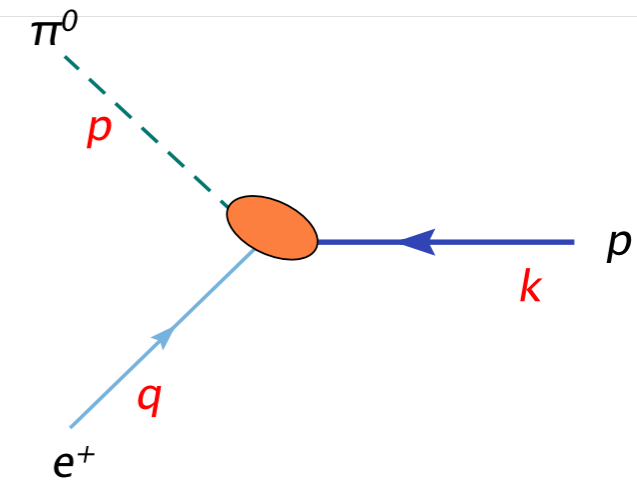
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QCD matrix element of nucleon decay

* $\langle \pi^0 | (ud)_{\Gamma} u_L | p \rangle = P_L [W_0 - \frac{i \not{q}}{m_p} W_1] u_p \quad \Gamma = R, L$

* indirect method: LO **approximation** of W_0 in ChPT: Claudson, Wise, Hall, 1982

$$W_0[\langle \pi^0 | (ud)_{R} u_L | p \rangle] \simeq \frac{\alpha}{\sqrt{2}f} (1 + D + F)$$

$$W_0[\langle \pi^0 | (ud)_{L} u_L | p \rangle] \simeq \frac{\beta}{\sqrt{2}f} (1 + D + F)$$



f : pion decay constant

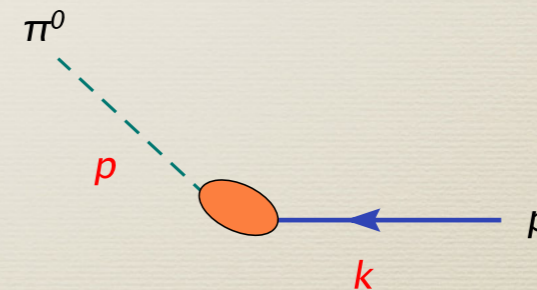
$D + F = g_A$: nucleon axial charge

 $\langle 0 | (ud)_{R} u_L | p \rangle = \alpha P_L u_p$

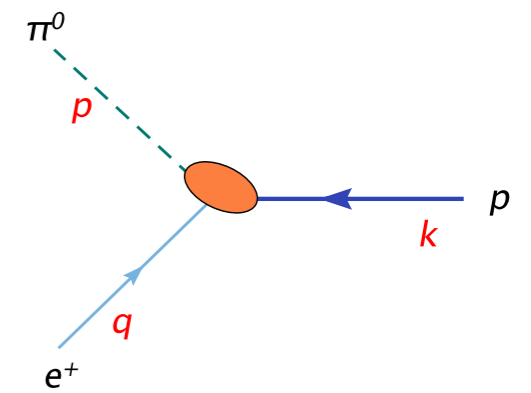
$\langle 0 | (ud)_{L} u_L | p \rangle = \beta P_L u_p$

* **direct method**: calculates the form factor W_0 of $p \rightarrow PS$ matrix elements directly

* comparison given later...



constraining GUT



• partial width

$$1/\tau(p \rightarrow \pi^0 + e^+) = \frac{m_p}{32\pi} \left[1 - \left(\frac{m_\pi}{m_p} \right)^2 \right]^2 \left| \sum_i C^i W_0^i(p \rightarrow \pi^0) \right|^2$$

• given GUT and $W_0^i(\mu)$ from lattice, $C^i(\mu)$ is constrained using experimental lower bound of proton lifetime

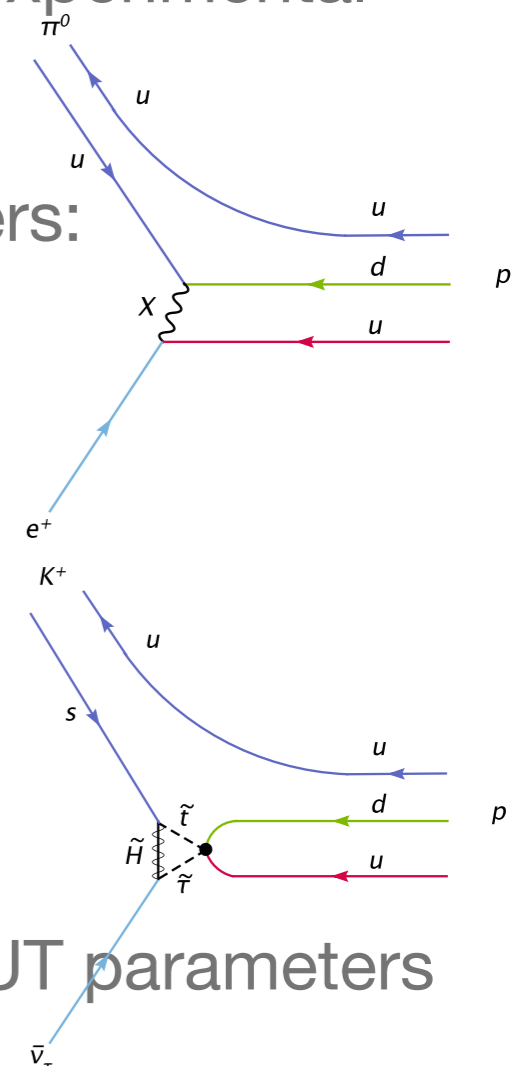
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• m_X, \dots for heavy boson mediated dim 6 nucleon decay

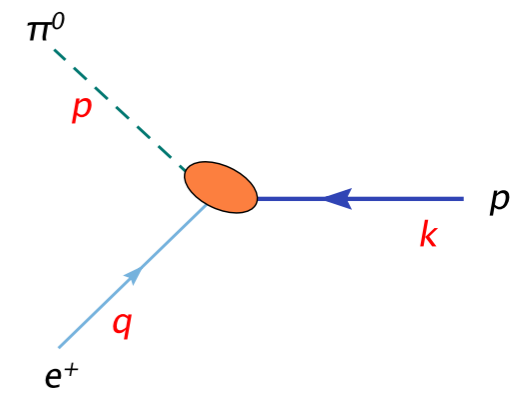
• m_C and spectrum of sparticles for colored higgs mediated dim 5 nucleon decay

➡ complement to LHC

• constraints on $C^i(\mu)$ may be transcribed into constraints of GUT parameters

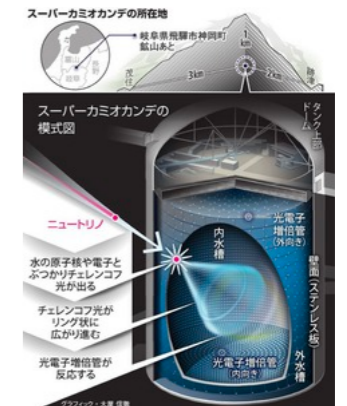


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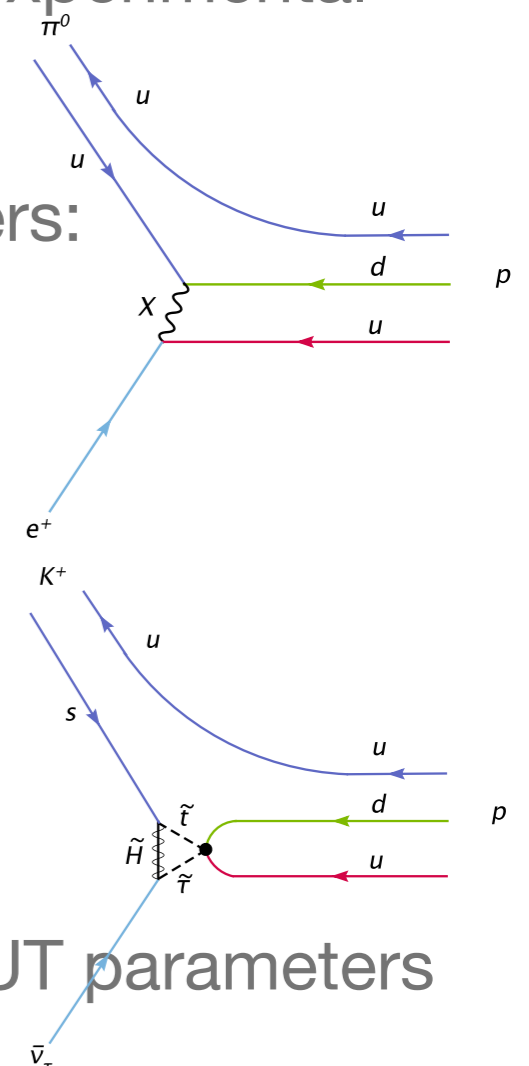
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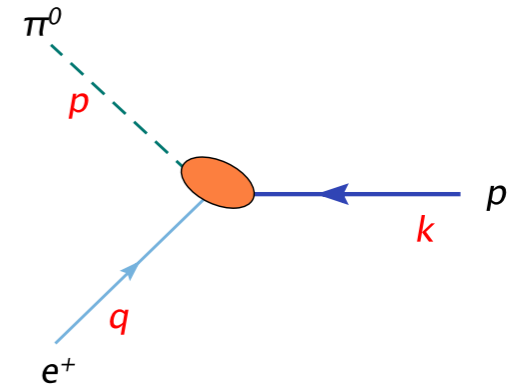
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DWF calculation of proton decay matrix elements after JLQCD 2000

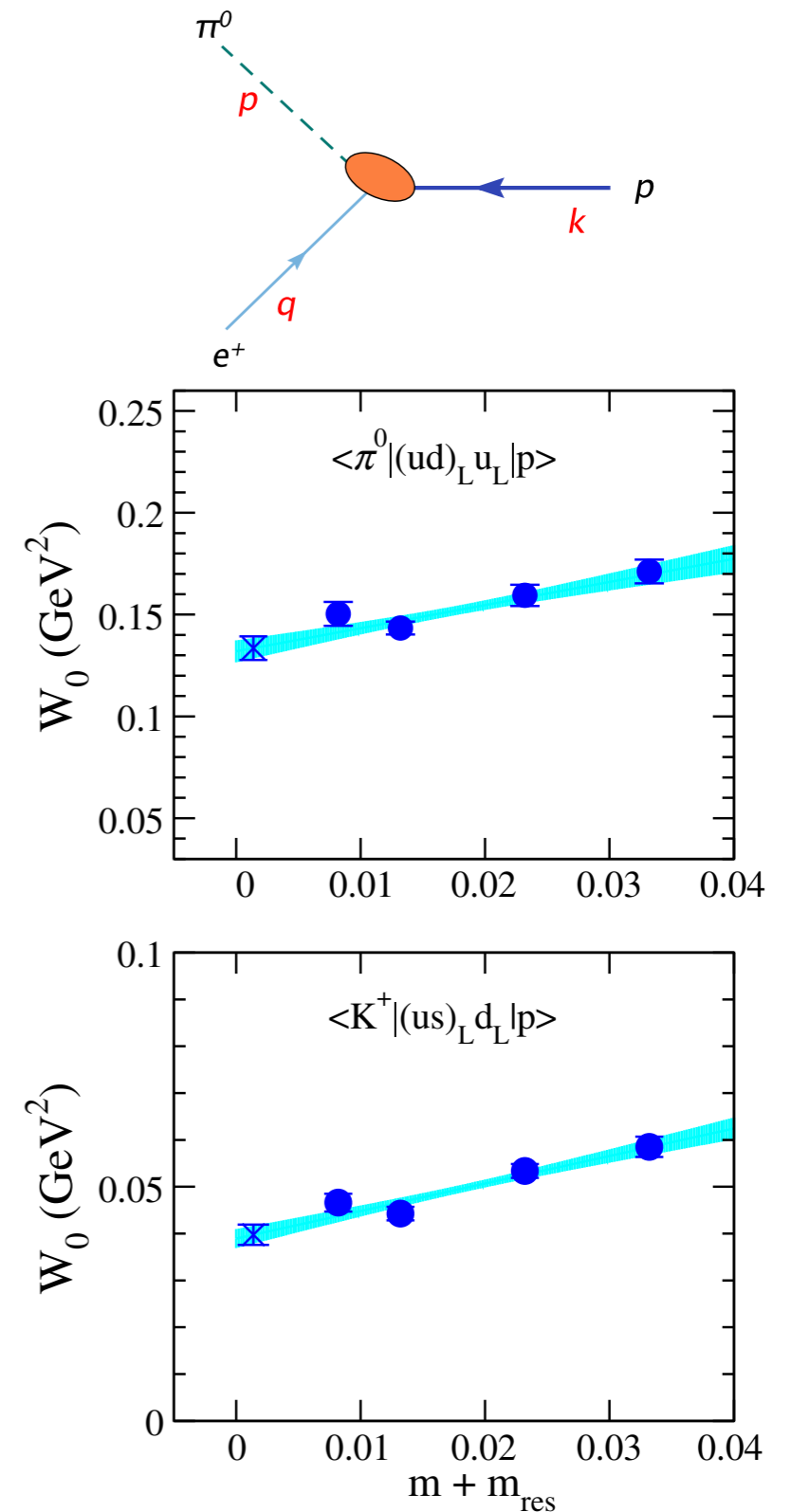
- quench direct & $N_f=2$ LEC (2006)
 - NPR scheme constructed
 - YA, C. Dawson, J. Noaki, A. Soni
- $N_f=2+1$ LEC (2008)
 - YA, P. Boyle, P. Cooney, L. Del Debbio, R. Kenway, C. Maynard, A. Soni, R. Tweedie
- $N_f=2+1$ direct (2014)
 - YA, T. Izubuchi, E. Shintani, A. Soni

- RBC & RBC/UKQCD collaborations
 - $N_f=2+1$ direct with AMA (2017)
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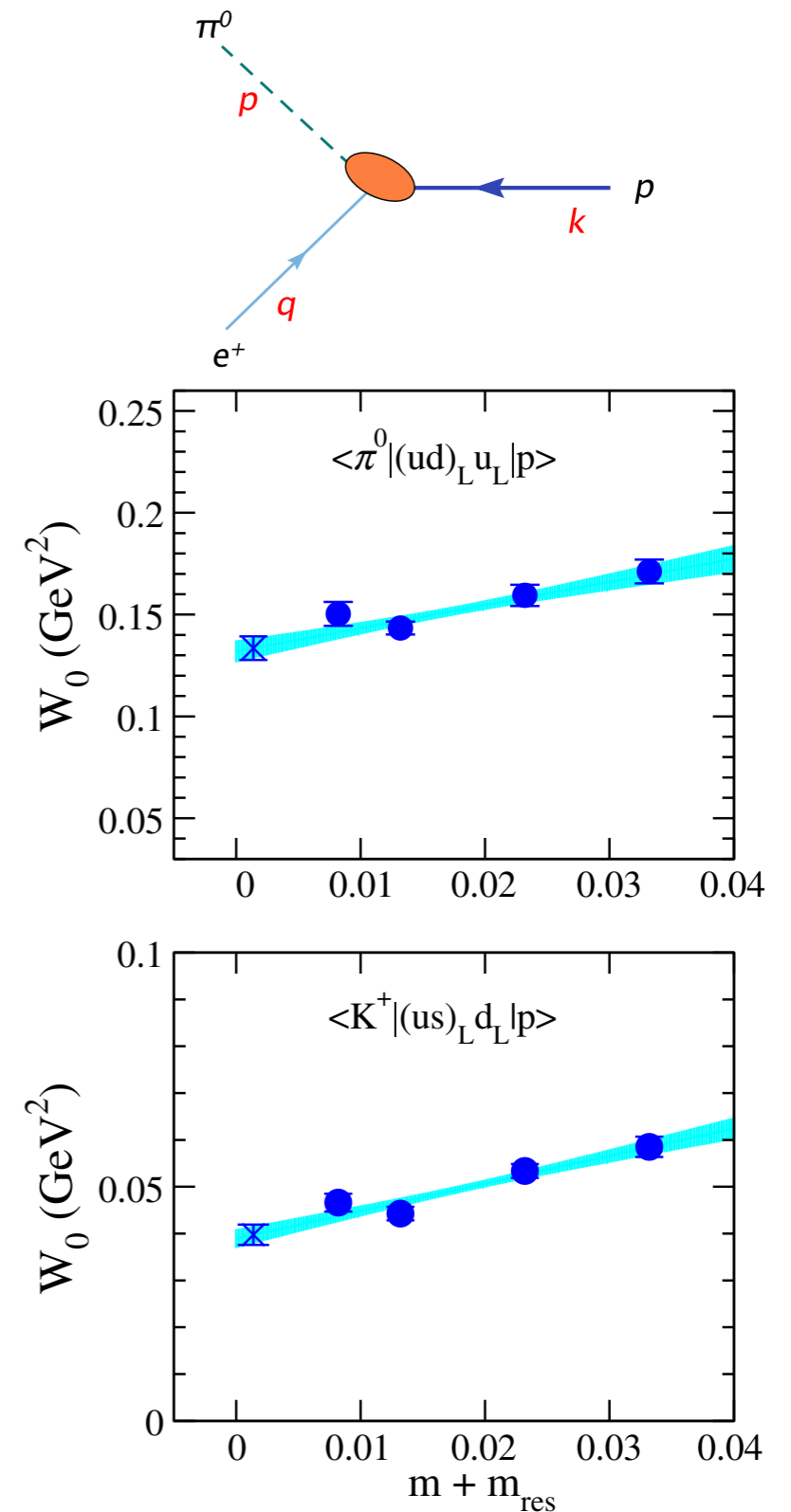
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- pros

- DWF: renormalization is simple
- AMA: statistically improved a lot

- cons

- lightest pion ~ 330 MeV
 - linear extrapolation may eventually fail



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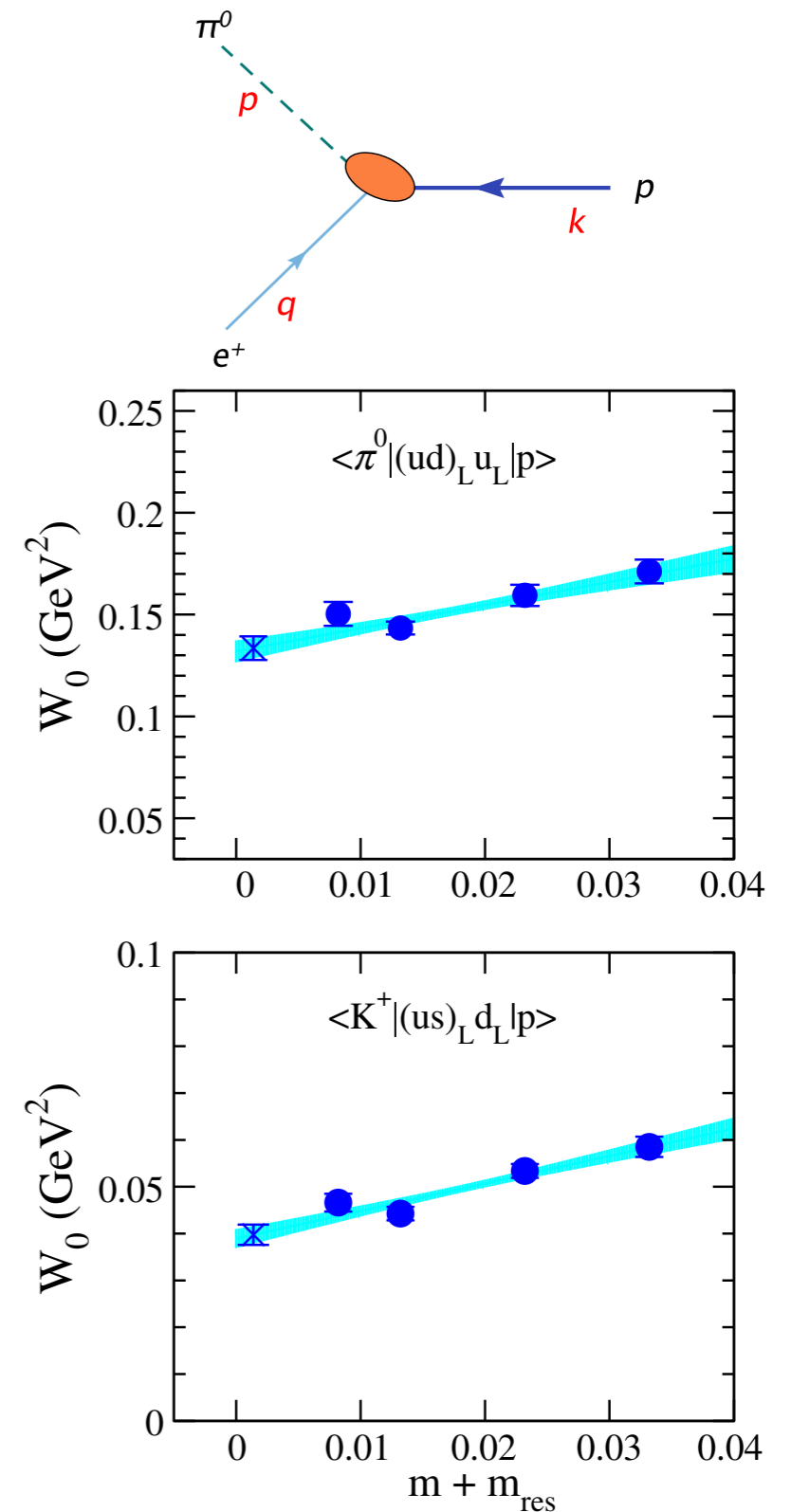
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Skirm chiral bag model (Martin-Stavenga 2016)

- drastic decrease towards chiral limit due to topological stabilization of proton

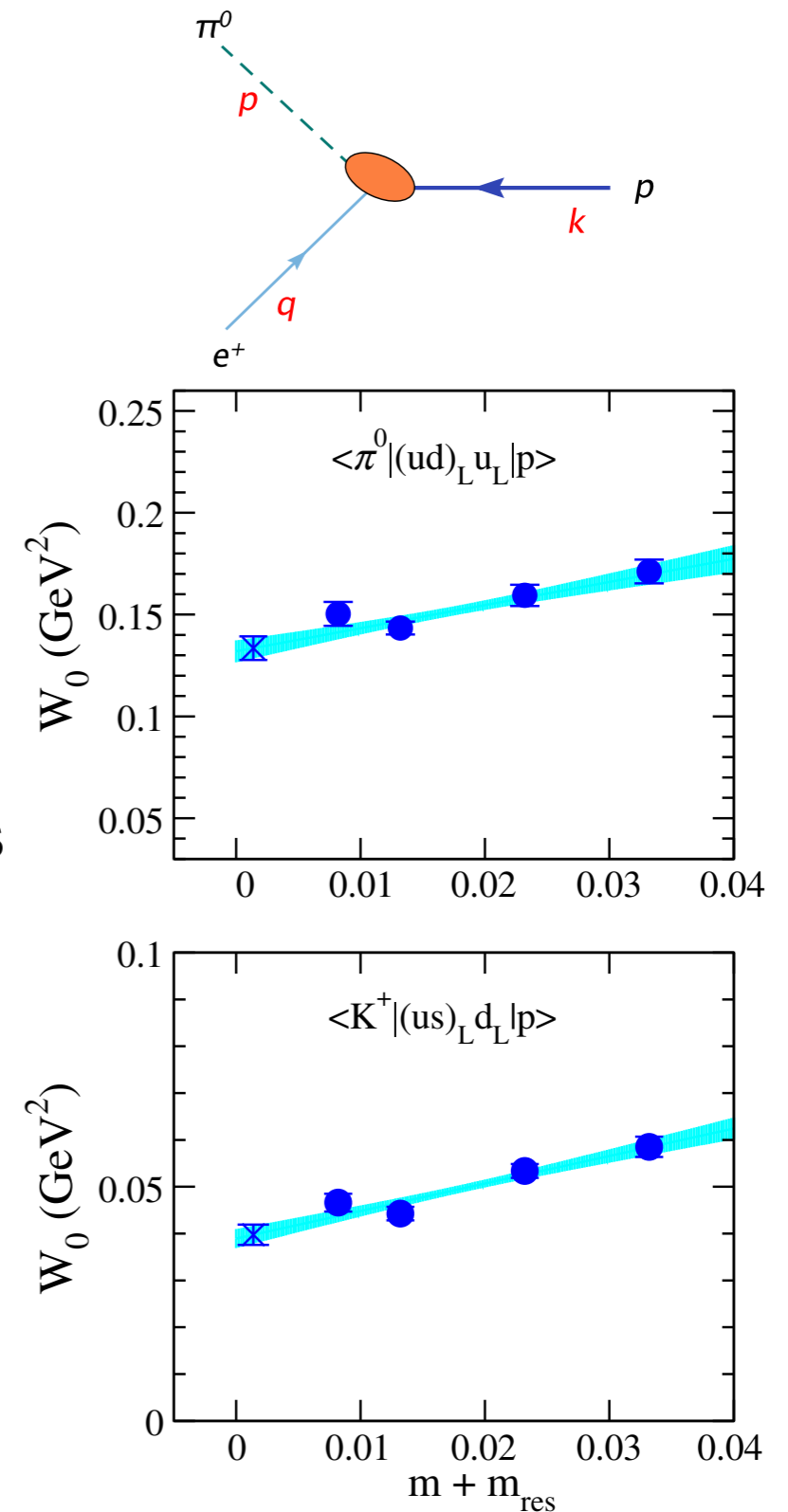


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- compare w/ polynomial w/ higher order
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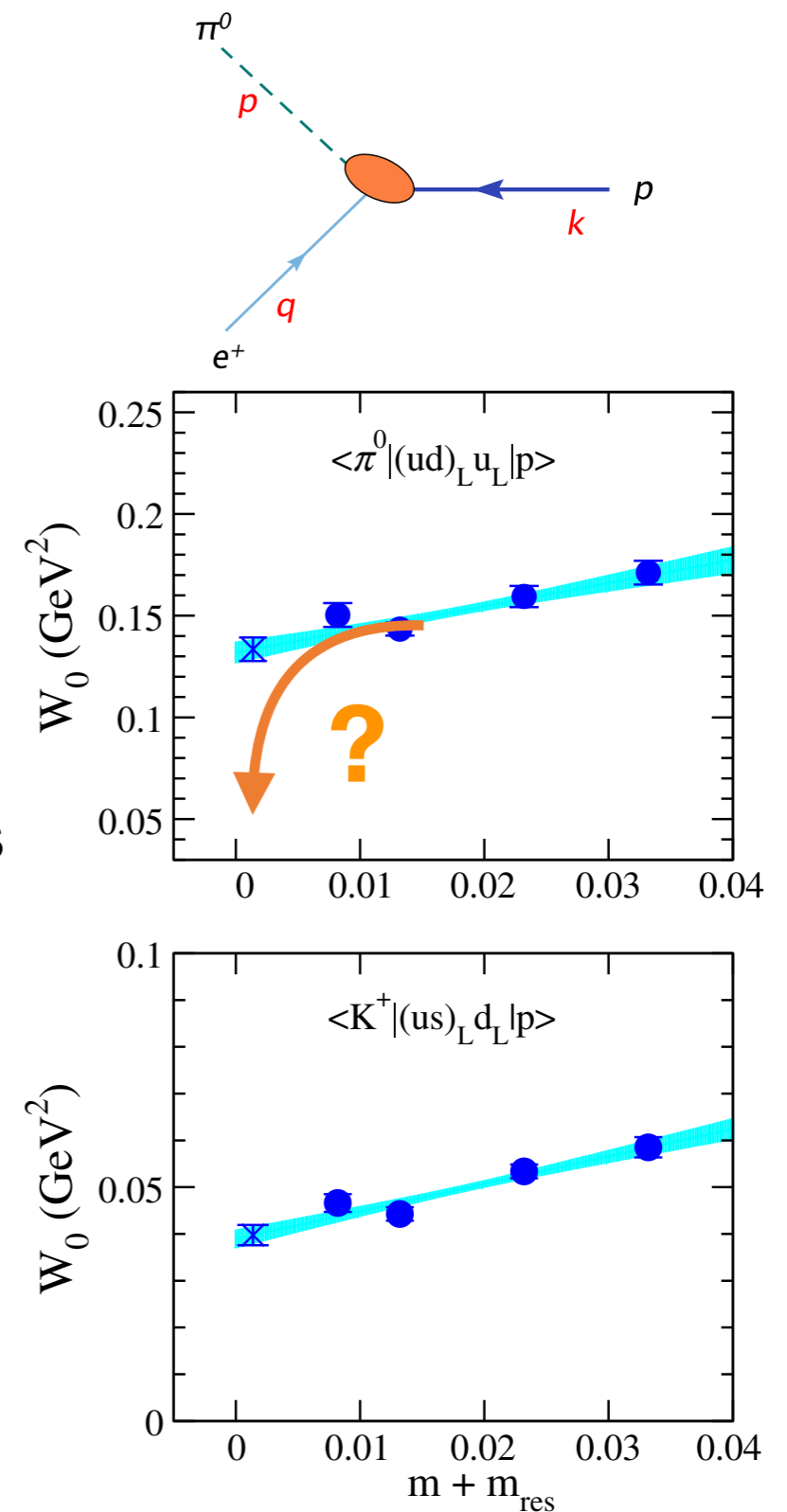


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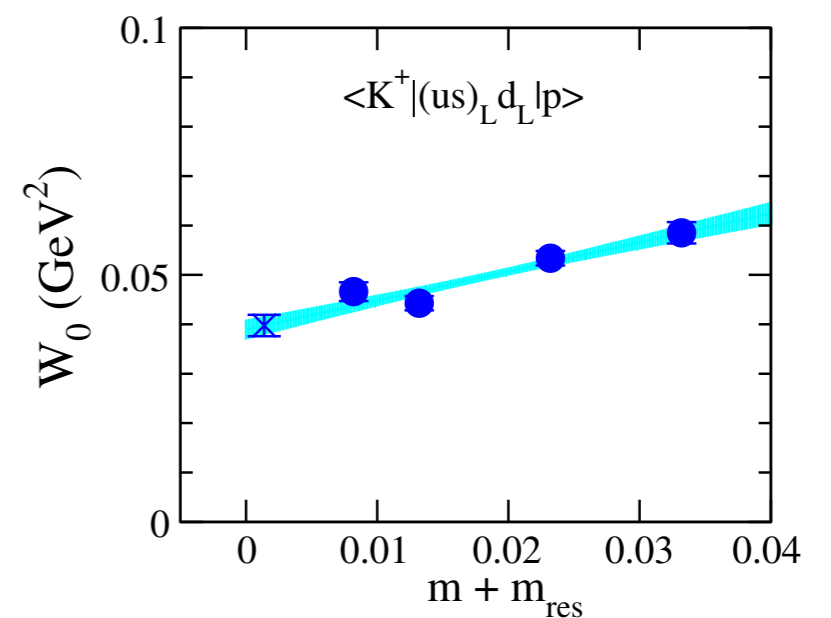
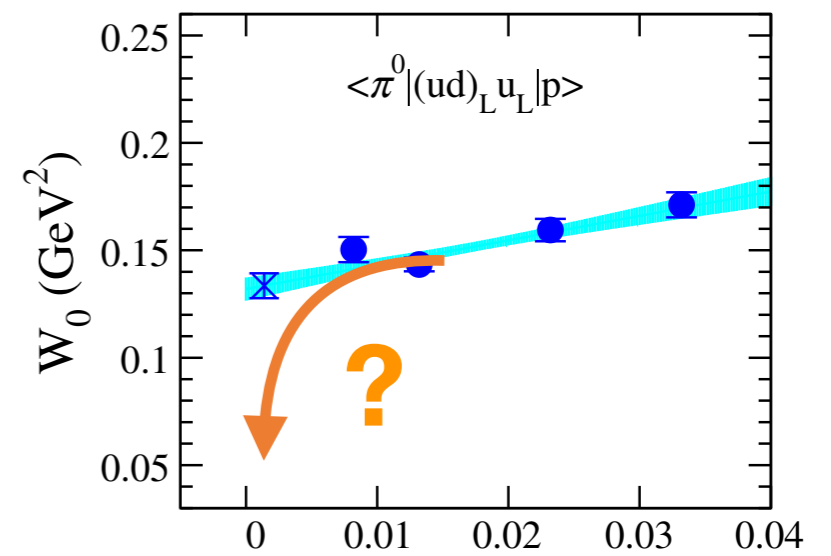
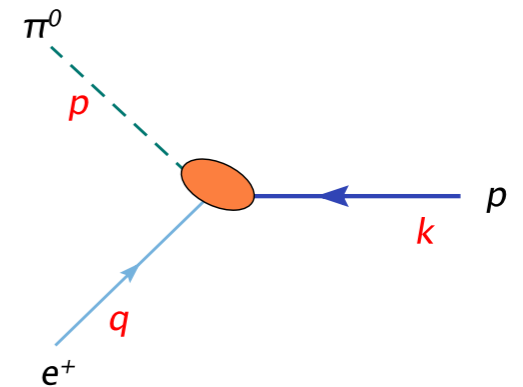
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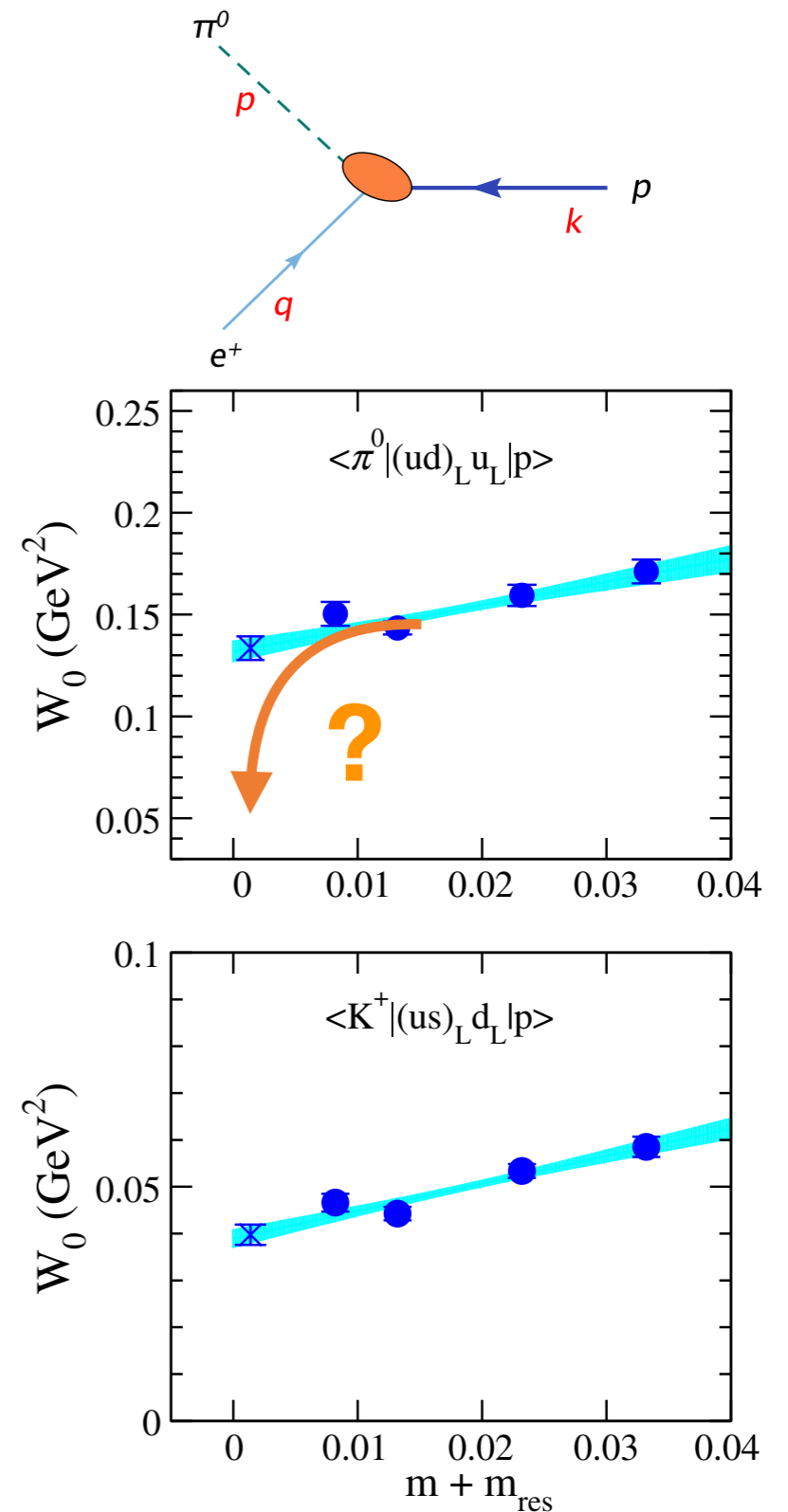
systematic error from the chiral extrapolation:

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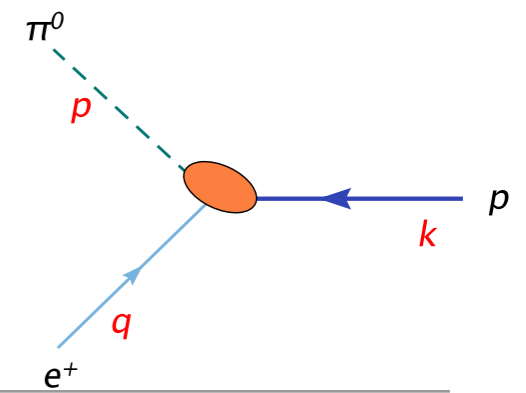
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No extrapolation needed if one use physical point ensembles

Let's use them



Error budget of RBC/UKQCD 2017



IMPROVED LATTICE COMPUTATION OF PROTON DECAY ...

PHYSICAL REVIEW D **96**, 014506 (2017)

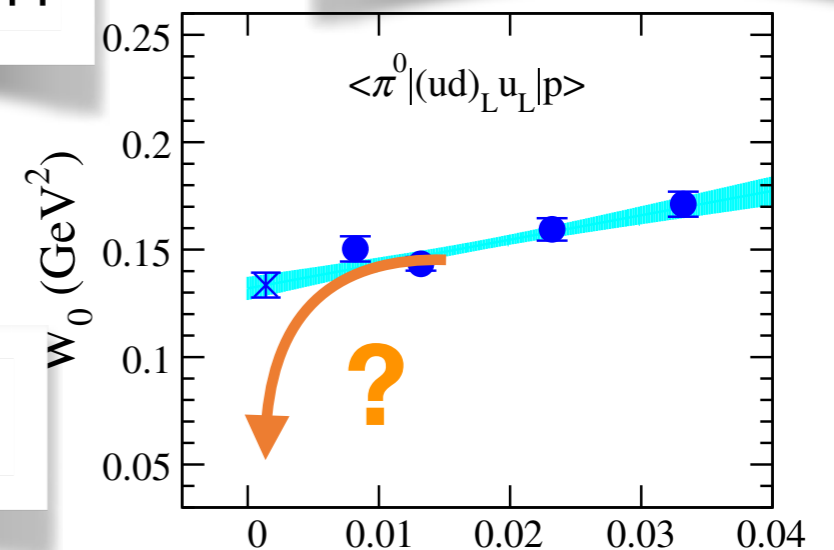
TABLE IV. Table of the renormalized W_0 in the physical kinematics at 2 GeV in the $\bar{M}\text{S}$ NDR scheme. The fourth column contains the relative error of the systematic uncertainties. “ χ ” comes from the chiral extrapolation given from three different fitting ranges as explained in the text. The “ q^4 ” and “ a^2 ” columns are the uncertainties of the higher-order correction than $\mathcal{O}(q^2)$ and the lattice artifact at $\mathcal{O}(a^2)$, respectively. The “ m_s ” column is the uncertainty coming from using the unphysical strange quark mass. Δ_Z and Δ_a are the errors of the renormalization factor and lattice scale estimate, respectively.

Matrix element	W_0 GeV ²	Statistical[%]	Total	Systematic error [%]						
				χ	q^4	mq^2	a^2	m_s	Δ_a	Δ_Z
$\langle \pi^0 (ud)_R u_L p \rangle$	-0.131(4)(13)	3.0	9.7	1.8	0.7	0.3	5.0	...	0.6	8.1
$\langle \pi^0 (ud)_L u_L p \rangle$	0.134(5)(16)	3.4	11.6	5.7	2.3	2.6				

chiral extrapolation

renormalization

this possibility is not taken into account →



Lattice computation

- * *Lattice gauge theory: gauge theory on discrete Euclidian space-time (lattice spacing a)*
- * *a regularization of gauge theory with manifest gauge invariance*
- * *with finite volume and “ a ”, path integral can be performed using (super) computer*
- * *continuum limit $a \rightarrow 0$ has to be performed, or discretization error must be estimated*
- *
$$\mathcal{L}(a) = \mathcal{L}_{QCD} + a \sum_i c_i^{(5)} O_i^{(5)} + a^2 \sum_j c_j^{(6)} O_j^{(6)} + \dots$$
- * *all $O^{(5)}$ break chiral symmetry*
- * *If the lattice action has chiral symmetry, no $O(a)$ error! \rightarrow more continuum like*
- * *Lattice action with chiral symmetry available*
- * ***Domain wall fermions (DWF)** (Kaplan, Furman-Shamin),
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lattice matrix element: direct \Leftrightarrow indirect

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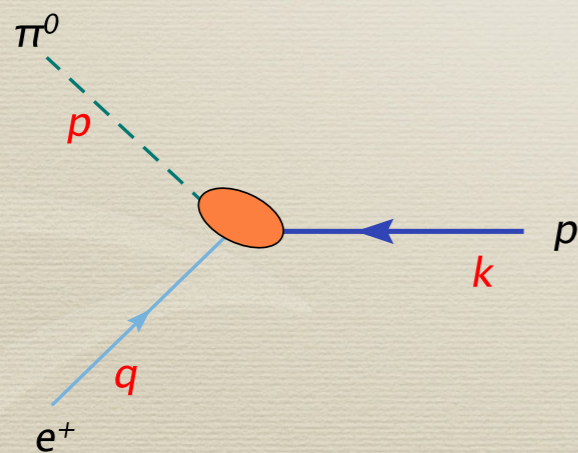
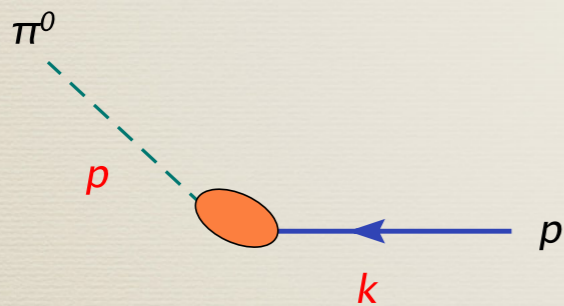
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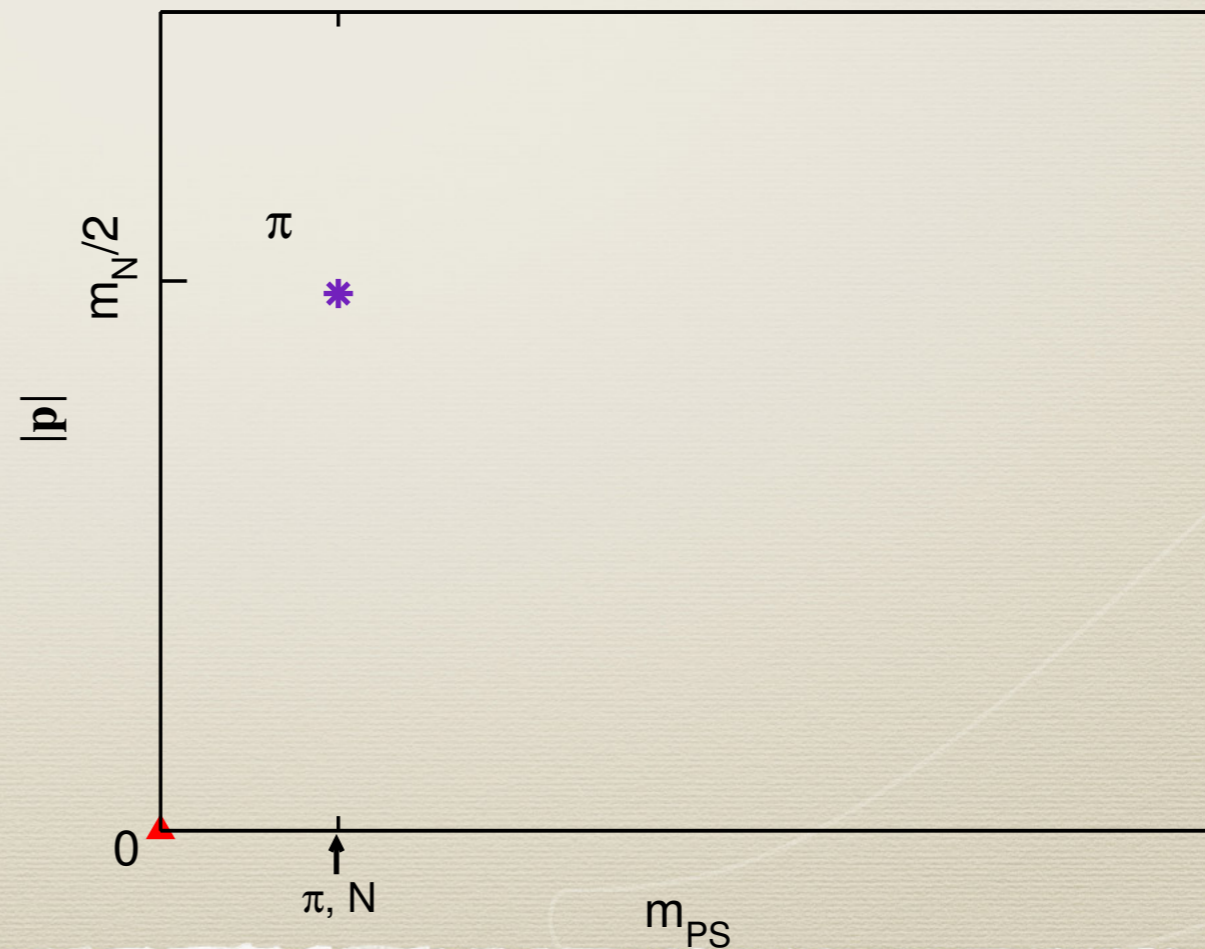
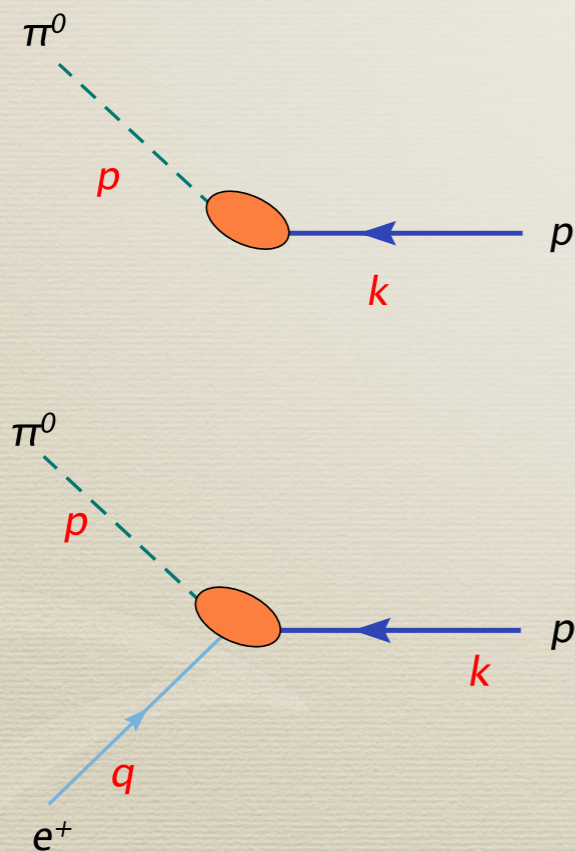
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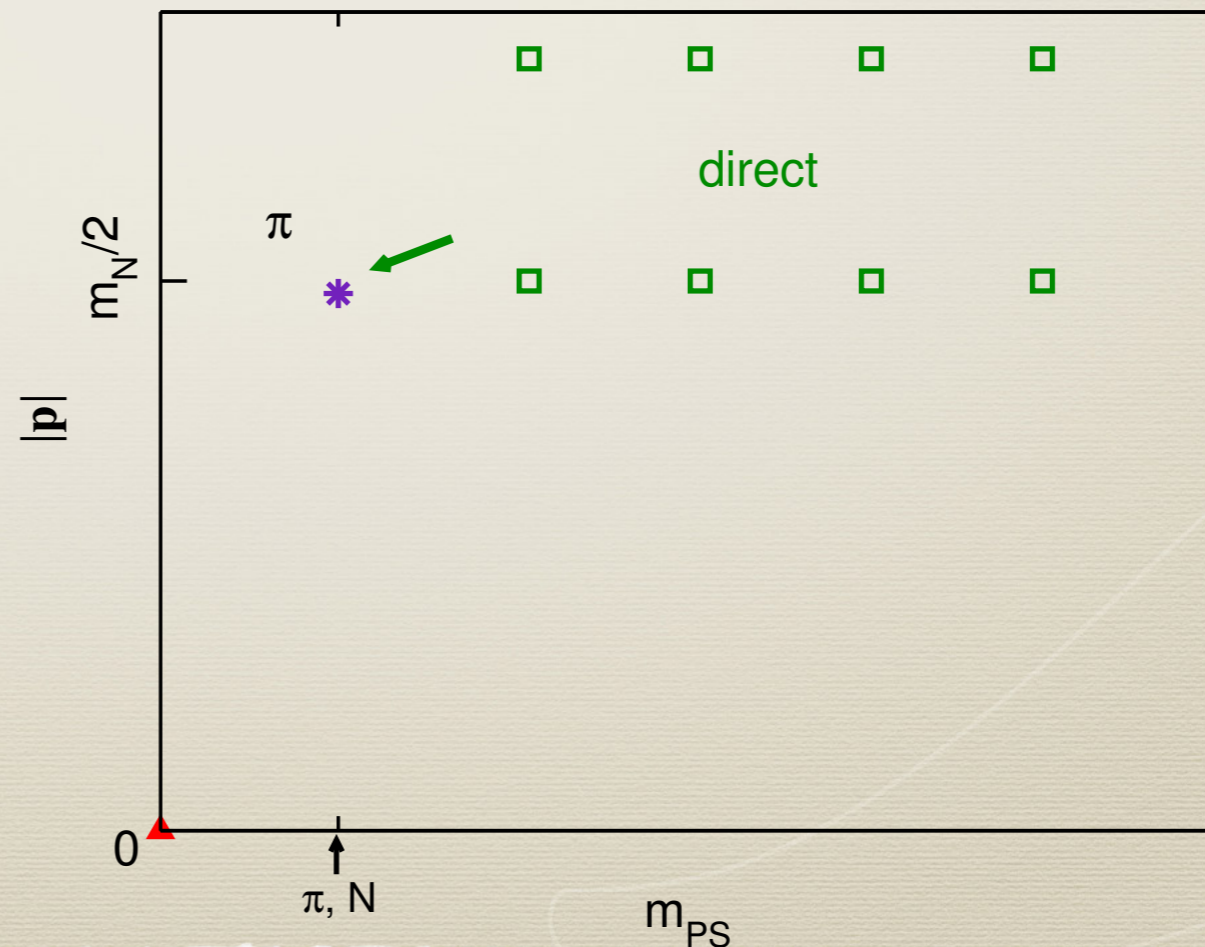
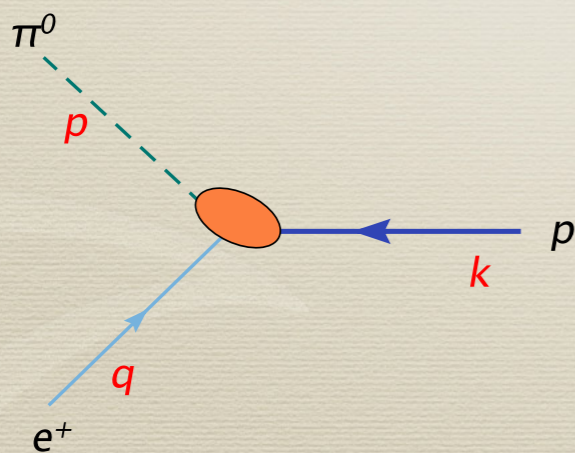
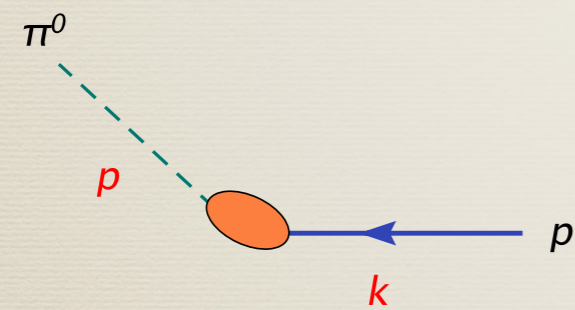
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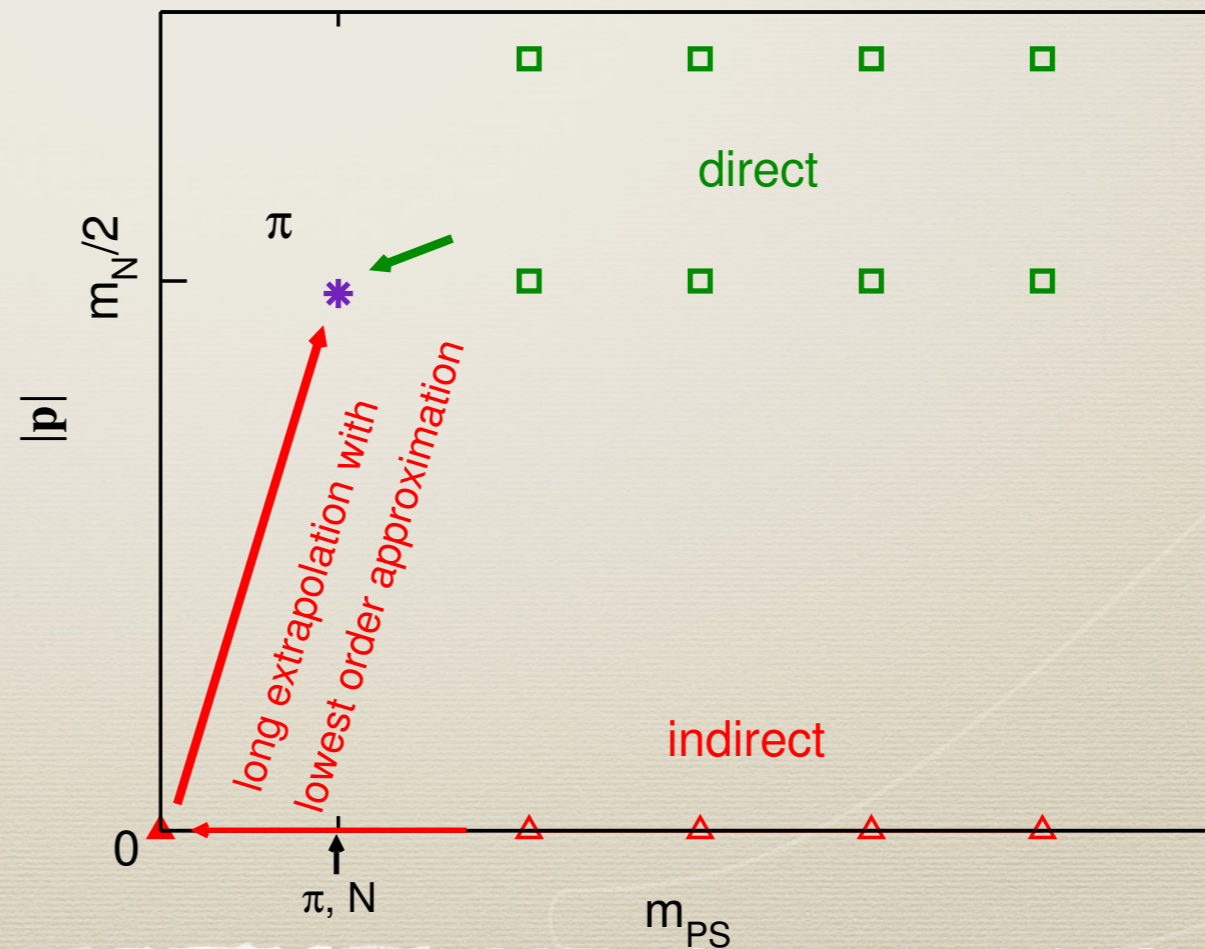
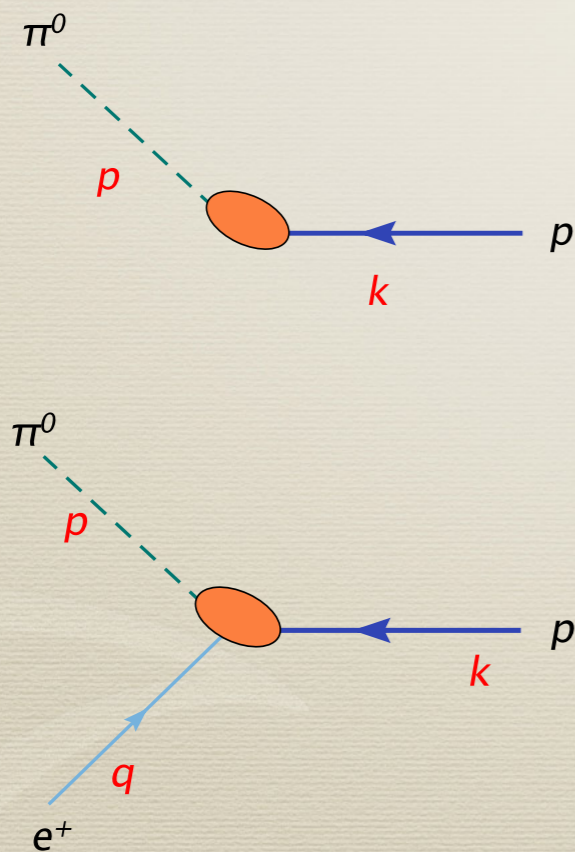
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how to calculate LEC's: α and β

* $\langle 0 | O_{RL} | p \rangle = \alpha u_p$

$O_{RL} = (\bar{u}^c P_R d) \cdot P_L u$ (color indices contracted with ϵ^{ijk} to make singlet)

$J_p = (\bar{u}^c \gamma_5 d) \cdot u$:proton interpolation operator

*
$$\sum_{\vec{x}} \langle 0 | O_{RL}(\vec{x}, t) \cdot \bar{J}_p(0) | 0 \rangle = \sum_{\vec{x}} \langle 0 | O_{RL}(\vec{x}, t) \cdot \sum_i | i \rangle \frac{1}{2E_i} \langle i | \cdot \bar{J}_p(0) | 0 \rangle$$

$$= \sum_i e^{-E_i t} \frac{1}{2E_i} \langle 0 | O_{RL}(0) | i \rangle \langle i | \bar{J}_p(0) | 0 \rangle$$

(large t) $\rightarrow e^{-m_p t} \frac{1}{2m_p} \langle 0 | O_{RL}(0) | p \rangle \langle p | \bar{J}_p(0) | 0 \rangle$

*
$$\sum_{\vec{x}} \langle 0 | J_p(\vec{x}, t) \cdot \bar{J}_p(0) | 0 \rangle \rightarrow e^{-m_p t} \frac{1}{2m_p} \langle 0 | J_p(0) | p \rangle \langle p | \bar{J}_p(0) | 0 \rangle$$

* linear combination of products of 3 quark propagators

* quark propagators: inverse of domain-wall fermion Dirac operator

* engineering the interpolation operator necessary to have good S/N

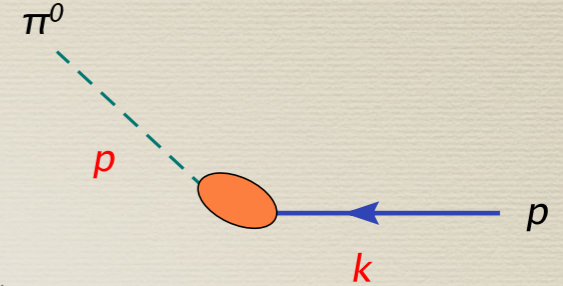
Direct method to calculate W_0

$$* \quad \langle \pi^0 | O_{RL} | p \rangle = P_L [W_0 - \frac{i \not{q}}{m_p} W_1] u_p$$

$$O_{RL} = (\bar{u}^c P_R d) \cdot P_L u$$

$$J_p = (\bar{u}^c \gamma_5 d) \cdot u \quad \text{:proton interpolation operator}$$

$$J_\pi^0 = \frac{1}{\sqrt{2}} (\bar{u} \gamma_5 u - \bar{d} \gamma_5 d) \quad \text{:pion interpolation operator}$$



* *three point function with momentum injection to pion in proton's rest frame*

$$\sum_{\vec{y}} \sum_{\vec{x}} e^{i\vec{p} \cdot (\vec{y} - \vec{x})} \langle 0 | J_{\pi^0}(\vec{y}, t') \cdot O_{RL}(\vec{x}, t) \cdot \bar{J}_p(0) | 0 \rangle$$

$$(t' \gg t \gg 0) \rightarrow e^{-E_\pi(t' - t)} e^{-m_p t} \frac{1}{2E_\pi} \frac{1}{2m_p} \langle 0 | J_{\pi^0} | \pi^0(\vec{p}) \rangle \langle \pi^0(\vec{p}) | O_{RL}(0) | p \rangle \langle p | \bar{J}_p(0) | 0 \rangle$$

$$* \quad \sum_{\vec{x}} \langle 0 | J_p(\vec{x}, t) \cdot \bar{J}_p(0) | 0 \rangle \rightarrow e^{-m_p t} \frac{1}{2m_p} \langle 0 | J_p(0) | p \rangle \langle p | \bar{J}_p(0) | 0 \rangle$$

$$* \quad \sum_{\vec{x}} e^{i\vec{p} \cdot \vec{x}} \langle 0 | J_{\pi^0}(\vec{x}, t) \cdot J_{\pi^0}(0) | 0 \rangle \rightarrow e^{-E_\pi t} \frac{1}{2E_\pi} \langle 0 | J_{\pi^0}(0) | \pi^0(\vec{p}) \rangle \langle \pi^0(\vec{p}) | J_{\pi^0}(0) | 0 \rangle$$

* *through some projection/subtraction, W_0 is obtained.*

Operator Property

- Operator $qqql$:

- Lorentz symmetry: $(\bar{q}^c \Gamma q)(\bar{l}^c \Gamma' q)$
- $SU(3)_c$ singlet: $\epsilon_{ijk} q^i q^j q^k l$
- $SU(2) \times U(1)$ symmetry determines the relative coefficients of the operators in low energy Lagrangian.
- relevant for nucleon decay: $q = u, d, s$ ($m_c > m_N$).
- at QCD scale, lepton is treated trivially, so we are left with

$$\mathcal{O} = \epsilon_{ijk} (u^{iT} C \Gamma d^j) \Gamma' s^k \quad \text{Lorentz spinor}$$

(u, d, s) simply labeling different flavors, not necessarily mean real flavors

- $\Gamma \Gamma'$ Lorentz structure variation with fixed flavor ordering:

- notation: $S = 1, P = \gamma_5, V = \gamma_\mu, A = \gamma_\mu \gamma_5, T = \sigma_{\mu\nu}$:
- $\mathcal{P}^-: SS, PP, AA, VV, TT.$
- $\mathcal{P}^+: SP, PS, AV, VA, T\tilde{T}.$

Operator

$$(\Gamma\Gamma')_{uds} = \epsilon_{ijk}(u^{iT} C \Gamma d^j) \Gamma' s^k$$

• \mathcal{P}^- operators

- $(SS)_{uds}, (SS)_{dsu}, (SS)_{sud}$

- $(PP)_{uds}, (PP)_{dsu}, (PP)_{sud}$

- $(AA)_{uds}$

- $(VV)_{uds}$

- $(TT)_{uds}$

- S : switching ($u \leftrightarrow d$): $(\Gamma\Gamma')_{dus} = \pm(\Gamma\Gamma')_{uds}$

- Any vector, tensor indices can be eliminated.

- $\Gamma(\Gamma') = S, P$ or L, R do everything:

common form in the low energy effective Lagrangian

→ Weinberg PRL 43 (1979) 1566.

- $(SS)_{uds} + (SS)_{dsu} + (SS)_{sud} + (PP)_{uds} + (PP)_{dsu} + (PP)_{sud} = 0$

- $(\Gamma\Gamma)_{uds}$ with (SS, PP, AA, VV, TT) can be used as a complete set:

We use them for the operator renormalization.

- Similar property for \mathcal{P}^- .

Renormalization: mixing

	$\Gamma \Gamma'$	
	\mathcal{S}^-	\mathcal{S}^+
\mathcal{P}^-	SS, PP, AA	VV, TT
\mathcal{P}^+	SP, PS, AV	$VA, T\tilde{T}$

$$\begin{pmatrix} SS \\ PP \\ AA \end{pmatrix}_{ren} = \begin{pmatrix} Z'_{ND}{}^{ij} \end{pmatrix} \begin{pmatrix} SS \\ PP \\ AA \end{pmatrix}_{latt}, \quad \begin{pmatrix} SP \\ PS \\ -AV \end{pmatrix}_{ren} = \begin{pmatrix} Z'_{ND}{}^{ij} \end{pmatrix} \begin{pmatrix} SP \\ PS \\ -AV \end{pmatrix}_{latt}$$

$$\begin{pmatrix} R \cdot L \\ L \cdot L \\ A \cdot LV \end{pmatrix}_{ren} = \begin{pmatrix} Z_{ND}{}^{ij} \end{pmatrix} \begin{pmatrix} R \cdot L \\ L \cdot L \\ A \cdot LV \end{pmatrix}_{latt}, \quad \begin{pmatrix} R \cdot R \\ L \cdot R \\ A \cdot RV \end{pmatrix} : \text{similar.}$$

• $(\Gamma\Gamma')_{udu} \equiv \epsilon_{ijk} (\bar{u}^{ci} \Gamma d^j) \Gamma' u^k$ is renormalized in the same way.

• Wilson fermion (Richards, Sachrajda, Scott):

$$(RL)_{ren} = Z(RL)_{latt} + \frac{\alpha_s}{4\pi} Z_{mix}(LL)_{latt} + \frac{\alpha_s}{4\pi} Z'_{mix}(A \cdot LV)_{latt}$$

no other terms appear at any order.

RI/MOM scheme renormalization

$$G^a(x_0, x_1, x_2, x_3) = \langle \mathcal{O}_{uds}^a(x_0) \bar{u}(x_1) \bar{d}(x_2) \bar{s}(x_3) \rangle.$$

- a labels chiral structure type $\Gamma\Gamma'$.

$$\mathcal{O}_{uds\ ren}^a = Z_{ND}^{ab} \mathcal{O}_{uds\ latt}^b$$

- Momentum p for all three external quarks, amputate it with three quark propagators:

$$\Lambda^a(p^2) = \text{F.T. } G^a(0, x_1, x_2, x_3)|_{\text{Amp}}.$$

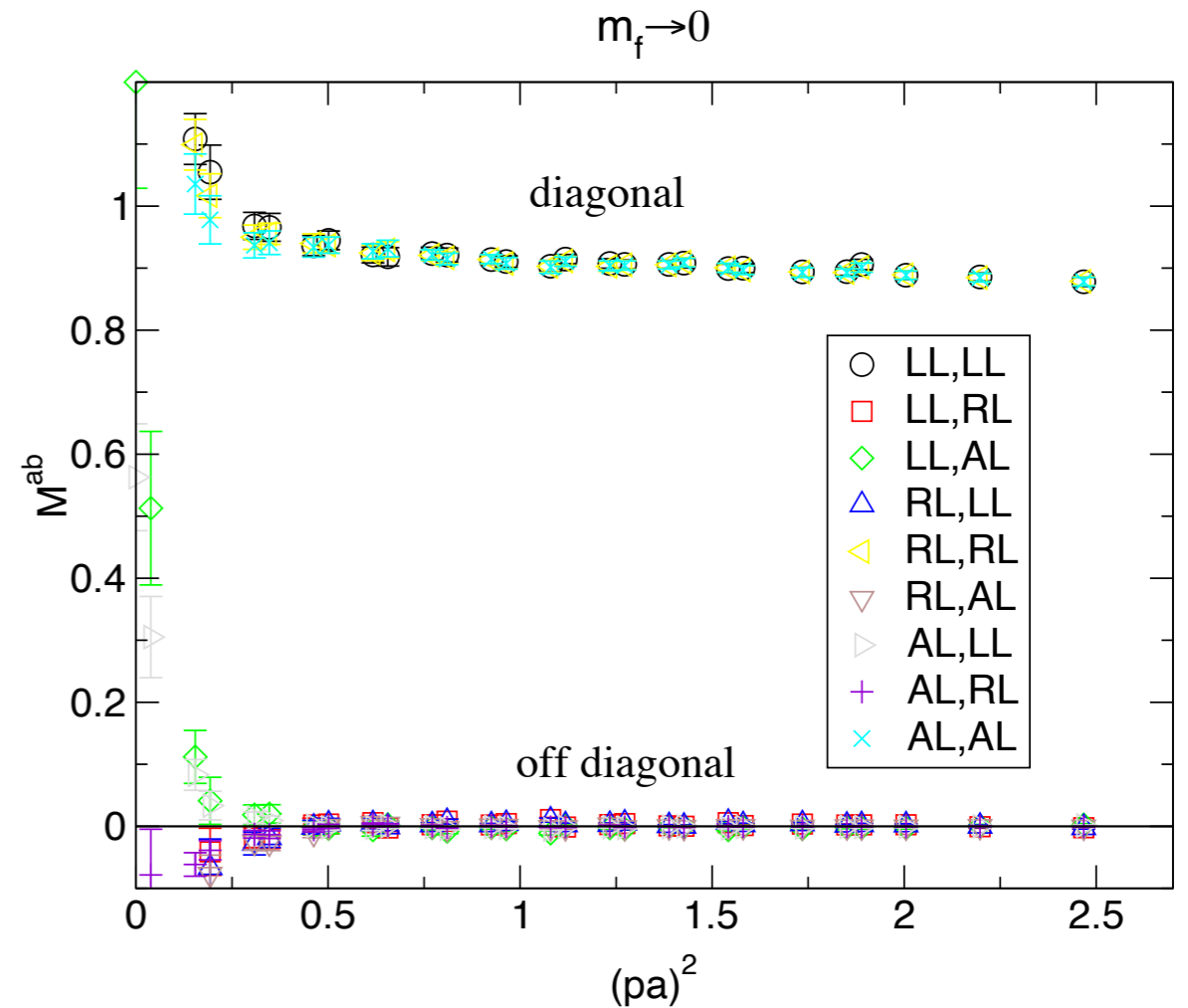
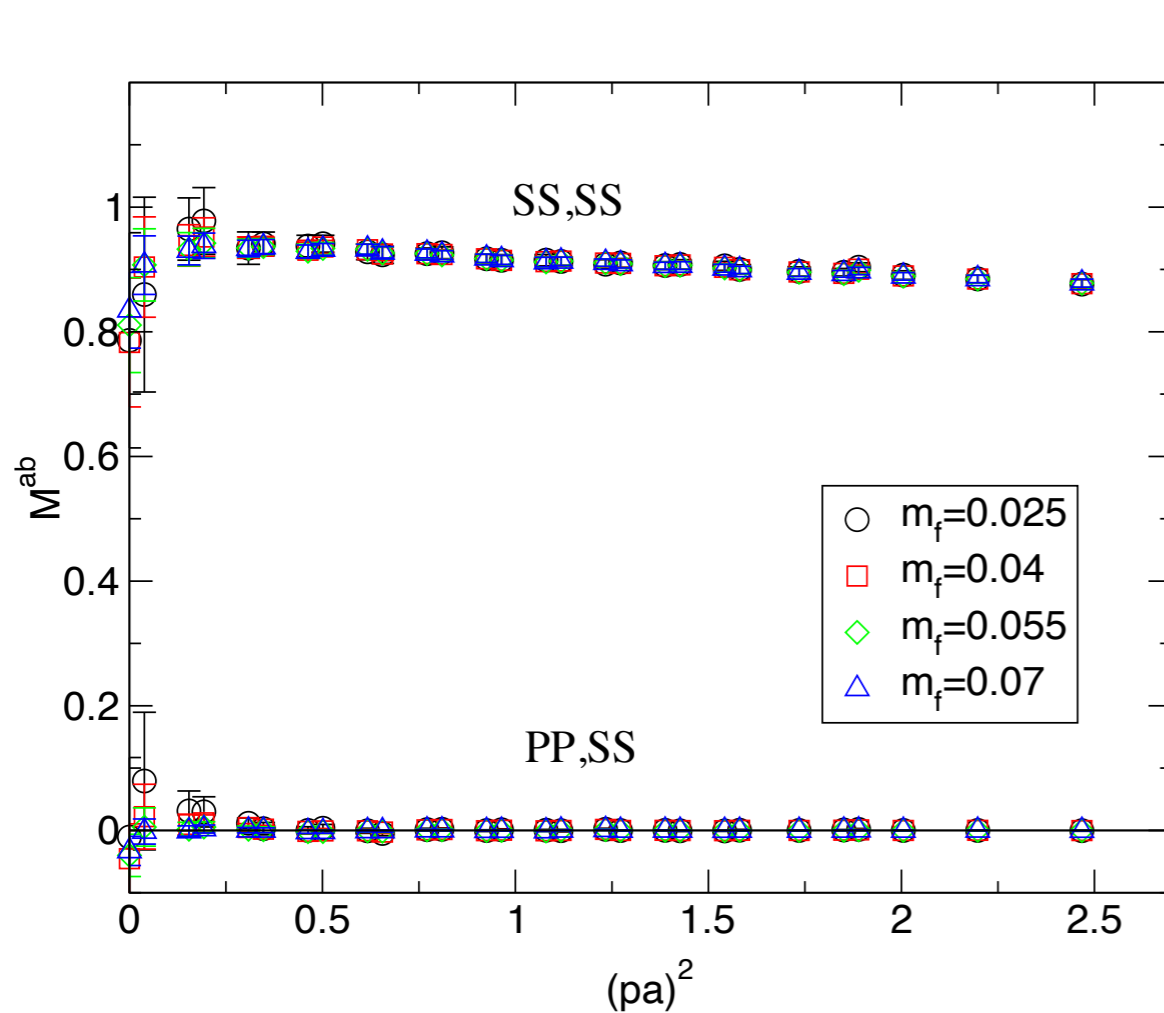
- Renormalization condition reads at scale p ,

$$P_{ijk\ \beta\alpha\ \delta\gamma}^a \cdot Z_q^{-3/2} Z_{ND}^{bc} \Lambda_{ijk\ \alpha\beta\ \gamma\delta}^c = \delta^{ab},$$

$$M^{ab} = P_{ijk\ \beta\alpha\ \delta\gamma}^a \cdot \Lambda_{ijk\ \alpha\beta\ \gamma\delta}^b \rightarrow Z_q^{3/2} (Z_{ND}^{-1})^{ab}.$$

Operator Mixing ?

$\beta = 0.87$ DBW2 glue ($a^{-1} = 1.3$ GeV), and DWF with $L_s = 12$ and $M_5 = 1.8$:
 $m_{\text{res}} \simeq 1.6 \text{ MeV}$.



● Small mass dependence.

● Off-diagonal elements negligible.
 → No mixing.

Treatment of Z_q

$$M^{LL,LL} \rightarrow Z_q^{3/2} / Z_{ND}^{LL,LL}.$$

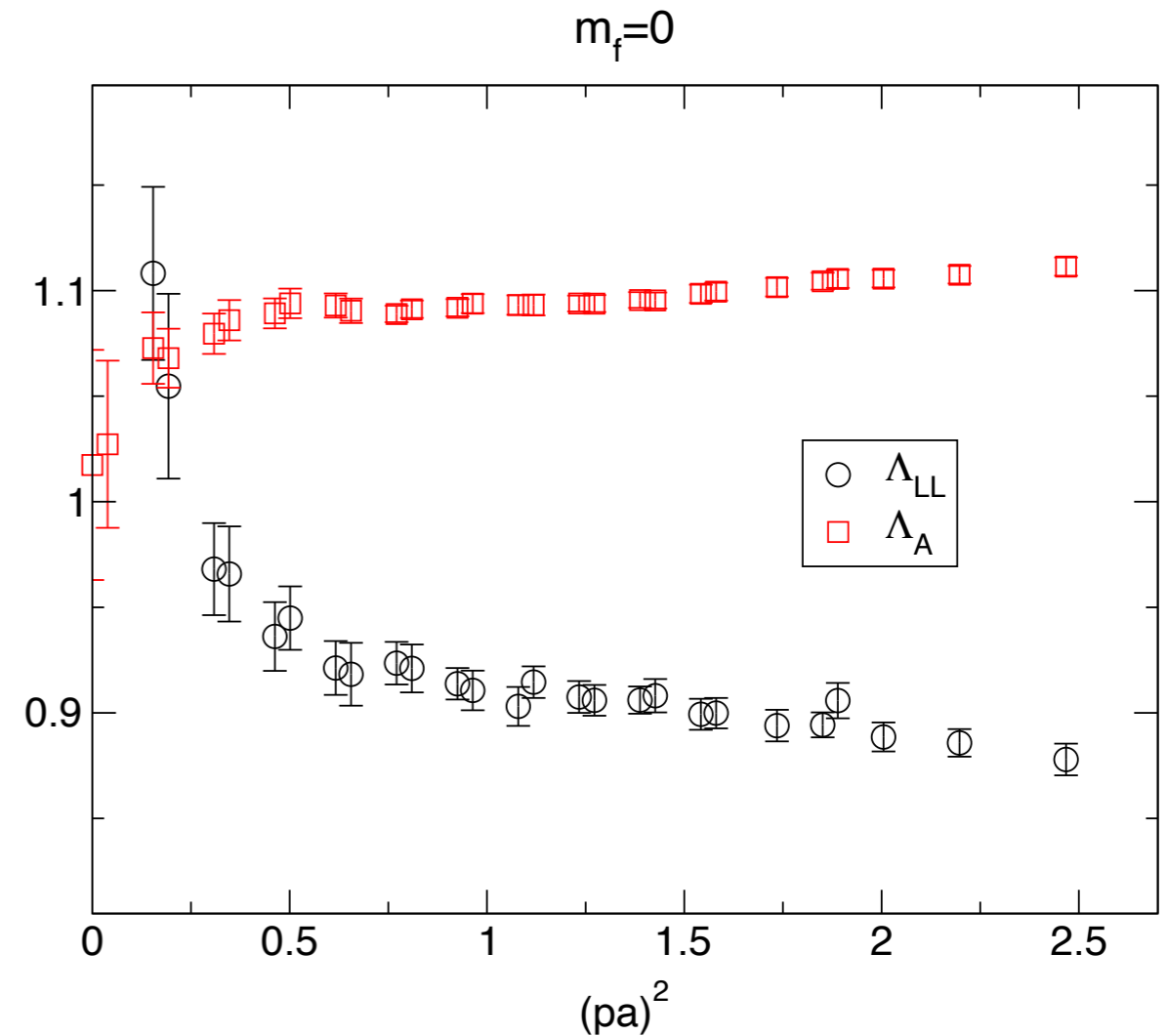
$$P_A \Lambda_A \rightarrow Z_q Z_A^{-1}.$$

$$(P_A \Lambda_A)^{3/2} / M^{LL,LL} \rightarrow Z_{ND}^{LL,LL} / Z_A^{3/2}.$$

$$Z_A = 0.7798(5) \leftarrow \frac{\langle \mathcal{A}_\mu P \rangle}{\langle A_\mu P \rangle}$$

• Z_A has no scale dependence (Null anomalous dimension for A_μ).

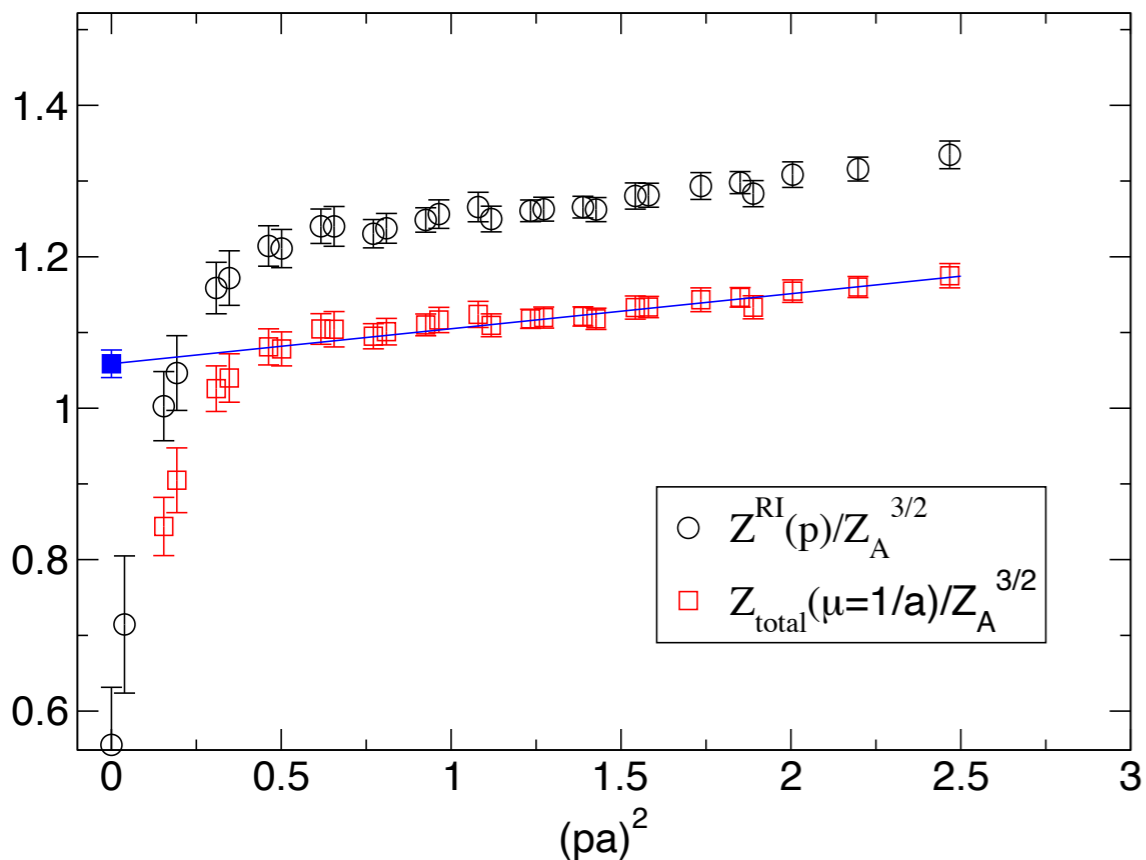
• $(P_A \Lambda_A)^{3/2} / M^{LL,LL}$ has same scale dependence as Z_{ND} and possible $O(a^2)$ discretization error.





Matching to \overline{MS} (NLO)

LL



$$\mathcal{O}^{\overline{MS}}(\mu) = \underbrace{U^{\overline{MS}}(\mu; p) \frac{Z^{\overline{MS}}(p)}{Z^{RI}(p)}}_{Z_{total}(\mu)} Z^{RI}(p) \mathcal{O}^{latt}$$

- 2-loop anomalous dimension: Nihei, Arafune (94)
- 1-loop matching (finite part): This work
- NPR: This work
- The product is independent of p . Let's set $\mu = 1/a$.

$$Z_A = 0.7798(5) \quad \rightarrow \quad Z_{total}^{LL,LL}(\mu = 1/a) = 0.73(1).$$

$\beta = 0.87$ DBW2 glue ($a^{-1} = 1.3$ GeV), and DWF with $L_s = 12$ and $M_5 = 1.8$:



Recent improvement in RI/MOM renormalization

- issues

- p must be large enough, to control perturbation theory
- p must be small enough, to control lattice artifact $(pa)^n$
- $\Lambda_{QCD} \ll p \ll 1/a$ Window problem

- solution

- super fine lattice
- perturbation theory: higher order
- scheme: less contamination of low energy physics

NLO → NNLO

RI/SMOM

Proton Decay Matrix Elements projects

Collaboration in Japan

- **Eigo Shintani** (Tsukuba)
- **Ryutaro Tsuji** (KEK)
- Yoshinobu Kuramashi (Tsukuba)
- YA

using PACS Wilson configurations

$$V=64^4 \quad L=5.5 \text{ fm}$$

$$1/a=2.3 \text{ GeV}$$

$$a = 0.08 \text{ fm}$$

Program for Promoting Researches
on the Supercomputer Fugaku

Large-scale lattice QCD simulation
and development of AI technology

Collaboration in USA

- **Jun-Sik Yoo** (KEK/Stony Brook)
- Taku Izubuchi (BNL/RBRC)
- Amarjit Soni (BNL)
- **Sergey Syritsyn** (Stony Brook/RBRC)
- YA

using RBC/UKQCD DWF configurations

$$V=24^3 \times 64, 32^3 \times 64 \quad L=4.7 \text{ fm}$$

$$1/a=1 \text{ and } 1.4 \text{ GeV}$$

$$a = 0.2, 0.14 \text{ fm} \rightarrow 0$$

[PRD 105, 0074501 (2022)]

use of **PACS** ensemble @ physical pion mass

$N_f=2+1$ PACS ensemble

- Iwasaki gauge $\beta=1.82$
- **stout smeared Wilson fermion**: $\rho=0.1$, $N=6$
- ud and s quarks are on **physical point**
- $1/a = 2.333(18)$ GeV
- 64^4 is mostly used in this study: $m_\pi L=3.8$

statistical note

- ~100 configurations
- for each config
 - matrix elements: **AMA**
 - one exact and
 - 256 sloppy solves
 - NPR:
 - single point source

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Renormalization:

Richards et al 1987 lattice perturbation theory
(Notation JLQCD2000)

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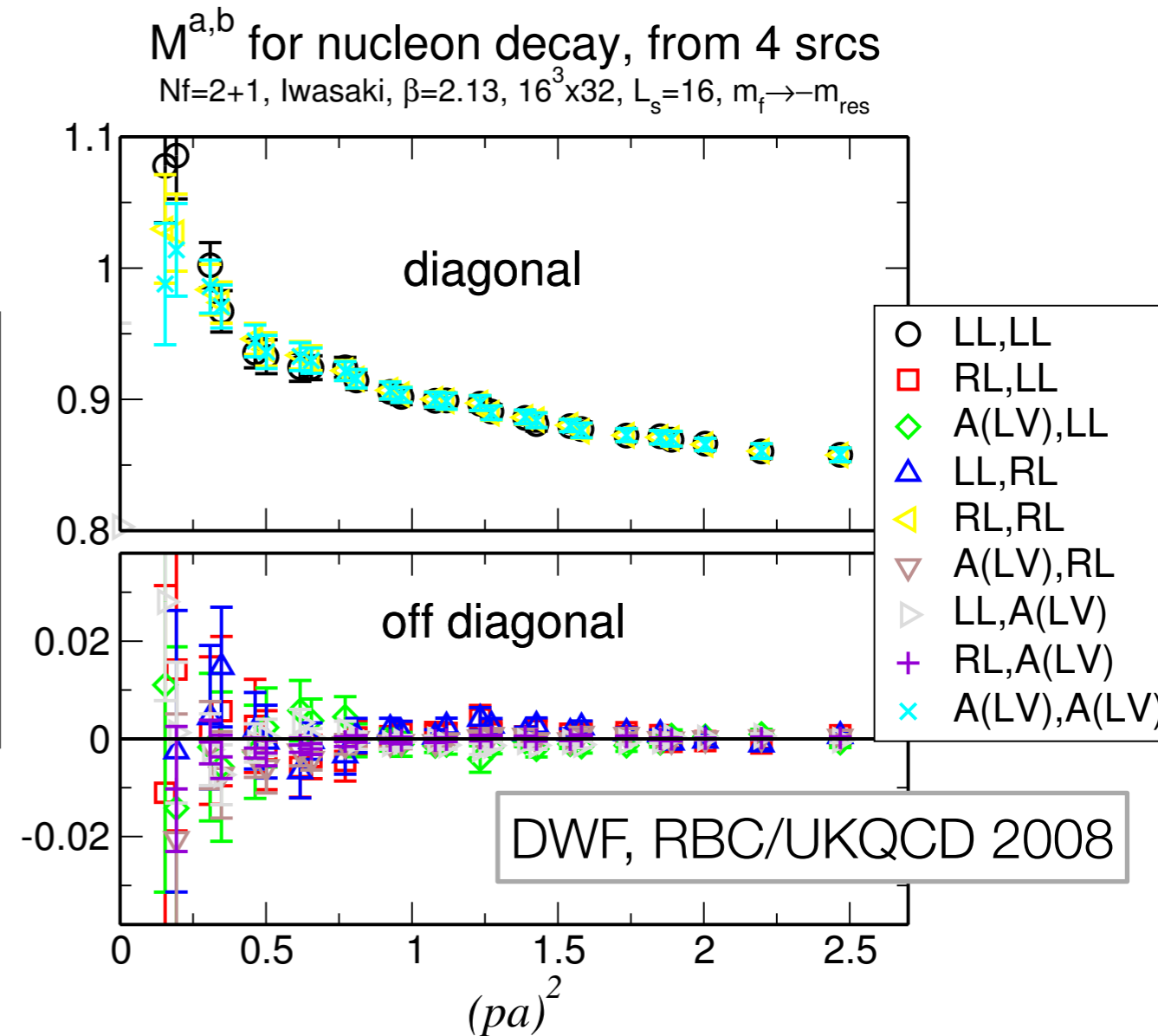
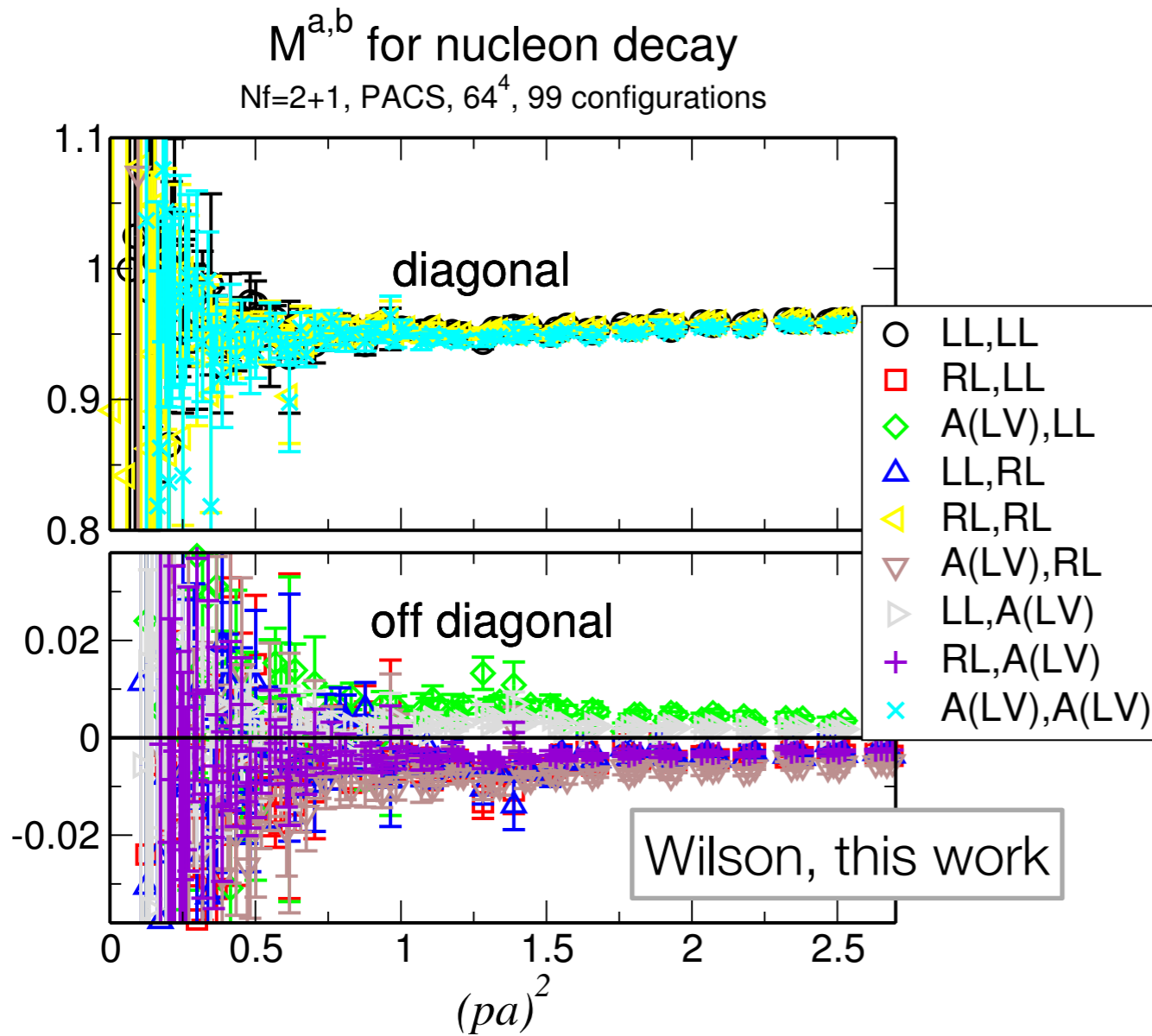
Renormalization:

Richards et al 1987 lattice perturbation theory
(Notation JLQCD2000)

$$\begin{aligned}\mathcal{O}_{RL}^{\text{cont}}(\mu) &= Z(\alpha_s, \mu a) \mathcal{O}_{RL}^{\text{latt}}(a) + \frac{\alpha_s}{4\pi} Z_{\text{mix}} \mathcal{O}_{LL}^{\text{latt}}(a) \\ &\quad - \frac{\alpha_s}{4\pi} Z'_{\text{mix}} \mathcal{O}_{\gamma_\mu L}^{\text{latt}}(a),\end{aligned}\tag{47}$$

$$\begin{aligned}\mathcal{O}_{LL}^{\text{cont}}(\mu) &= Z(\alpha_s, \mu a) \mathcal{O}_{LL}^{\text{latt}}(a) + \frac{\alpha_s}{4\pi} Z_{\text{mix}} \mathcal{O}_{RL}^{\text{latt}}(a) \\ &\quad + \frac{\alpha_s}{4\pi} Z'_{\text{mix}} \mathcal{O}_{\gamma_\mu L}^{\text{latt}}(a),\end{aligned}\tag{48}$$

RI/MOM 3q vertex matrix: comparison with DWF



• $(pa)^2$ dep bit different \rightarrow lattice artifact dominant

• Wilson:

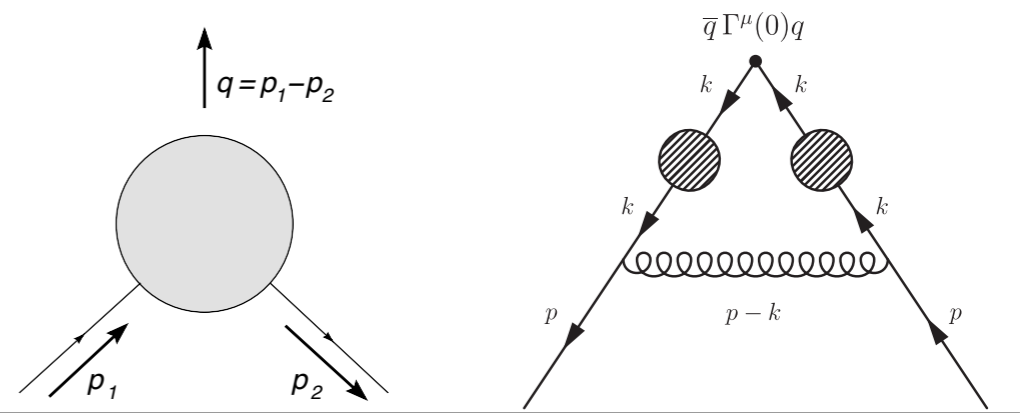
$$\begin{cases} O_{RL} &= (\bar{u}^c P_R d) \cdot P_L s \\ O_{LL} &= (\bar{u}^c P_L d) \cdot P_L s \\ O_{A(LV)} &= (\bar{u}^c \gamma_\mu \gamma_5 d) \cdot P_L \gamma_\mu s \end{cases}$$

off-diagonal larger than DWF, but, $\approx 1\%$ \rightarrow treated as negligible below

renormalization

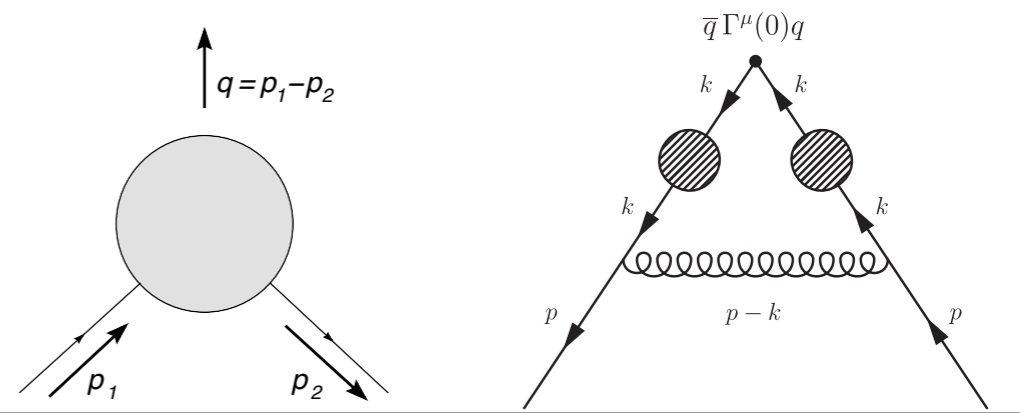
- ratio of 3q, 2q vertex function $\rightarrow Z_{ND}/Z_A^{3/2}, Z_{ND}/Z_V^{3/2}$
 - input Z_V or Z_A from SF scheme $\rightarrow Z_{PD}$

renormalization



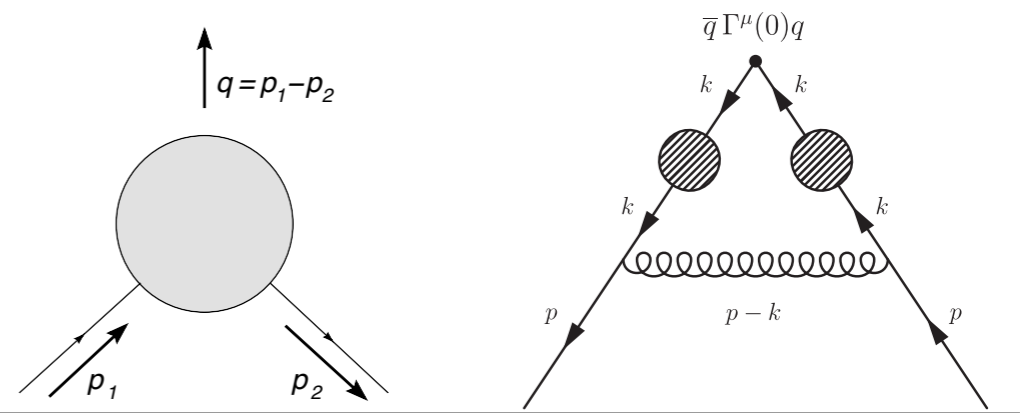
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- DWF (RBC/UKQCD) [2006~2017]
 - wave function renormalization (bilinear) uses exceptional momenta ($q_\mu=0$)
 - sensitive to the change of m_f and SSB \rightarrow source of systematic error

renormalization



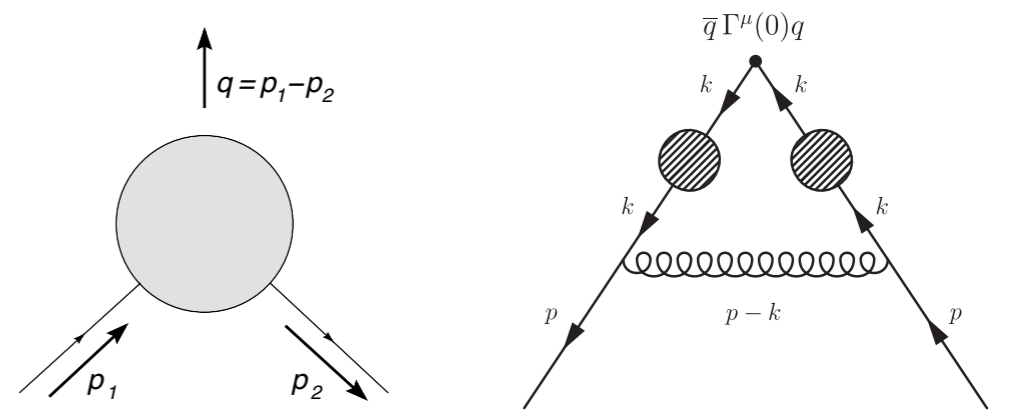
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- utilize RI/**SMOM** wave function renormalization: A_μ, V_μ
 - ➔ proton decay **SMOM** schemes: SMOM, SMOM $_{\gamma_\mu}$
 - had expected effect on bilinear and 4q operators
 - $\overline{\text{MS}}$ matching w/ NLO perturbation theory
 - totally: $2 \times 2 = 4$ schemes

renormalization



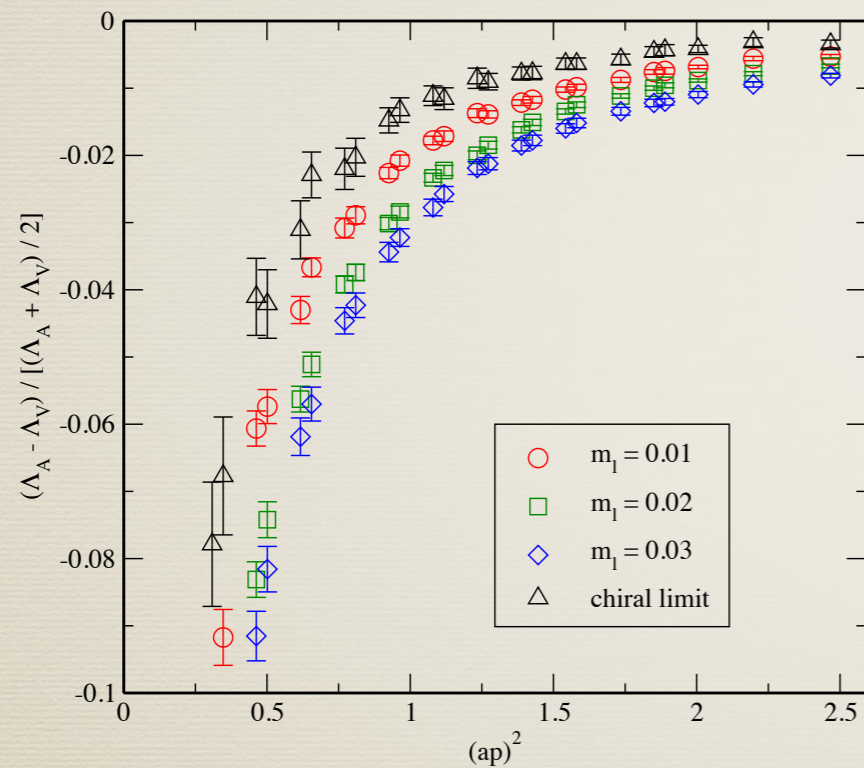
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SMOM schemes:
fully utilize
non-exceptional momenta

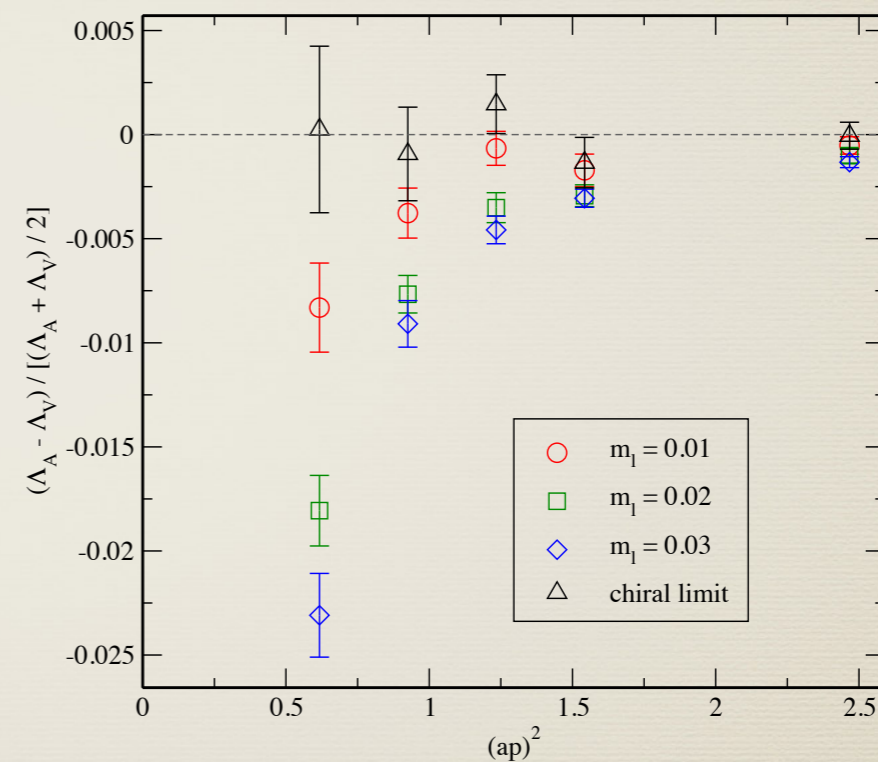
A test of non-exceptional mom

$\Lambda_A - \Lambda_V$: RBC/UKQCD [PRD 2008]

exceptional

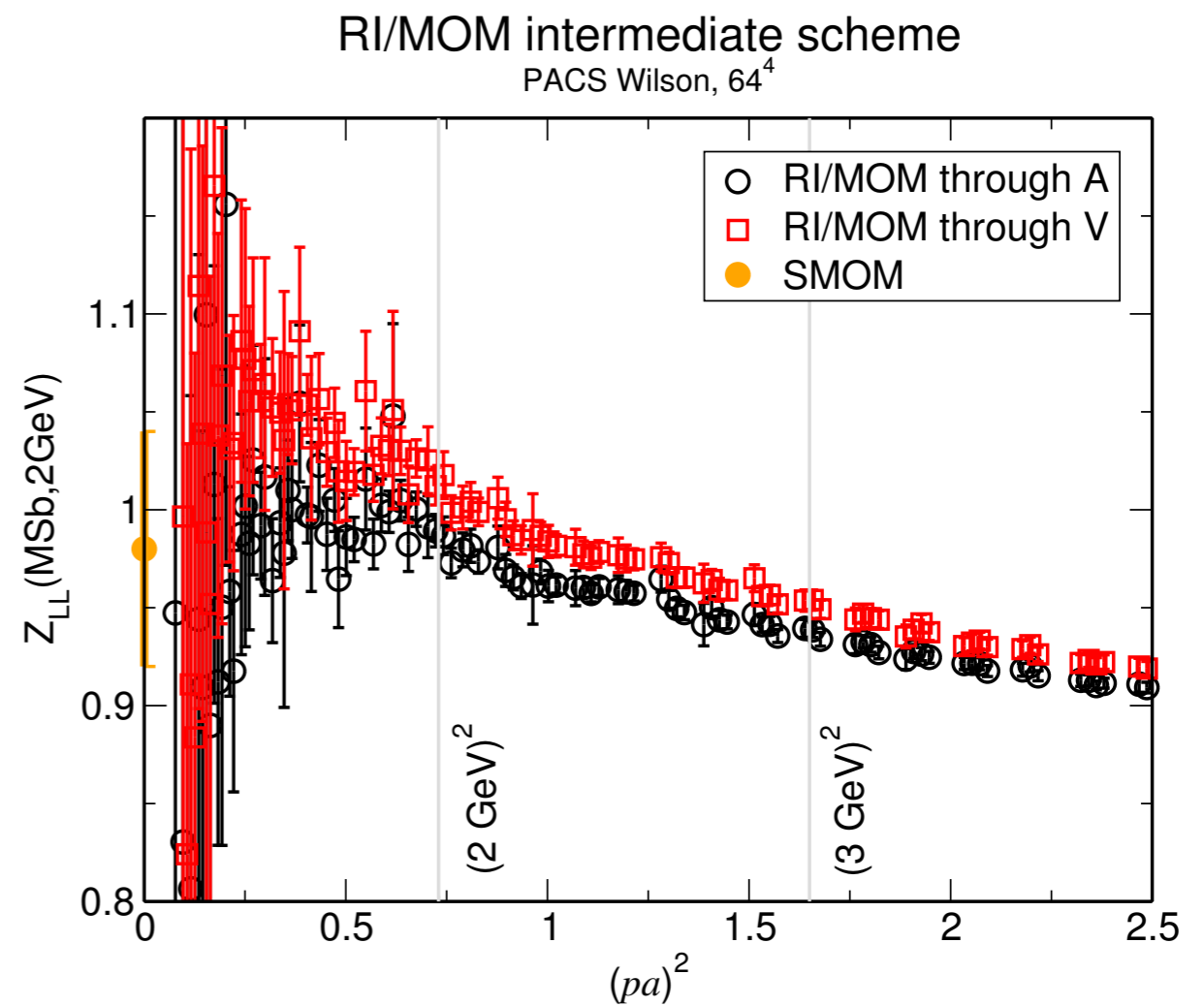
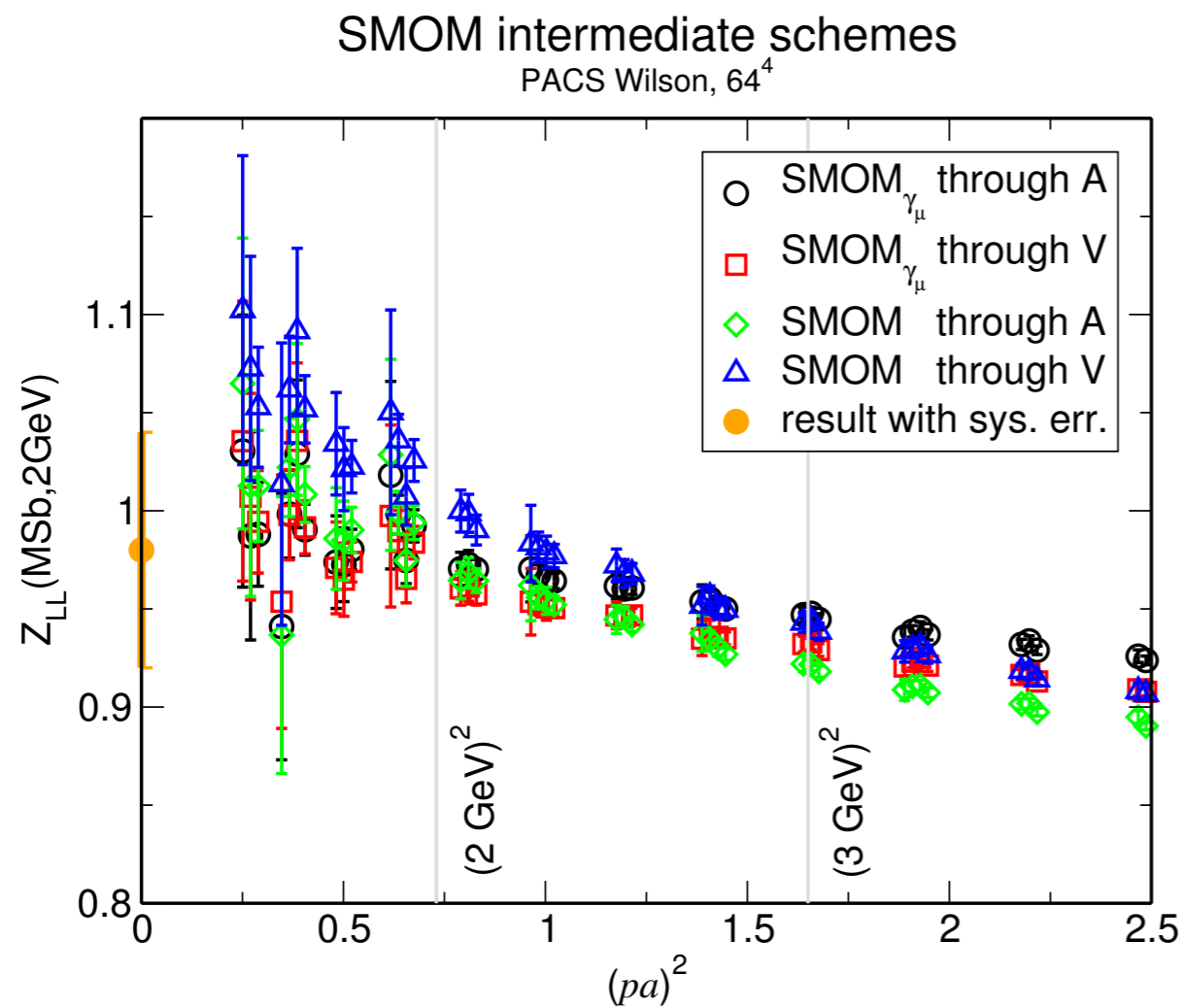


non-exceptional



* The success created a very good motivation to invest in non-exceptional momenta

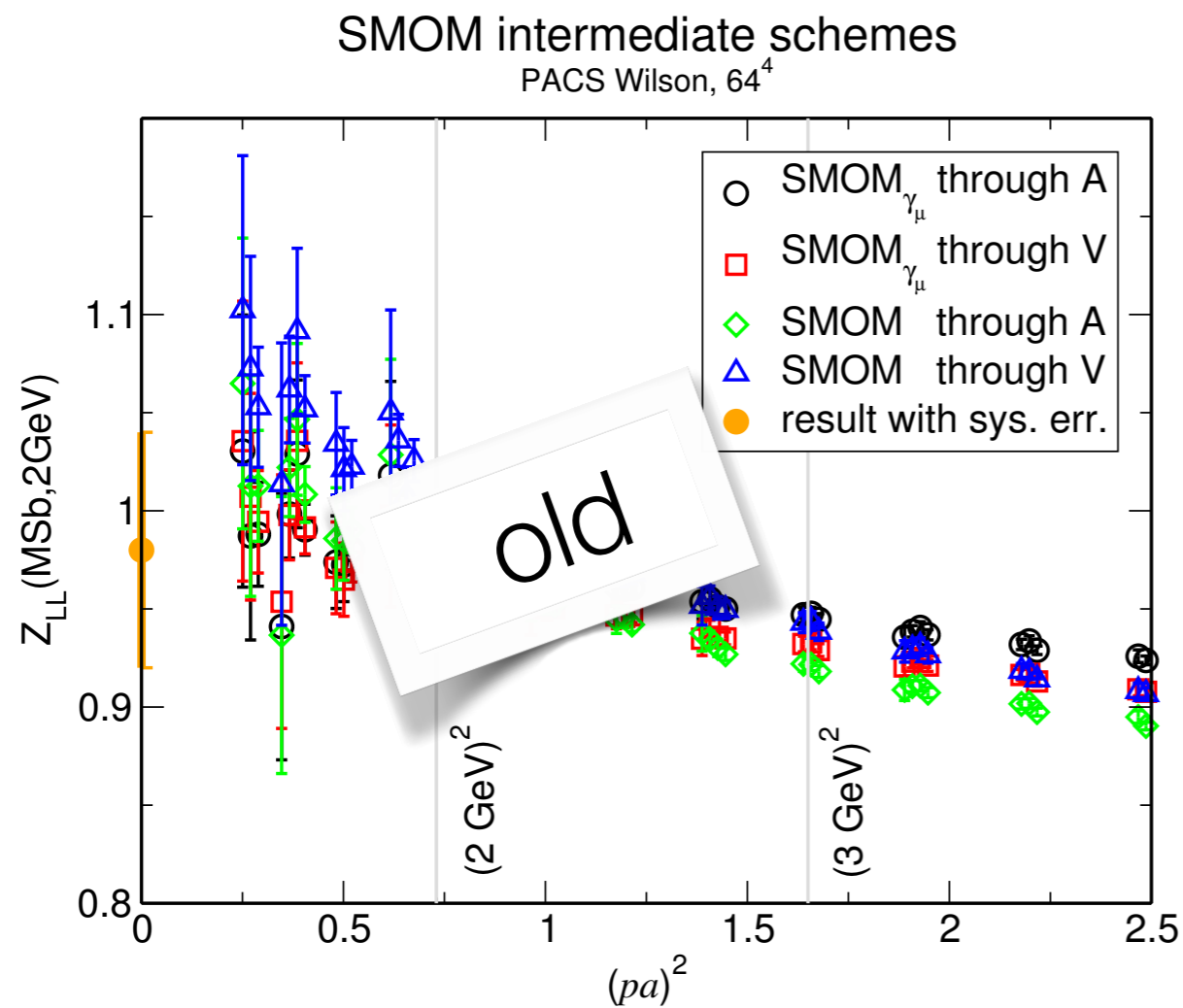
$\overline{\text{MS}}$ $Z(2\text{GeV})$ from RI/SMOM schemes ver.1



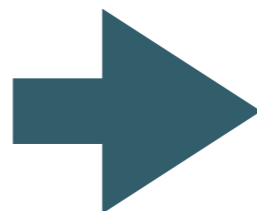
- $Z_{LL}(\text{MSb}, 2\text{GeV}) = 0.98 (6)$
- $Z_{RL}(\text{MSb}, 2\text{GeV}) = 0.98 (7)$

@ Lattice 2019

$\overline{\text{MS}}$ $Z(2\text{GeV})$ from RI/SMOM schemes improved



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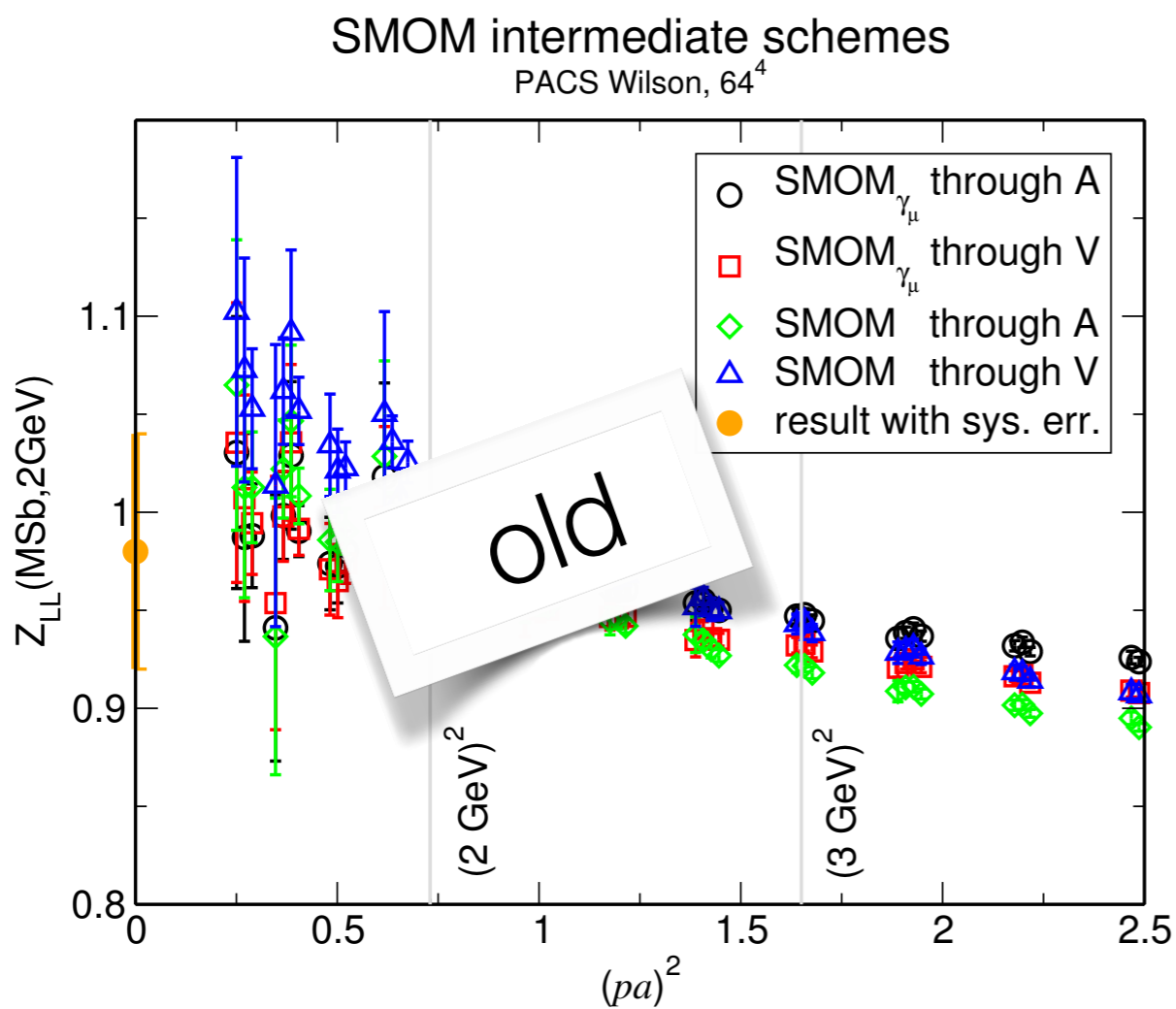
Improvement:

[Tsuji et al Lattice 2024]

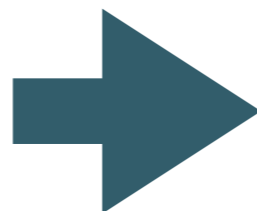
- (improved stat.)
- use of SYM3q scheme
 - ➔ NNLO available
- remove $(pa)^2$ and higher
- remove non-pert. eff.
 - ➔ fit variation
- estimate PT truncation
 - ➔ several interm. scheme
- $Z_{\overline{\text{MS}}}^{LL} = 1.018(6)_{\text{stat}}(37)_{\text{sys}}$
- $Z_{\overline{\text{MS}}}^{RL} = 1.016(5)_{\text{stat}}(41)_{\text{sys}}$

- note: $Z_{LL}(\text{MSb}, 2\text{GeV}) = Z_{RL}(\text{MSb}, 2\text{GeV}) \approx 1 \rightarrow$ bare ME \approx ren. ME

$\overline{\text{MS}}$ $Z(2\text{GeV})$ from RI/SMOM schemes improved



- $Z_{LL}(\text{MSb}, 2\text{GeV}) = 0.98 (6)$
- $Z_{RL}(\text{MSb}, 2\text{GeV}) = 0.98 (7)$



Improvement:

[Tsuji et al Lattice 2024]

- (improved stat.)
- use of SYM₂ scheme
- \rightarrow N^2L
- \rightarrow N^2L and higher
- remove non-pert. eff.
- \rightarrow fit variation
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NPR schemes and matching - utilizing RI/MOM

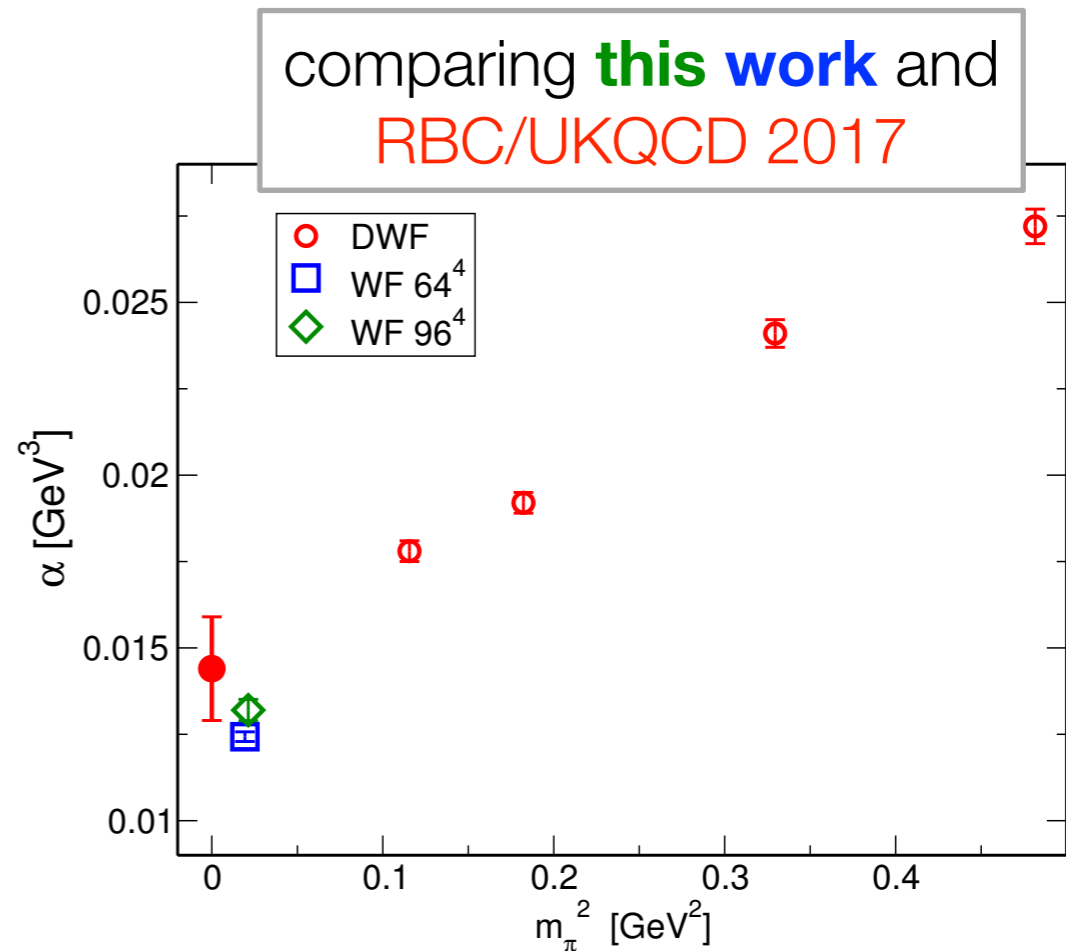
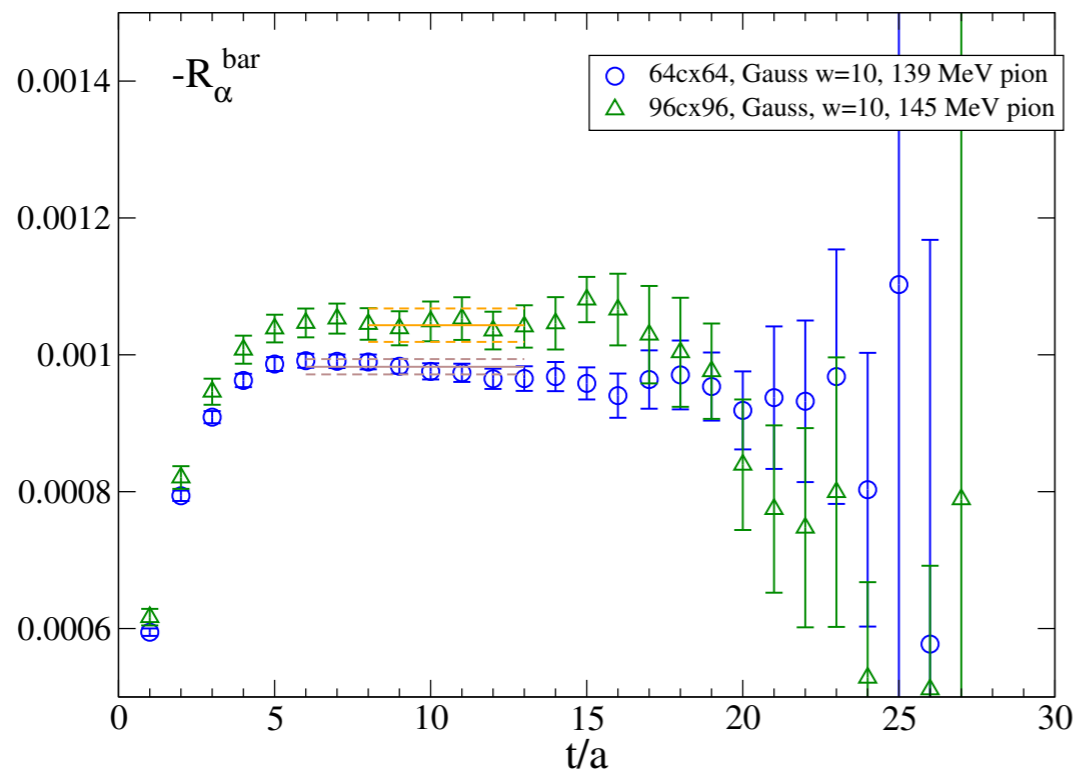
NPR schemes and matching - utilizing RI/MOM

- three quark vertex
 - original:
 - $p_1 = p_2 = p_3 = p, q = 3p, \mu^2 = p^2$ (non-exceptional) [RBC 2006] NLO
 - symmetric 3q:
 - $p_1 + p_2 + p_3 = q = 0, p_1^2 = p_2^2 = p_3^2 = p^2 = \mu^2$ (exceptional) [Gracey 2013] NNLO
 - SMOM
 - $p_1^2 = p_2^2 = p_3^2 = q^2 = \mu^2$ NEW! (yet) (non-exceptional) [Kniehl, Veretin 2023] NNLO
- bilinear (q-qbar) vertex for wave function renormalization
 - RI/MOM [Martinelli et al 1993, 1995]
 - $p_1 = p_2, q = 0$ (exceptional)
 - RI/SMOM (symmetric momentum configuration)
 - $p_1^2 = p_2^2 = q^2$ (non-exceptional) [RBC/UKQCD 2007]
 - two schemes: $SMOM$ and $SMOM_{\gamma_\mu}$ [Sturm et al 2009] NLO, [Almeida, Sturm 10] NNLO

proton decay LEC: α , β

Low energy constants of proton decay

- LO ChPT $W_0[\langle \pi^0 | (ud)_{R u_L} | p \rangle] \simeq \frac{\alpha}{\sqrt{2}f} (1 + D + F)$
- ratio of 2pt functions



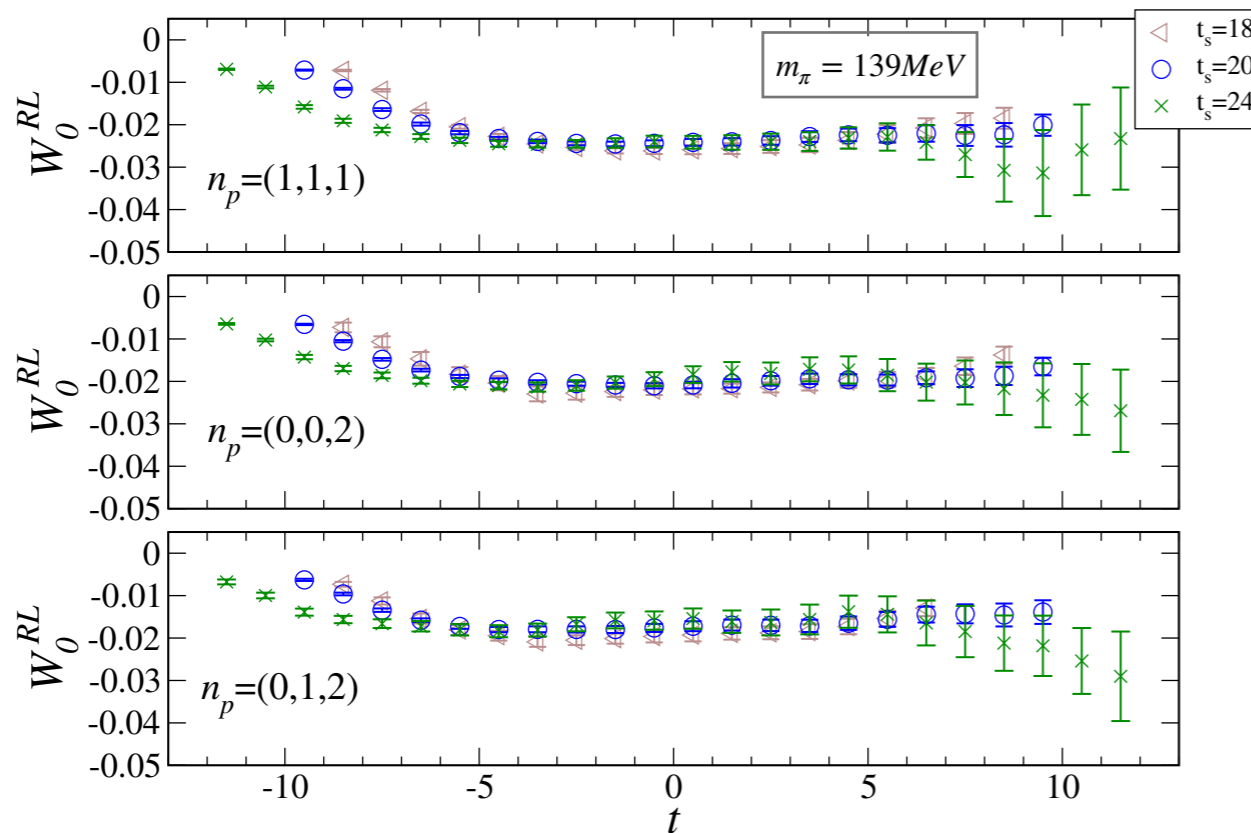
- α consistent with earlier DWF computation w/ long chiral extrapol.
- β as well
- no big surprise happening when going down to physical ud mass

proton decay form factor W_0 for pion final state

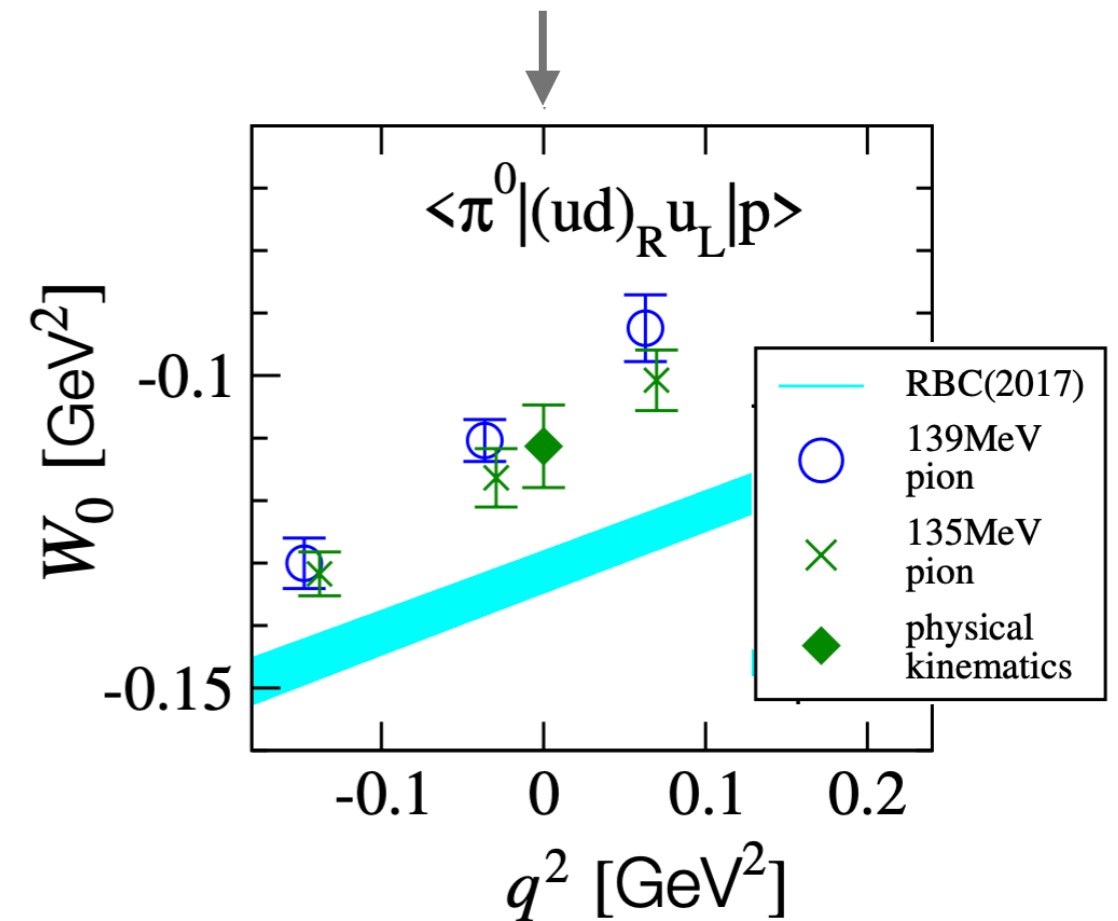
Relevant form factor $W_0 \langle \pi^0 | (ud)_R u_L | p \rangle$ as an example

- from ratio of 3 and 2 point functions

- meson momentum $\vec{p} = \frac{2\pi}{L} \vec{n}_p$



on-shell lepton: $-q^2 = m_l^2 = 0$



- $|W_0|$ ~20% smaller than DWF (with a long chiral extrapolation) at $q^2=0$
- consistent with sys. error ! no big surprise found for $m_f \rightarrow m_{ud}$
- 10% total error is not a dream...

RBC/UKQCD study

Phys. Rev. D 105, 074501 (2022)

Advantage

- New renormalization scheme (subtraction point)
 - Matching available one order higher (NNLO) : Gracey (2012)
 - → reduced systematic error ~1%
- Two lattice spacings → continuum limit
- Chiral symmetry

Disadvantage

- Coarse lattice $a=0.2,0.14\text{fm}$ → large error after continuum extrapolation ~20%

JUN-SIK YOO *et al.*

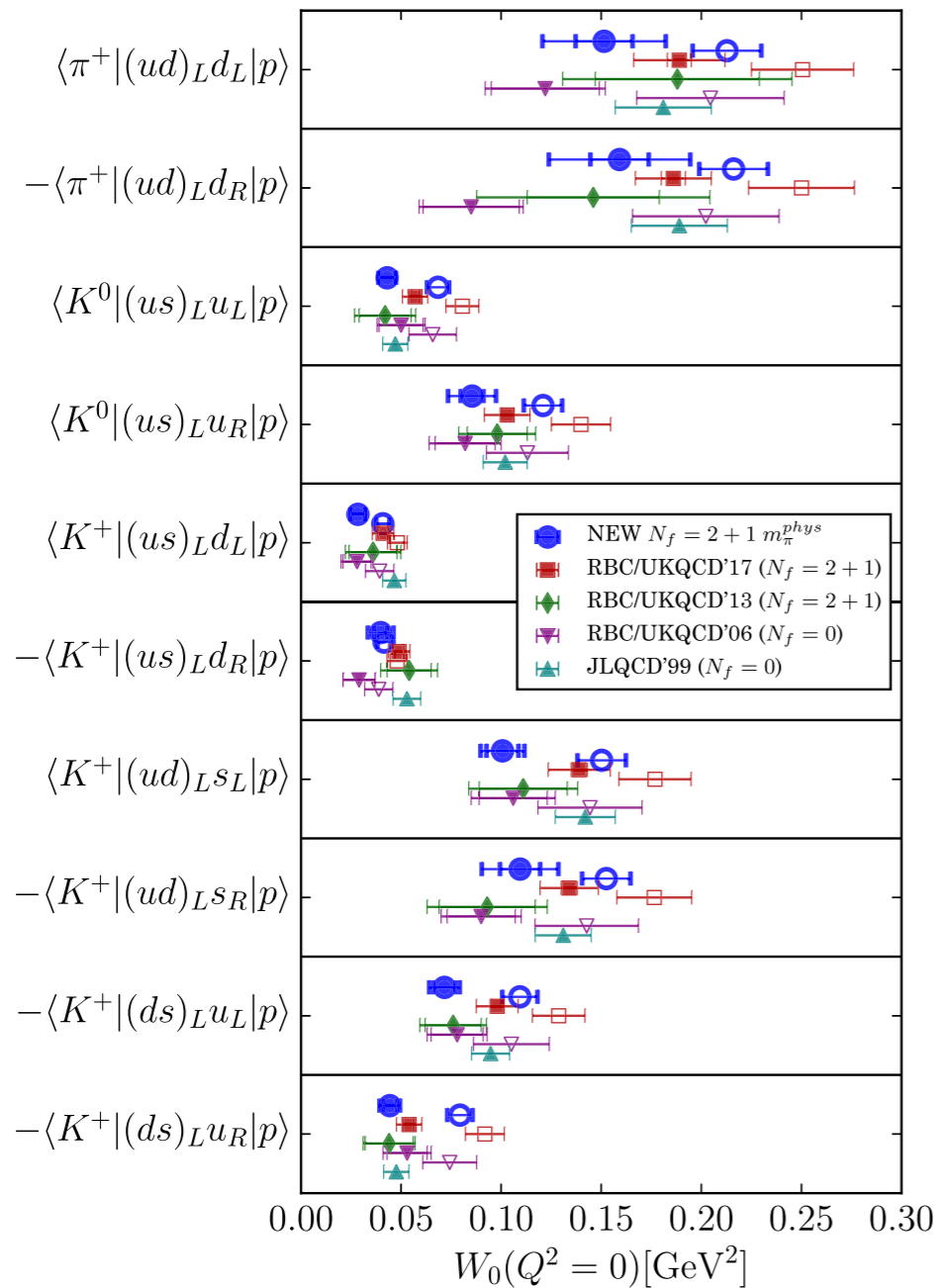
PHYS. REV. D **105**, 074501 (2022)

TABLE VIII. Results for the form factors $W_{0,1}$ on the two ensembles and in the continuum limit at the two kinematic points $Q^2 = 0$ (first line) and $Q = -m_\mu^2$ (second line) renormalized to $\overline{\text{MS}}(2 \text{ GeV})$. The first uncertainty is statistical, the second is systematic due to excited states, and the third is the uncertainty of the continuum extrapolation.

	$W_0[\text{GeV}^2]$		
	24ID	32ID	Cont.
$\langle \pi^+ (ud)_L d_L p \rangle$	0.1032(86)(26)	0.1252(48)(50)	0.151(14)(8)(26)
	0.1050(87)(36)	0.1271(49)(50)	0.153(14)(7)(26)
$\langle \pi^+ (ud)_L d_R p \rangle$	-0.1125(78)(41)	-0.134(5)(11)	-0.159(15)(20)(25)
	-0.1139(78)(45)	-0.136(5)(12)	-0.161(15)(20)(26)

RBC/UKQCD results

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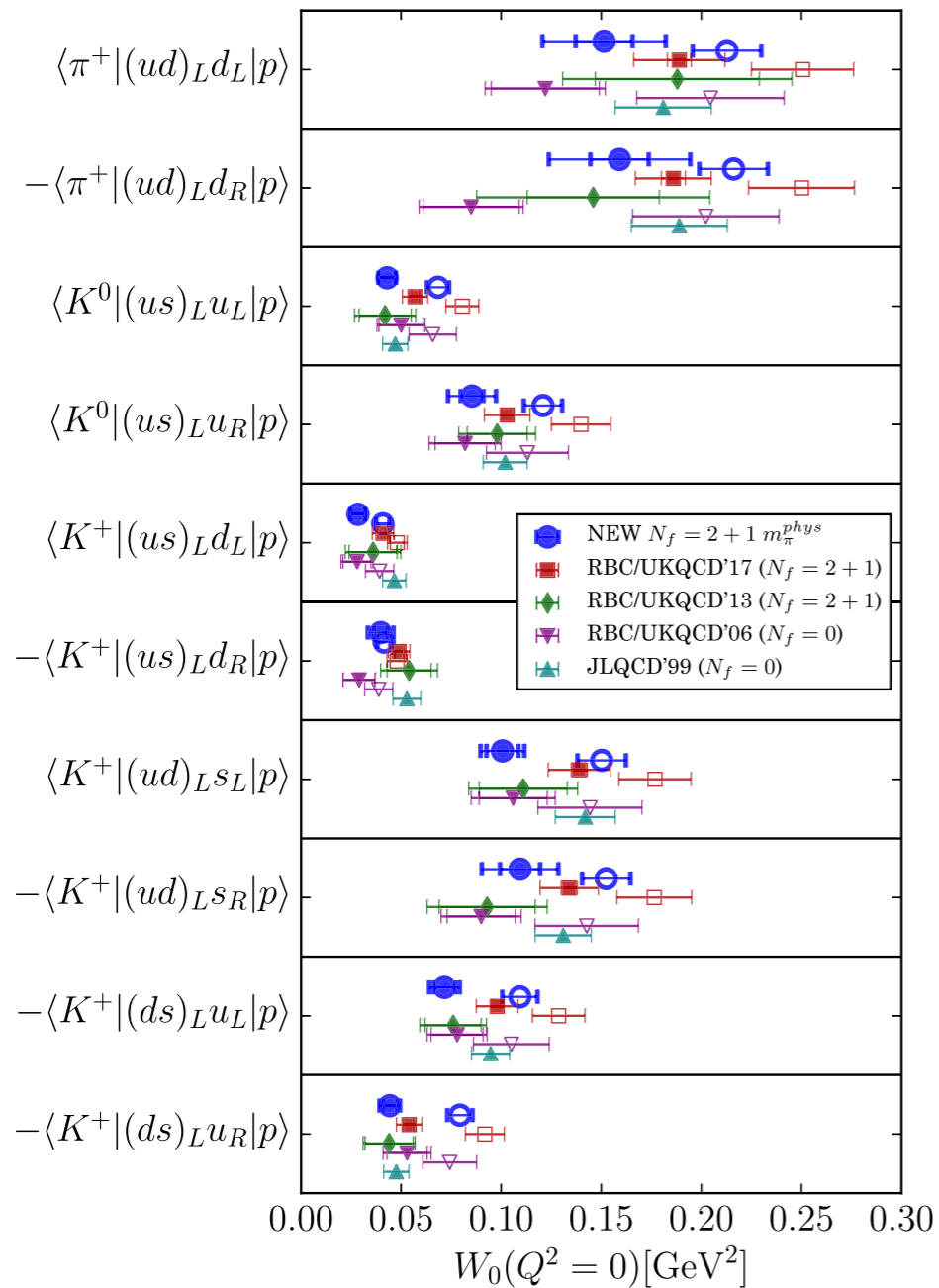


Blue solid symbols are the current best estimate

FIG. 14. Comparison of our results (“NEW”) for the proton decay amplitudes $W_0(0)$ computed directly (filled symbols) and indirectly (open symbols) to previous determinations [38,40,42]. All results are renormalized to the $\overline{\text{MS}}(2 \text{ GeV})$ scheme.

RBC/UKQCD results

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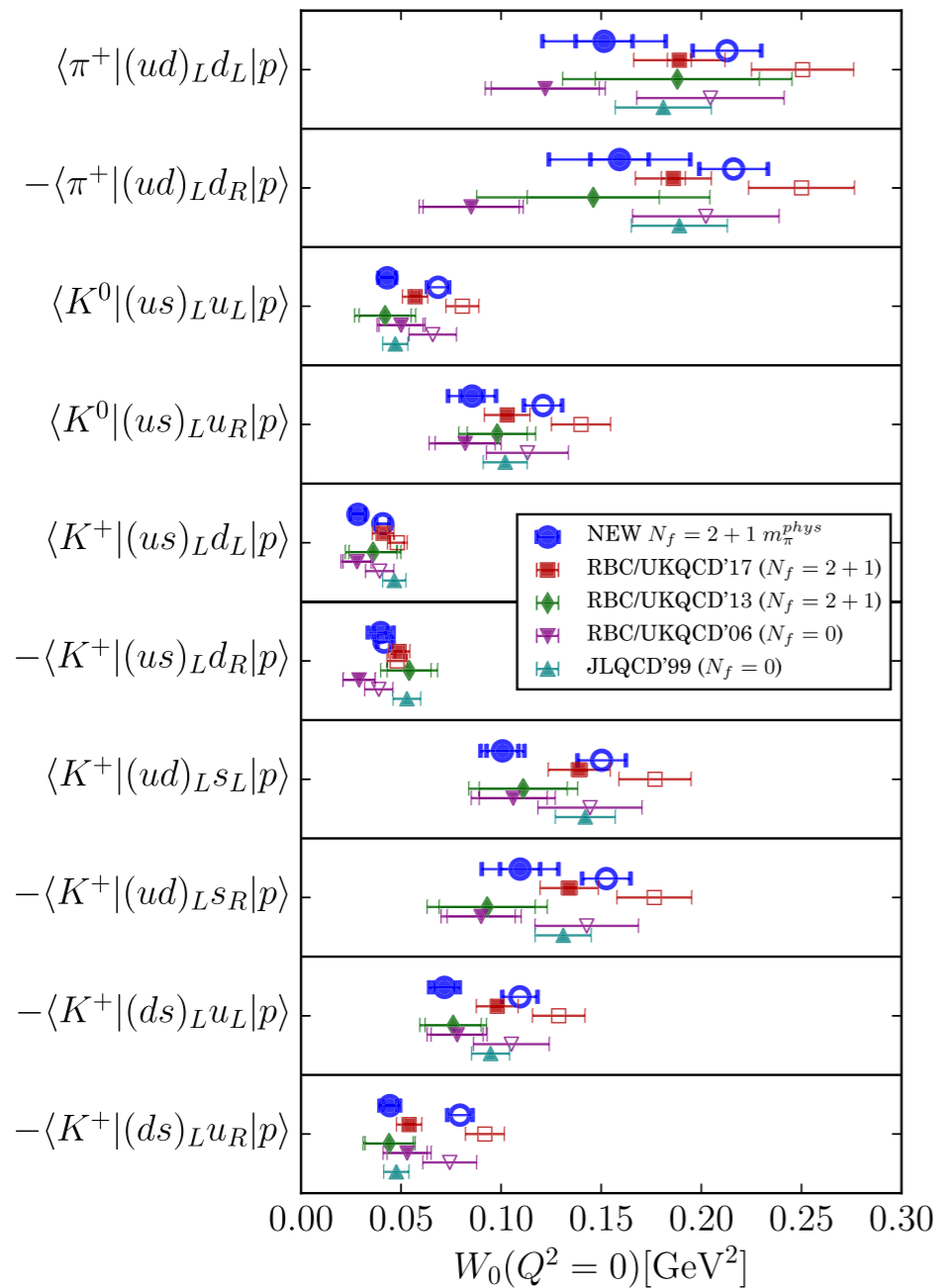
Blue solid symbols are the current best estimate

LEC ?

FIG. 14. Comparison of our results (“NEW”) for the proton decay amplitudes $W_0(0)$ computed directly (filled symbols) and indirectly (open symbols) to previous determinations [38,40,42]. All results are renormalized to the $\overline{\text{MS}}(2 \text{ GeV})$ scheme.

RBC/UKQCD results

Phys. Rev. D 105, 074501 (2022)



Blue solid symbols are the current best estimate

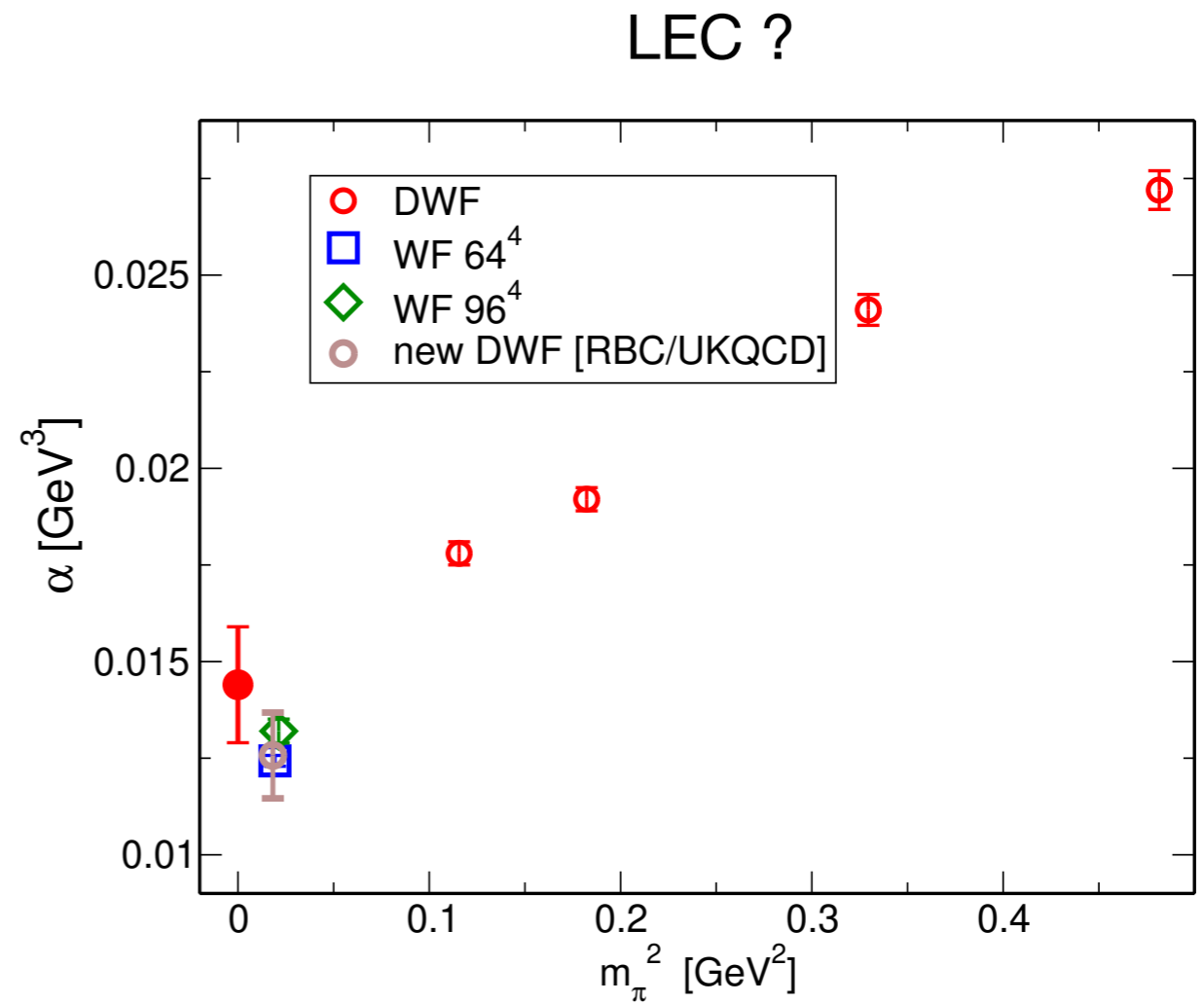


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RBC/UKQCD results

Phys. Rev. D 105, 074501 (2022)

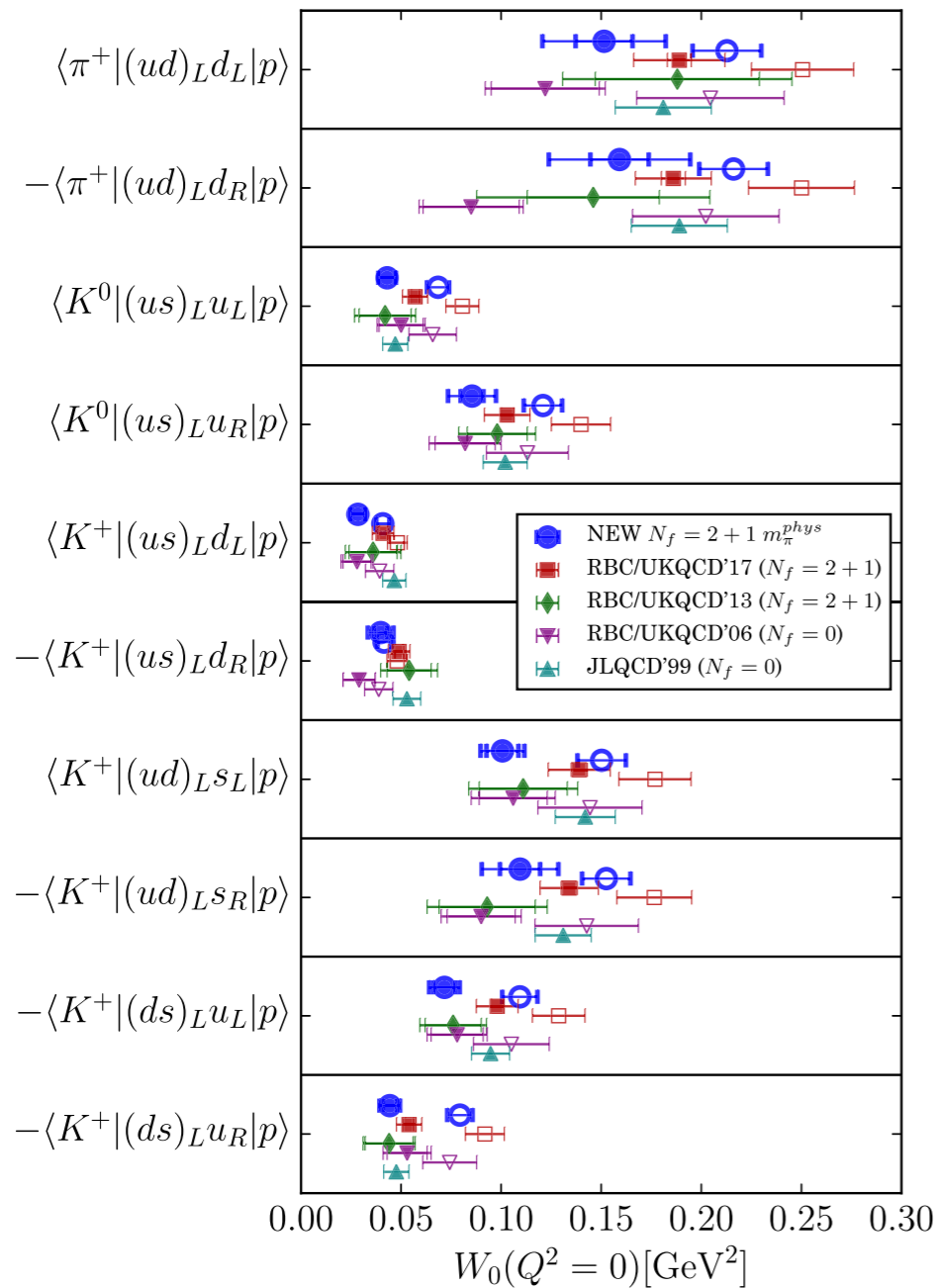
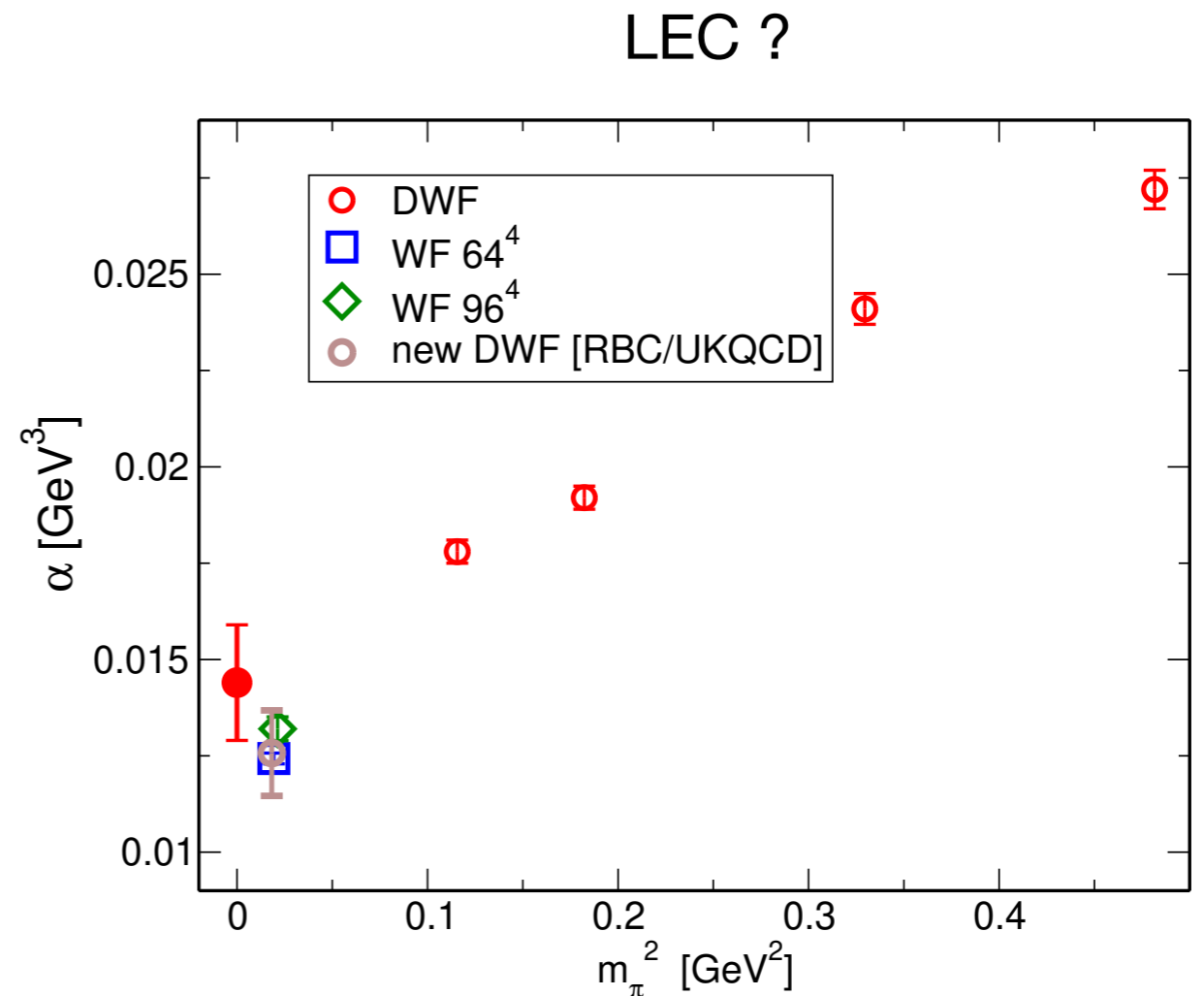


FIG. 14. Comparison of our results (“NEW”) for the proton decay amplitudes $W_0(0)$ computed directly (filled symbols) and indirectly (open symbols) to previous determinations [38,40,42]. All results are renormalized to the $\overline{\text{MS}}(2 \text{ GeV})$ scheme.

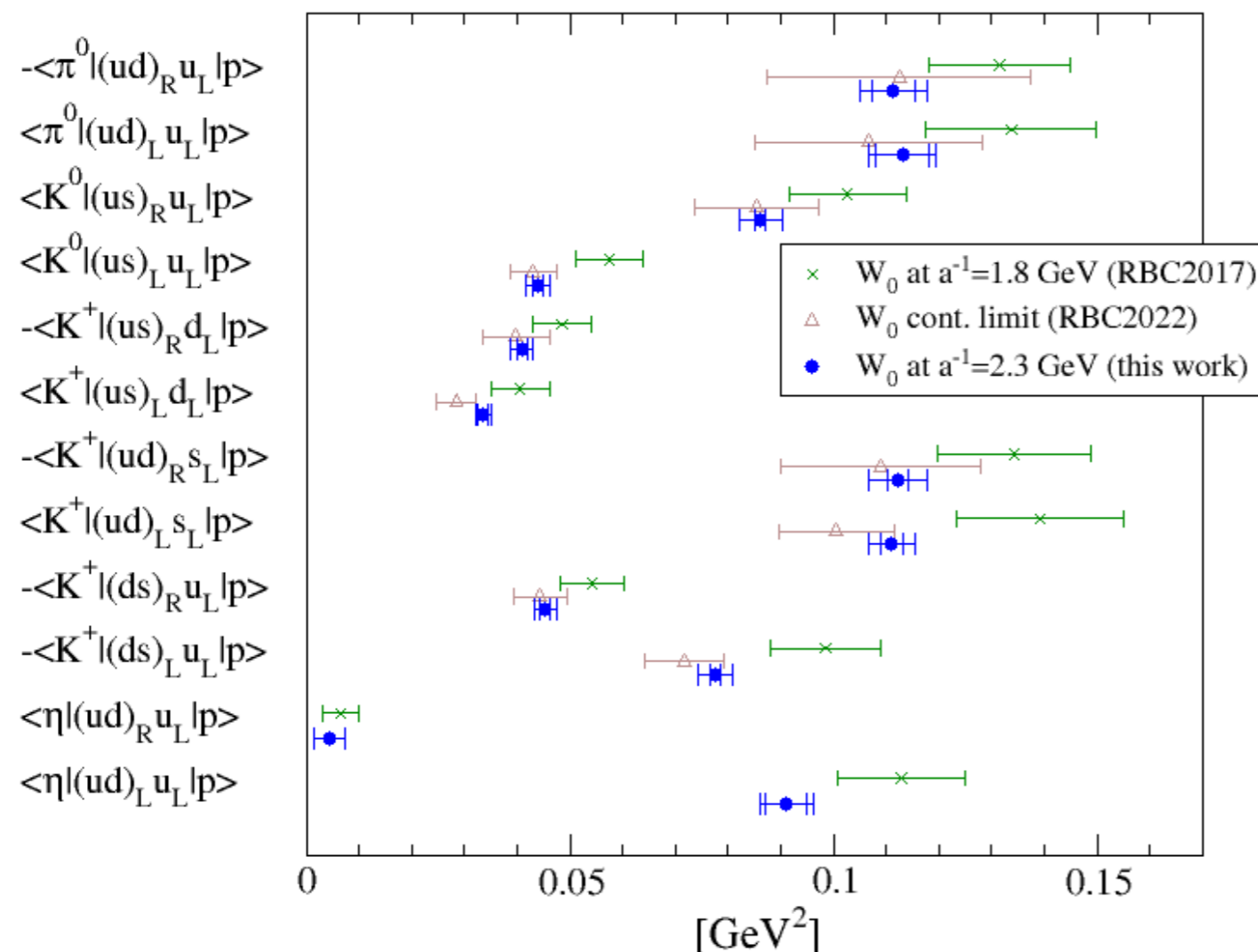
Blue solid symbols are the current best estimate



No surprise happened at physical point !

Preliminary PACS results of proton decay FF $W_0 (\overline{MS}, 2\text{GeV})$

- First Wilson fermion physical point (preliminary) results
- compared with DWF (RBC 2017, RBC 2022(physical point, continuum))



- (L)QCD parity invariance:
 - $W_0^{LR} = W_0^{RL}$
 - $W_0^{RR} = W_0^{LL}$
- PACS yet to do:
 - discretization error est.
 - no large error expected
- Consistent with DWF

Summary and outlook

Summary and outlook

- proton decay FFs till 2017 may suffer from chiral extrapolation error
- Now we can do computations on physical mass (no extrapolation)
 - aiming to remove final loose end
- Using PACS Wilson ensemble
 - RI/SMOM non-perturbative renormalization schemes applied
 - robust against SSB and mass effect
 - W_0 ($p \rightarrow \pi^0$), LEC α and β consistent with DWF(2017) [preliminary]
- RBC/UKQCD with domain wall fermion
 - New renormalization scheme help reduce the systematic error
 - W_0 ($p \rightarrow \pi^0$), LEC α and β consistent with DWF(2017)
 - \rightarrow No chiral limit surprise! $\leftarrow \rightarrow$ Martin & Stavenga (skirm chiral bag)

PACS analysis to be finalized, envisioning comparable /better accuracy
plan to use PACS10c configurations w/ continuum scaling study

acknowledgements

- Thanks for various help from other PACS members esp
 - Takeshi Yamazaki, Shoichi Sasaki, Natsuki Tsukamoto
- Used Computers
 - K-computer,
 - Oakforest PACS,
 - Cygnus,
 - Supercomputer Fugaku,
 - Wisteria/BDEC-01
- Programs
 - Multidisciplinary Cooperative Research Program in CCS Tsukuba
 - HPCI: hp200062, hp200167, hp210112, hp220079, hp230199

Program for Promoting Researches
on the Supercomputer Fugaku

Large-scale lattice QCD simulation
and development of AI technology

JPMXP1020230409

Thank you very much for your attention

appendix

use of RI/SMOM schemes [for PACS analysis]

- original NPR for proton decay
 - wave function renormalization from (axial) vector currents with exceptional momentum configuration
 - which is sensitive to low energy parameters like mass
 - RI/MOM defined $m \rightarrow 0$; we have only physical point $0 \simeq m_{ud} \ll m_s$
- new attempt using SMOM schemes for wave function
 - 3q vertex has not been using exceptional momenta
 - now no exceptional momenta at all
 - safer to use with physical point ensemble
- matching from SMOM schemes to MSb to NLO available
 - (wave function renormalization known to NNLO, 3q unchanged)

RI/SMOM procedure for proton decay operators

- calculate 3q vertex: same momentum p injected from all quark legs
- cancel the wave function renormalization by dividing with local vector or axialvector vertex in SMOM schemes with proper power
- this will provide: $Z_{ND}/Z_A^{3/2}$, $Z_{ND}/Z_V^{3/2}$
- take Z_A , Z_V factor away by inputting those calculated in SF scheme
- Z_{ND} (SMOM) matched to MSb and ran to 2 GeV by NLO PT
- final value and error estimate:
 - error from truncation of PT, $(pa)^2$, contamination of SSB
 - center = $(\max+\min)/2$, error = $(\max-\min)/2$
 - max, min determined in
 - $(pa)^2$ linear extrapolation: 4 vales: (A,V) x (SMOM, SMOM $\gamma\mu$)
 - scatter of all data in the window 2-3 GeV

Lattice setup: input and output

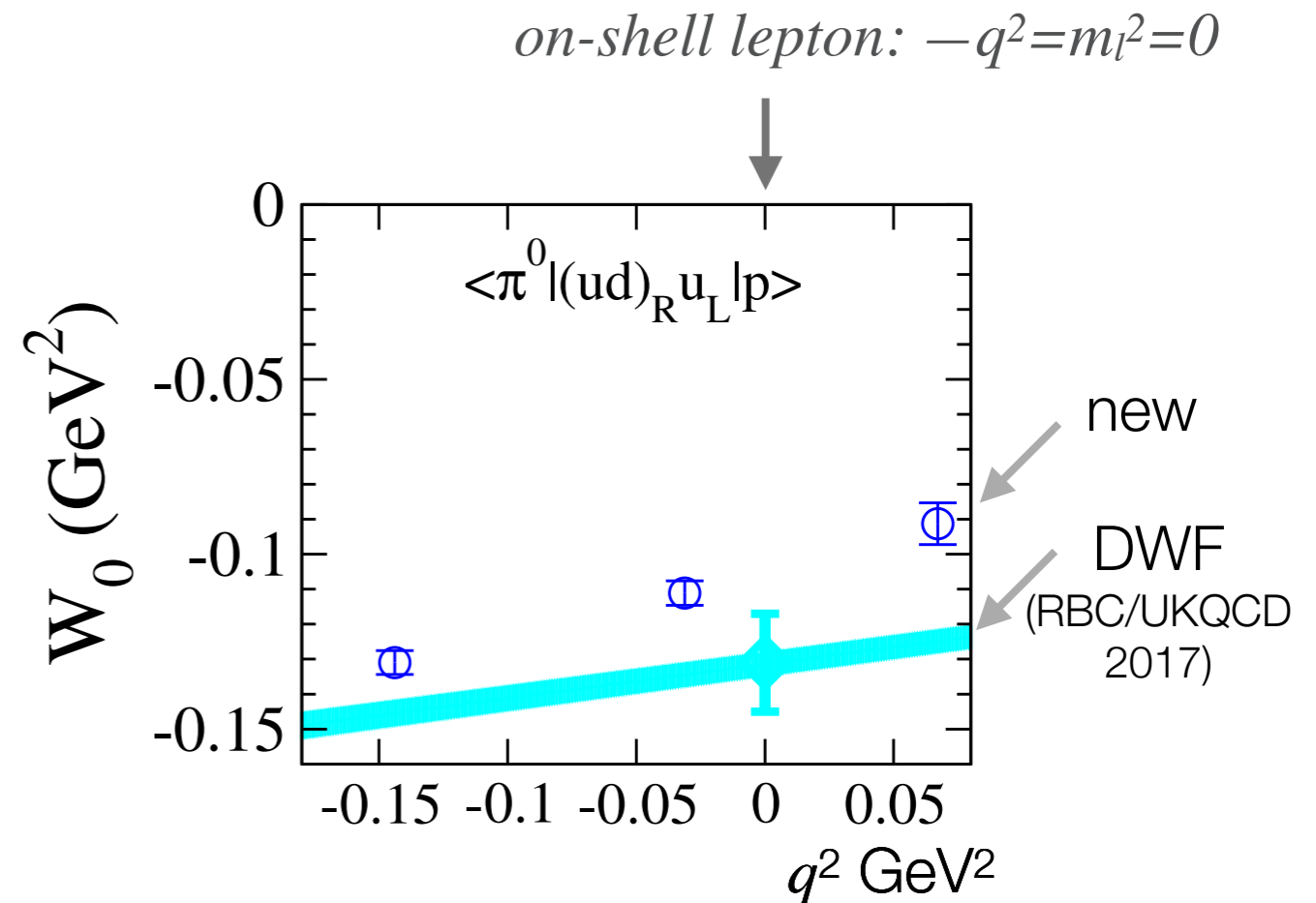
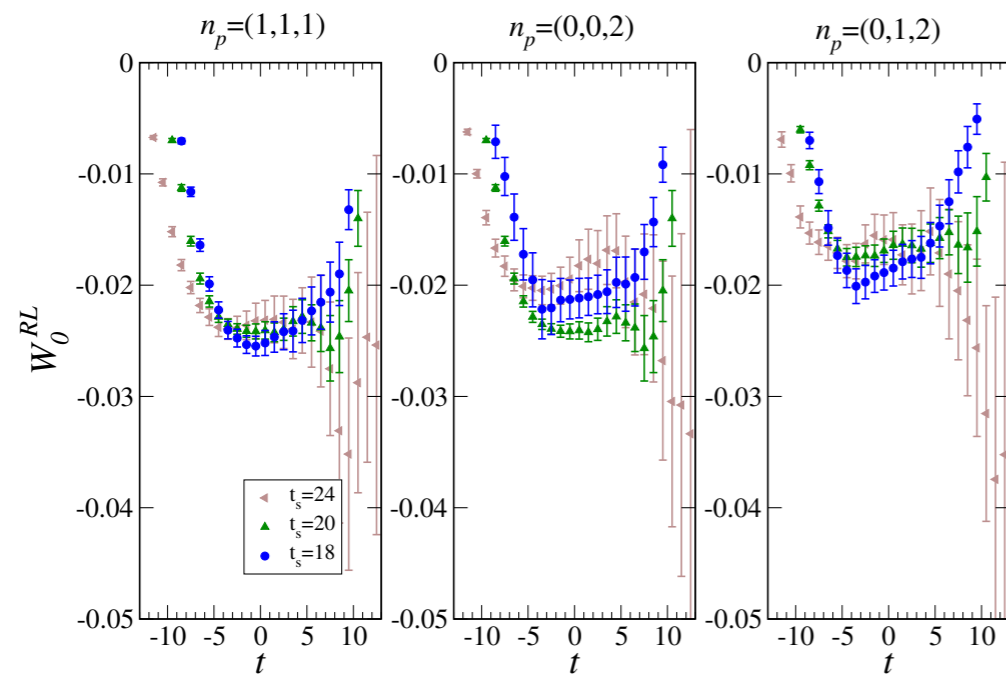
- * $2+1$ flavor lattice QCD computation parameters:
 - * gauge coupling, $m_{ud}(\text{degenerate})$, m_s
 - * Monte Carlo simulations are done at $m_s^{(sim)} \approx m_s^{(phys)}$, $m_{ud}^{(sim)} > m_s/5$
 - * small mass is demanding: cost $\sim 1/m^x$: $x > 1$
 - * at a fixed gauge coupling, tune m 's so that it reproduces ratios of π , K , Ω mass.
 - * π , K : quark mass dependence best known: NNLO ChPT / analytic also tested
 - * Ω (sss): no pion chiral logs at NLO: safe to apply linear chiral extrapolation
- * all other quantities are predictions, ex: [RBC/UKQCD PRD78(08)114509]
 - * $f_{\pi} = 124.1(6.9) \text{ MeV} \leftrightarrow 130.7(0.1)(0.36)[\text{exp}]$, $f_K/f_{\pi} = 1.205(18) \leftrightarrow 1.223(12)[\text{exp}]$
 - * one lattice spacing, estimate of $O(a^2)$ systematic error was added.
 - * quark masses, B_K ...
- ➔ continuum limit results available: [RBC/UKQCD PRD83(11)074508, PRD84(11)014503]

proton decay form factor W_0 for pion final state

Relevant form factor $W_0 \langle \pi^0 | (ud)_{RUL} | p \rangle$ as an example

- from ratio of 3 and 2 point functions

- meson momentum $\vec{p} = \frac{2\pi}{L} \vec{n}_p$



- $|W_0|$ ~20% smaller than DWF (with a long chiral extrapolation) at $q^2=0$
- consistent with sys. error ! no big surprise found for $m_f \rightarrow m_{ud}$
- 10% total error is not a dream...