

# Modification of Hard Probes of the QGP due to Flow-Induced Jet Drift

arXiv preprint coming soon...

*Supported in part by DoE Grant (DE-SC0024560)  
Supported in part by a start-up grant from NMSU*

INT  
6<sup>th</sup> August, 2024

Jo Bahder

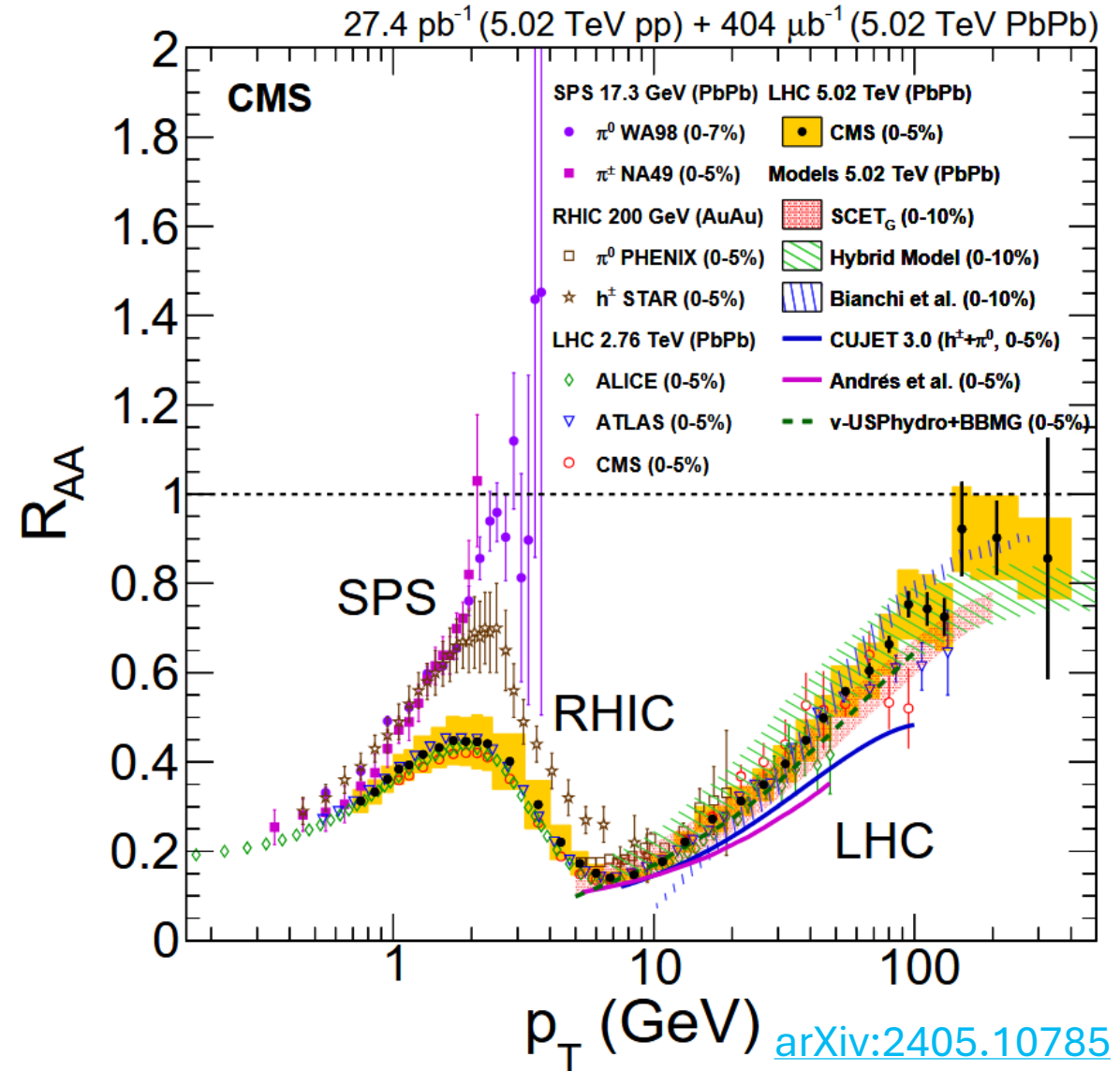


In collaboration with  
Hasan Rahman, Matthew Sievert,  
and Ivan Vitev

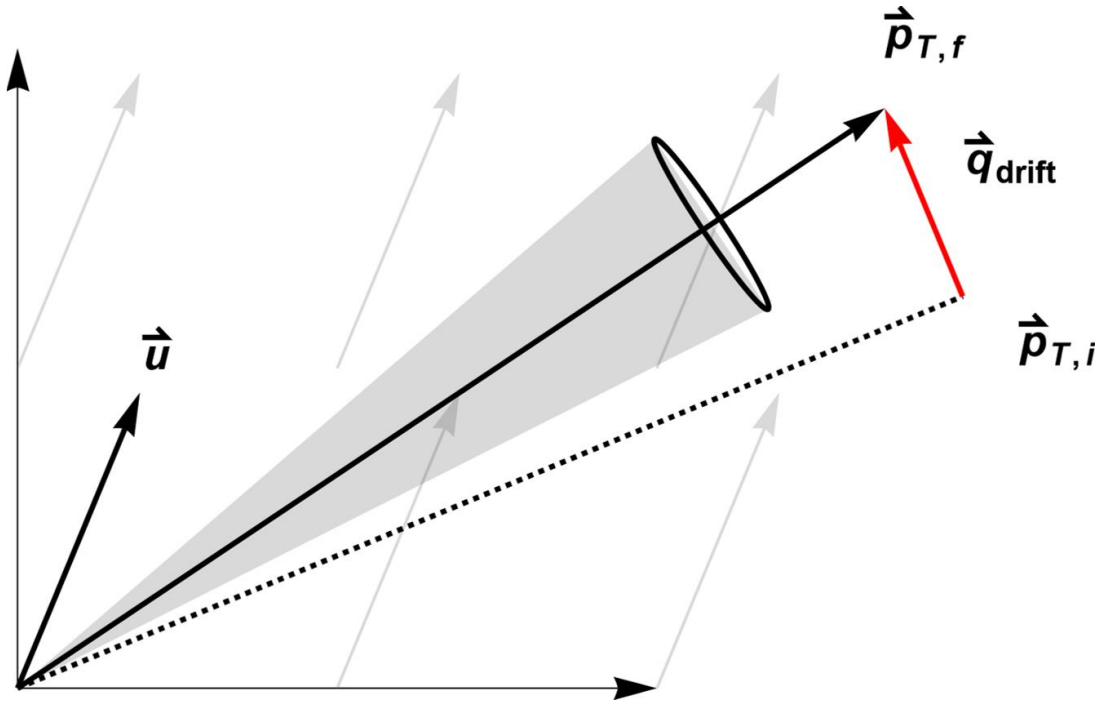


# The Specialness of Anisotropic Effects

- Isotropic modifications of hard probes are not strongly model-discriminate
  - Compete with many sources of background
  - Many microscopic models obtain results similar to data (e.g. ch. had.  $R_{AA}$ )
- Few, if any, effects can mimic the coupling of hard particles to anisotropies in the medium
  - Few vector directions: Collective Flow & Gradients



# Anisotropic Jet Broadening: "Jet Drift"



- Part of a "New" class of pQCD effects:  
asymmetric / anisotropic
- Preferential broadening in direction of medium flow

Flow enhanced  
and flow direction  
controlled

A. V. Sadofyev, I. Vitev,  
& M. D. Sievert  
Phys.Rev.D 104 (2021)  
[arXiv:2104.09513](https://arxiv.org/abs/2104.09513)

$$\langle \vec{q}_{drift} \rangle = \hat{e}_\perp \int d\tau \begin{matrix} 3 \\ E(\tau) \end{matrix} \begin{matrix} \mu^2(\tau) \\ \lambda(\tau) \end{matrix} \ln \frac{E(\tau)}{\mu(\tau)} \begin{matrix} u_\perp(\tau) \\ 1 - u_\parallel(\tau) \end{matrix}$$

$$\frac{1}{\lambda} = \sigma \rho$$

$$\rho \propto T^3$$

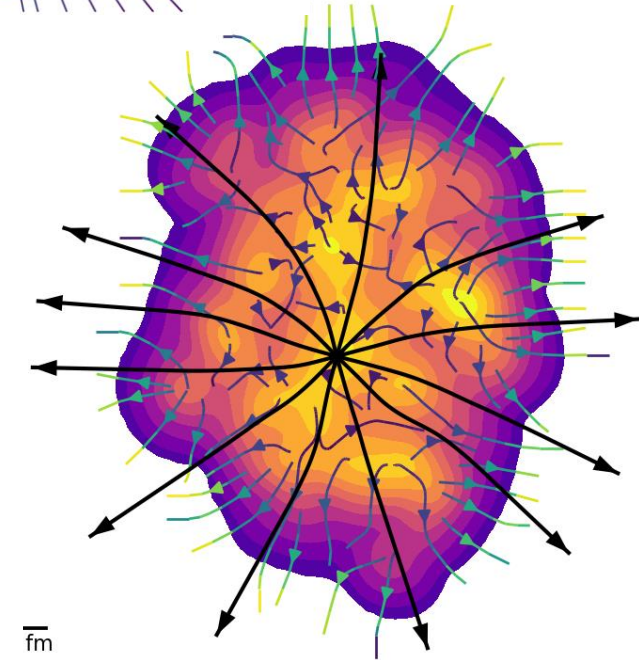
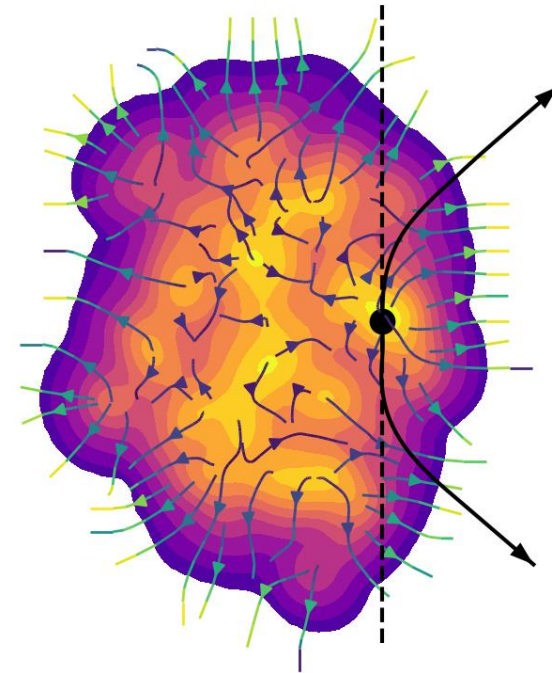
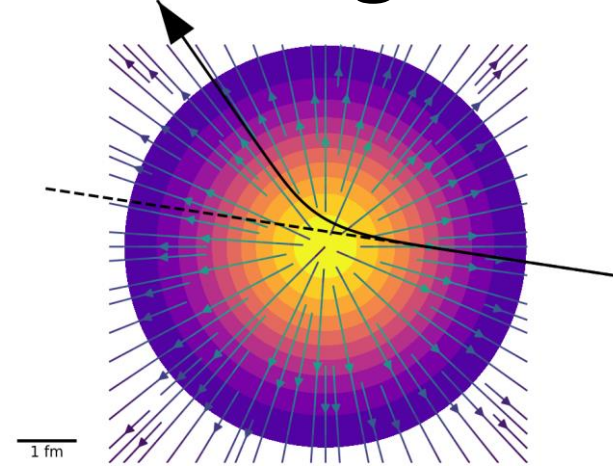
$$\mu \propto T$$

Energy suppressed

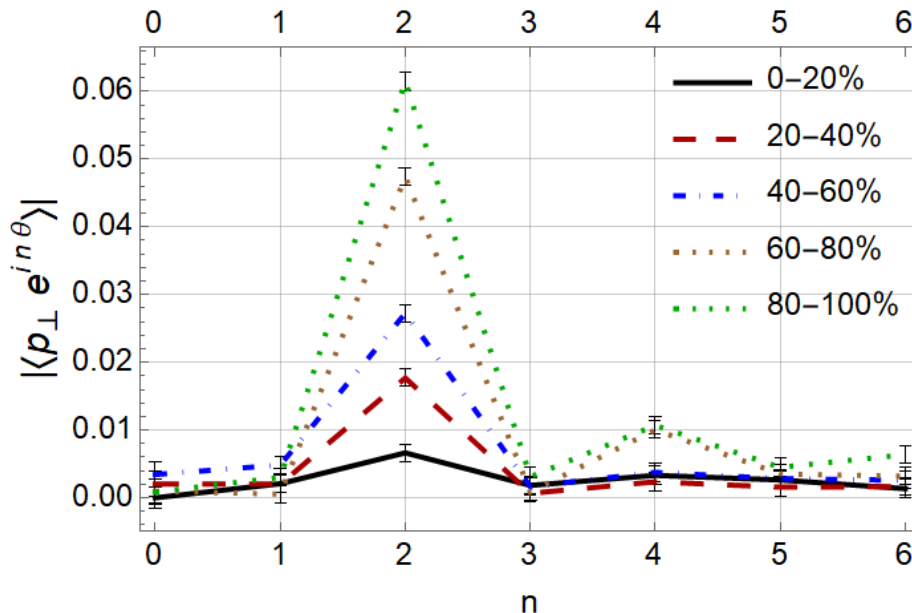
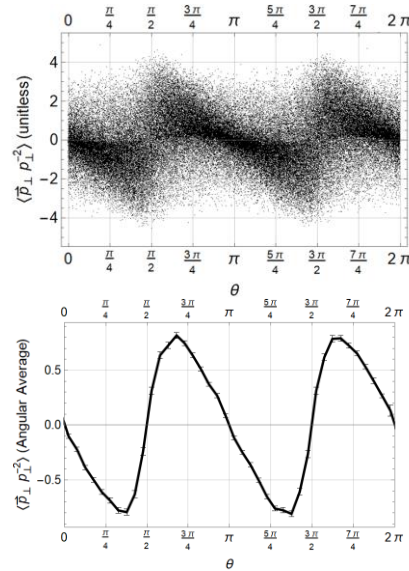
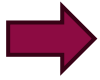
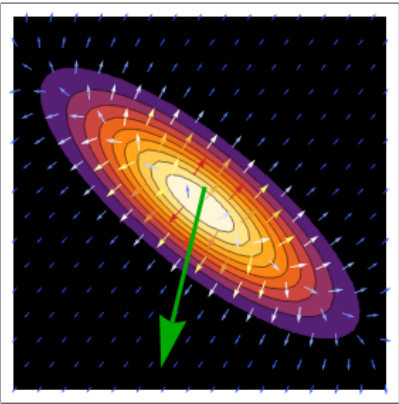
Temperature/Density Enhanced

# Possible Signatures of Jet Drift in Leading Hadrons

- Modification of suppression
  - Deflecting away from flowing hot spots
  - Particle sees reduced integrated density
- Anisotropic flow modification
  - Particles couple to soft anisotropic flow
- Acoplanarity enhancements
  - Particles couple to same “attractor” in the medium



# Does Drift Survive Event Averaging?



- Naively, one might expect cancellation via event averaging
- Coupling to event anisotropic flow shown in glauber elliptic geometry to preserve effect
  - Fluctuating event plane, centrality
  - L. Antiporda, J. Bahder, H. Rahman & M. D. Sievert Phys.Rev.D 105 (2022) ([arXiv: 2110.03590](https://arxiv.org/abs/2110.03590))
- What about in realistic heavy ion collisions?

# APE: A Monte Carlo to Study Drift

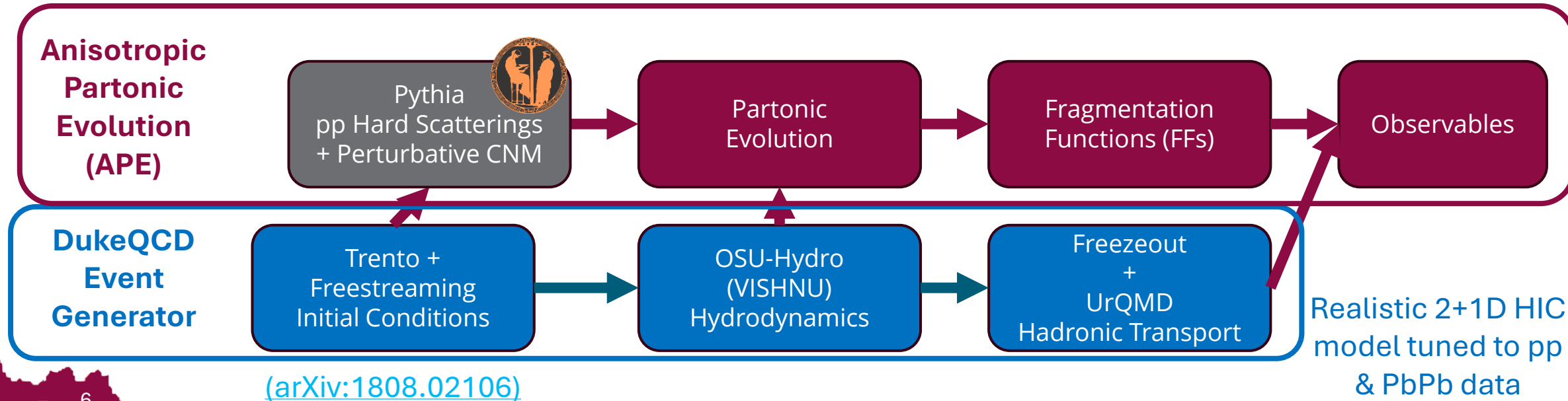
Pert. Cronin  
Effect, Nuclear  
Shadowing, &  
Iso-spin Effects  
([arXiv:0201078](https://arxiv.org/abs/0201078))  
([arXiv:0309094](https://arxiv.org/abs/0309094))

Test effect of addition of jet drift to realistic event-by-event  
jet-medium simulations on hard probes of heavy ion collisions

Highly Tuned pp  
Hard Scatterings  
([arXiv:2203.11601](https://arxiv.org/abs/2203.11601))

Perturbative  
Partonic  
Evolution

Perturbative  
Fragmentation





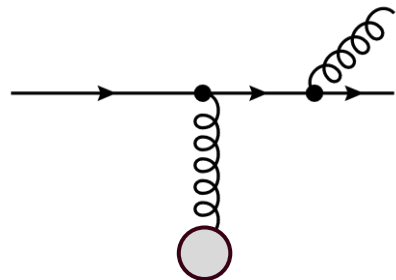
# Energy Loss Effects

## Radiative

$$\frac{dE}{d\ell} = -\frac{d}{dL} \left( \frac{2C_R\alpha_s}{\pi} \frac{L}{\lambda} E \int_{k_{min}}^{k_{max}} \frac{dk}{k} \int_0^{q_{max}} dq q \int_0^{2\pi} d\phi \right. \\ \left. \times \frac{\mu^2}{\pi(q^2 + \mu^2)^2} \frac{2\mathbf{k} \cdot \mathbf{q} (\mathbf{k} - \mathbf{q})^2 L^2}{16x^2 E^2 + (\mathbf{k} - \mathbf{q})^4 L^2} \right)_{L=\ell}$$

- Single Emission GLV @ 1<sup>st</sup> order in opacity w/ finite kinematic bounds (q, k)
- Interpolated tabulated results
- Gyulassy, Levai, Vitev (2000) ([arXiv:0006010](https://arxiv.org/abs/0006010))

E.g.

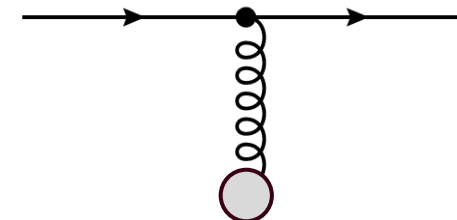


## Collisional

$$\frac{dE}{d\ell} = -C_R \frac{1}{2} \mu^2 \ln \left( 2^{\frac{N_f}{2(6+N_f)}} 0.920 \frac{\sqrt{3ET}}{\mu} \right)$$

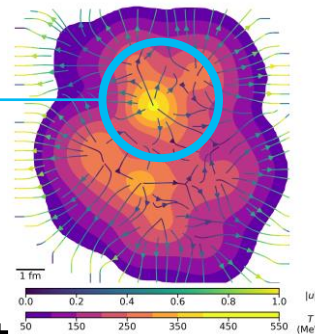
- Braaten & Thoma (1991) ([INSPIRE:317898](https://arxiv.org/abs/hep-ph/9103017))
- Light quarks:  $E \gg m^2/T$  regime
- Gluons:  $CA/CF = 9/4$  ([arXiv:2305.13182](https://arxiv.org/abs/2305.13182))

E.g.



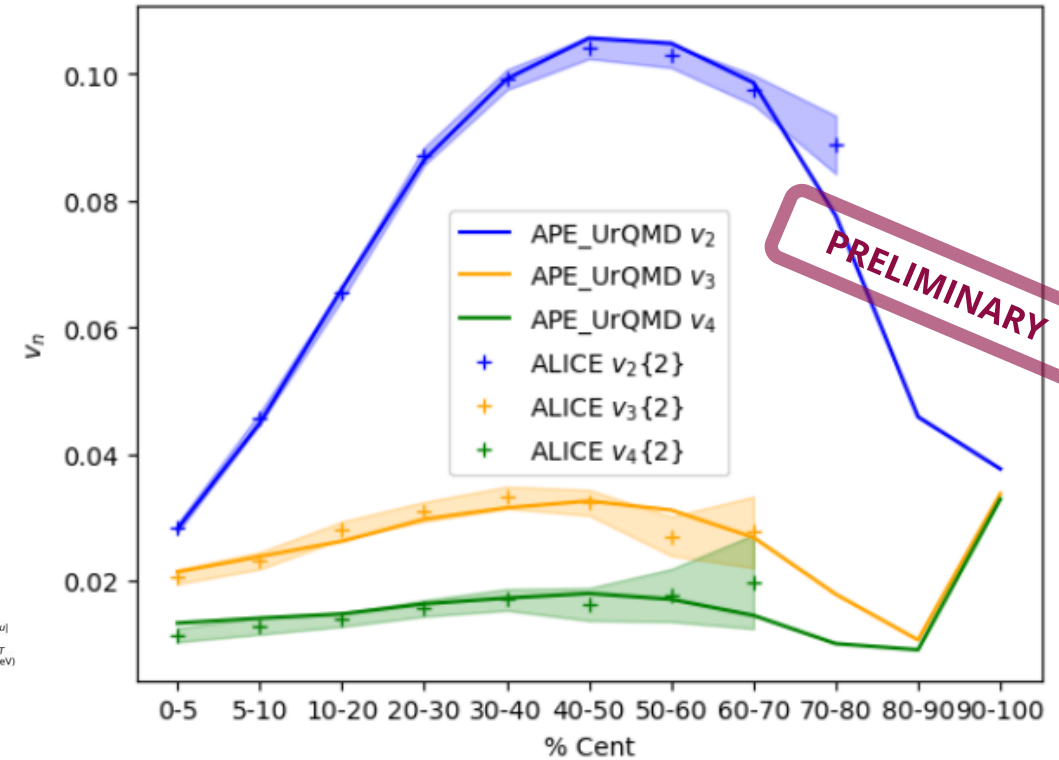
# DukeQCD's HIC Event Generator

- Highly tuned to soft sector observables
  - Reliable picture of final state
  - Not necessarily good description of intermediate dynamics – drift distinguishes
- Large event-by-event fluctuations
  - Nucleon substructure fluctuations included
  - Maximizes disruption of event-correlated drift



$$\langle \vec{q}_{drift} \rangle = \hat{e}_\perp \int d\tau \frac{3}{E(\tau)} \frac{\mu^2(\tau)}{\lambda(\tau)} \ln \frac{E(\tau)}{\mu(\tau)} \frac{u_\perp(\tau)}{1 - u_\parallel(\tau)}$$

Energy suppressed

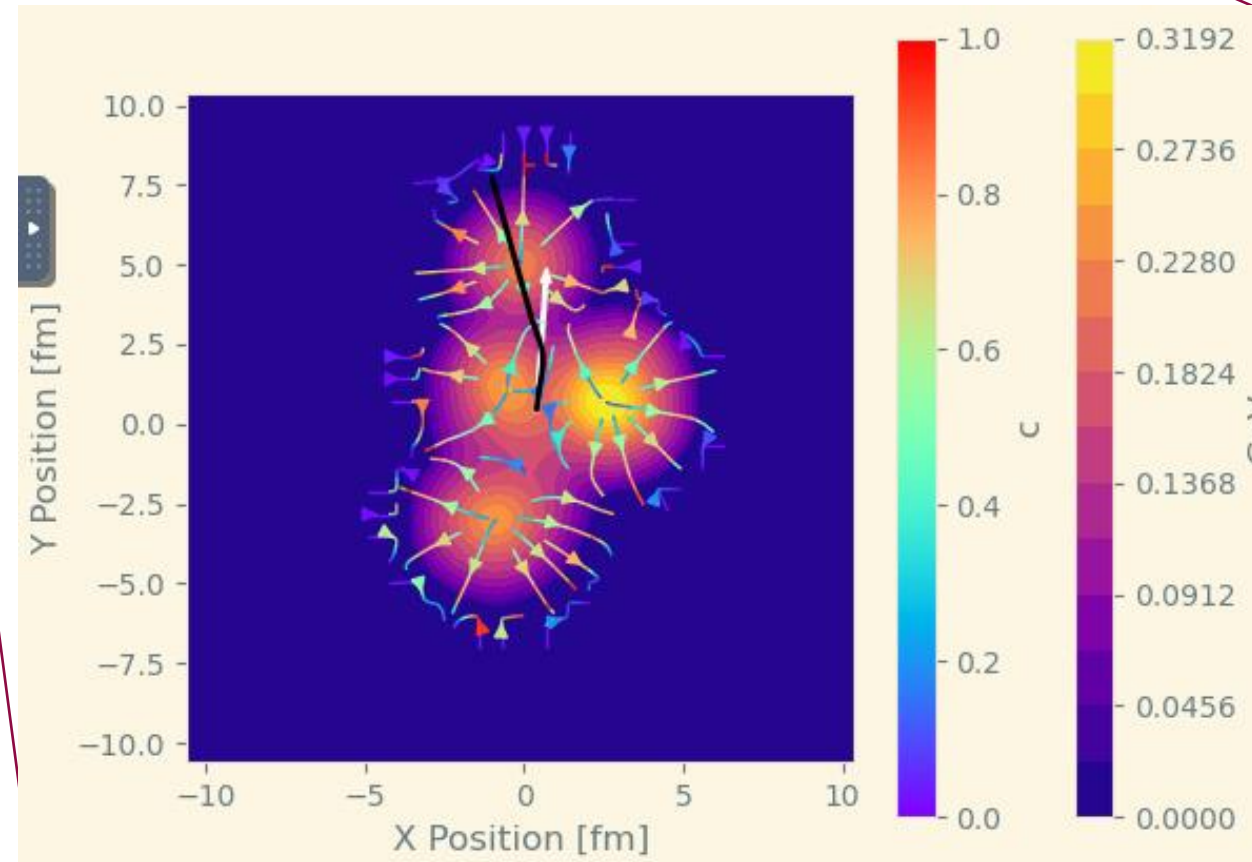
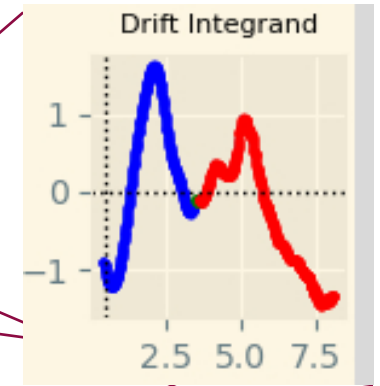
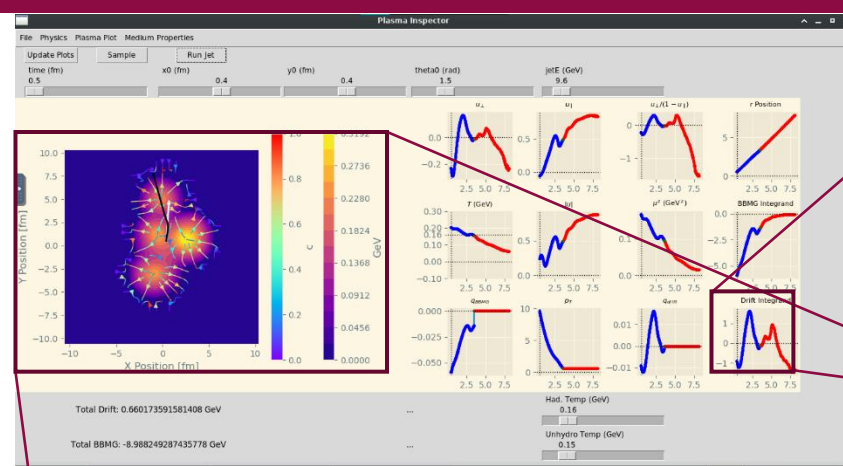


See also multiplicities, mean pT, pT fluctuation, etc.: ([arXiv:1808.02106](https://arxiv.org/abs/1808.02106))

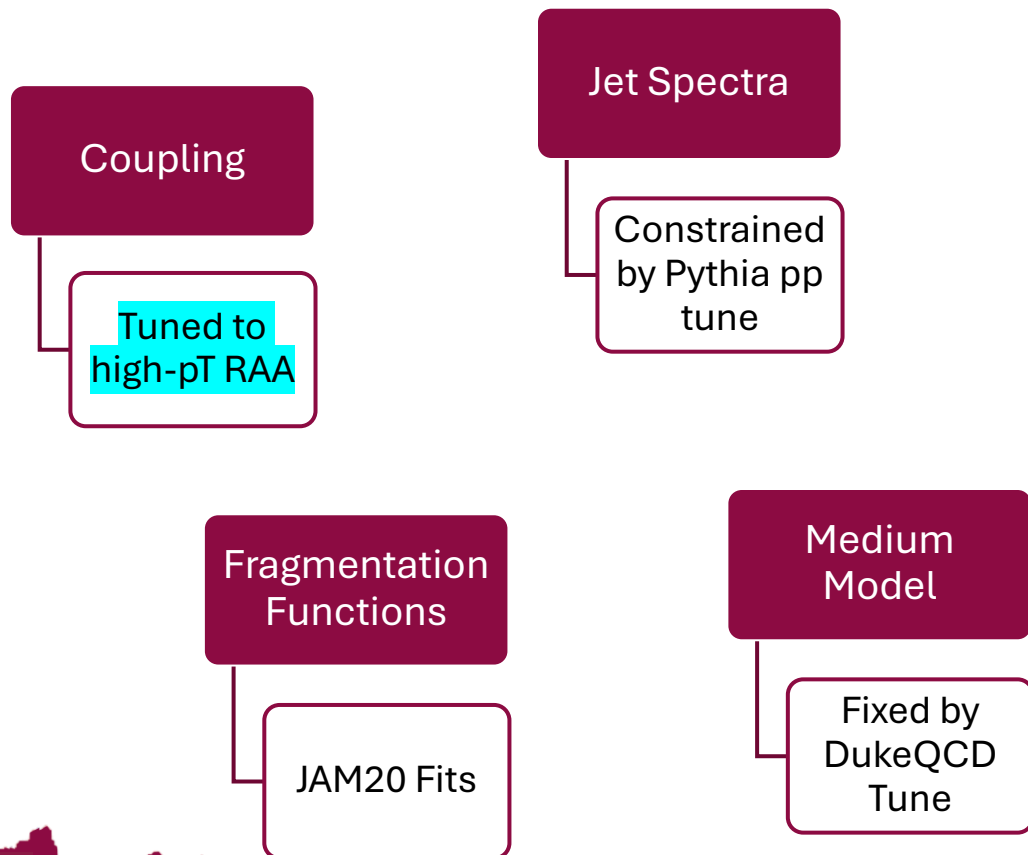


# Ape Trajectories

- Approx. Binary collision density weighting of production points
- Computed within QGP phase of hydro backgrounds
  - EL & Drift cut off at  $T < 155$  MeV
  - Cuts off highest flow region!!!
- Dynamic trajectories respond to medium flow
  - Deflections, zigzags, weirdness



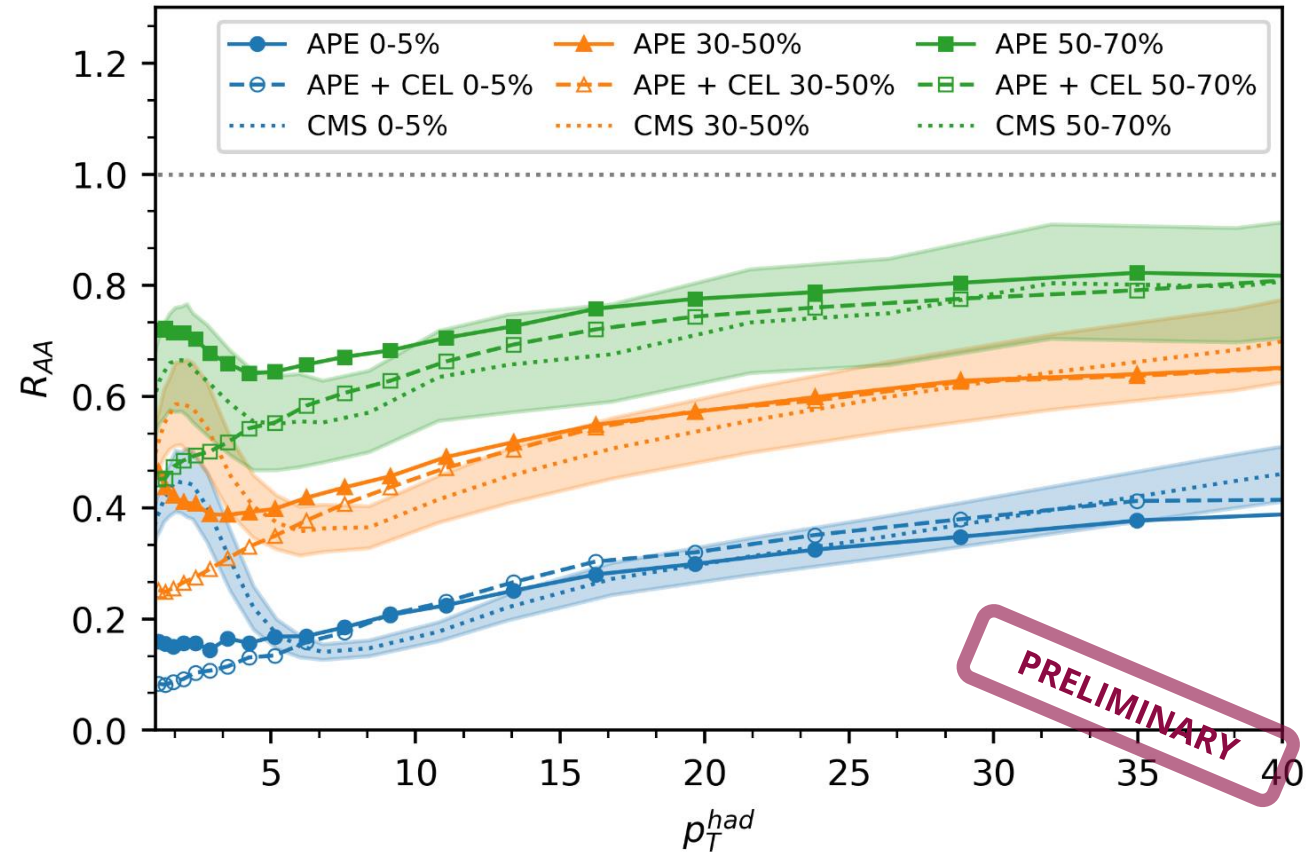
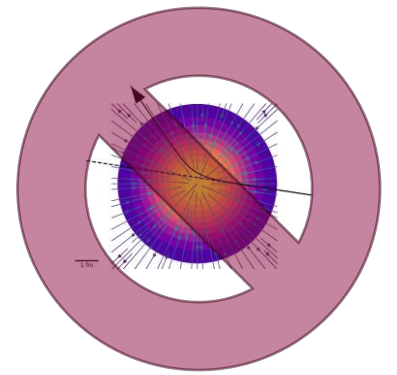
# Parameter Fits & Model Choices



- DukeQCD “hic-eventgen” medium model parameters set by Bayesian parameter estimation  
[\(\[arXiv:1804.06469\]\(https://arxiv.org/abs/1804.06469\)\)](https://arxiv.org/abs/1804.06469)
- Pythia input + pCNM determines partonic spectra
- Coupling from high pT RAA (30-50%)
- Choice of fragmentation function fits
  - Large change to scale of results!

# R\_AA Tunes Coupling

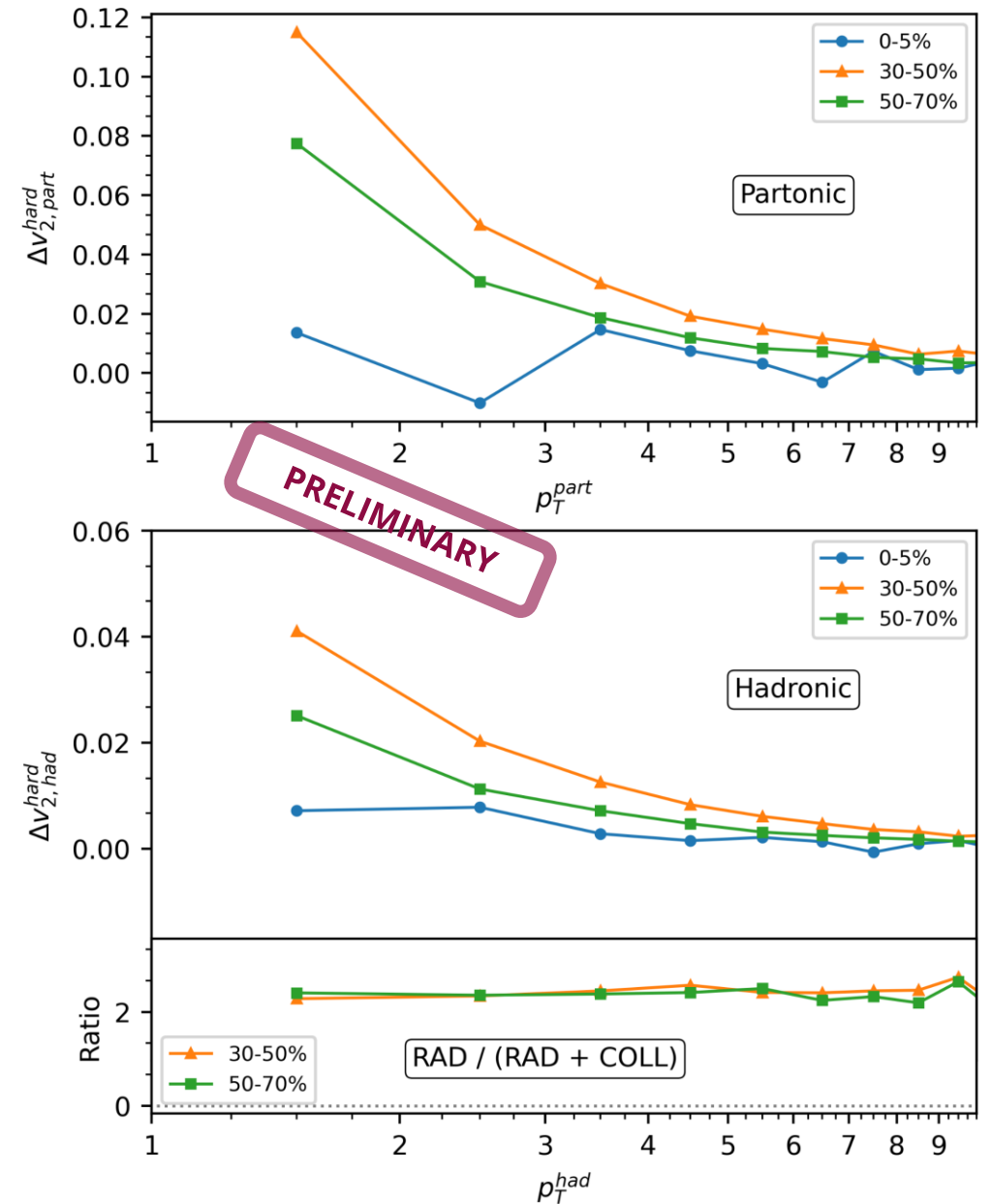
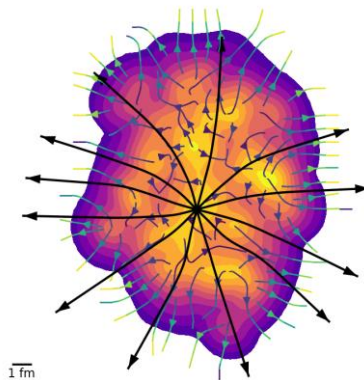
- Drift does not measurably modify  $R_{AA}(p_T)$ 
  - Very useful – can tune coupling
  - Rad:  $g=2.0$ , Rad + Coll:  $g=1.6$
- Additional sources of high- $p_T$  energy loss reduce effective drift coupling
- Better performance at low- $p_T$  can only enhance drift



PRELIMINARY

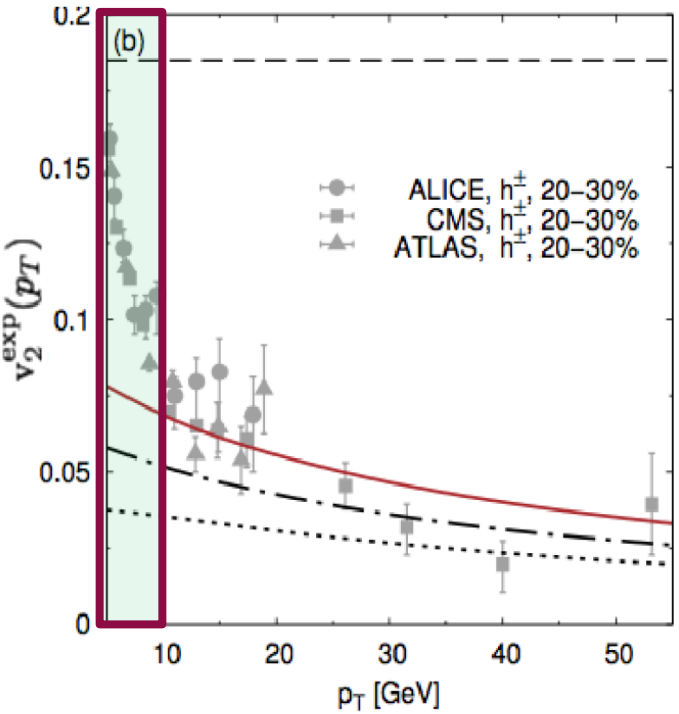
# V2 is Enhanced by Drift

- Large surviving  $v_2$  modulation at low- $p_T$ 
  - Compare to:  $\pm 1\%$  exp. uncertainty
- Conservative estimate of drift
  - Low temp cutoff removes large drift region
  - CNM effects + Coll. Energy loss reduce relative strength
  - No coalescence-type effects

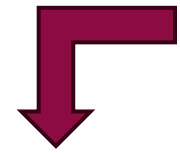


# Importance for the $R_{AA} \times v_2$ Puzzle

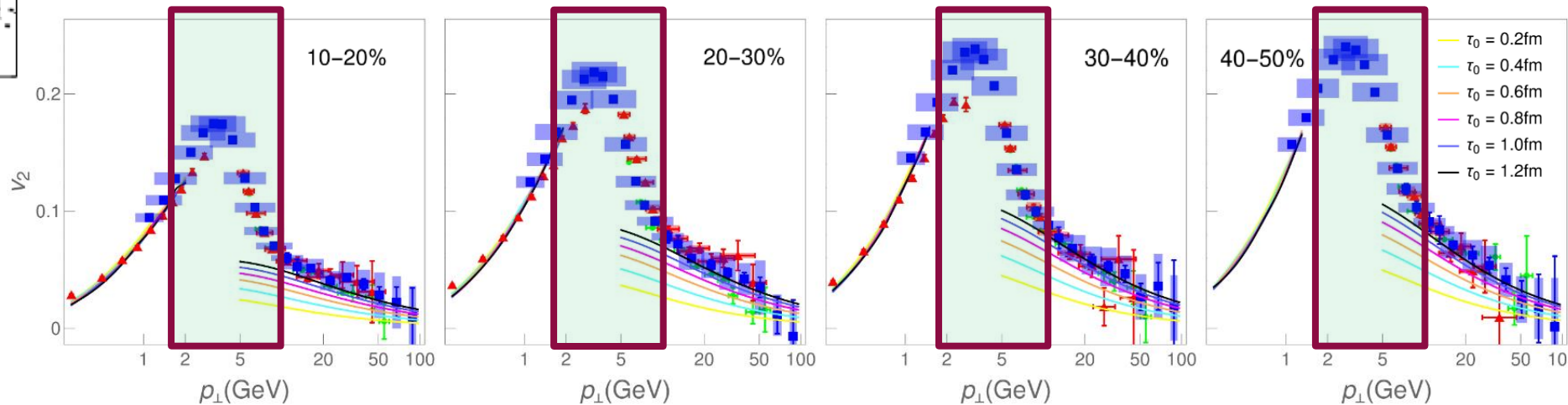
Drift produces measurable elliptic modulation at low  $p_T$  that qualitatively matches elliptic flow missing in perturbative & similar calculations!



Parametric Energy Loss  
Noronha-Hostler et al:  
[arXiv: 1602.03788](https://arxiv.org/abs/1602.03788)

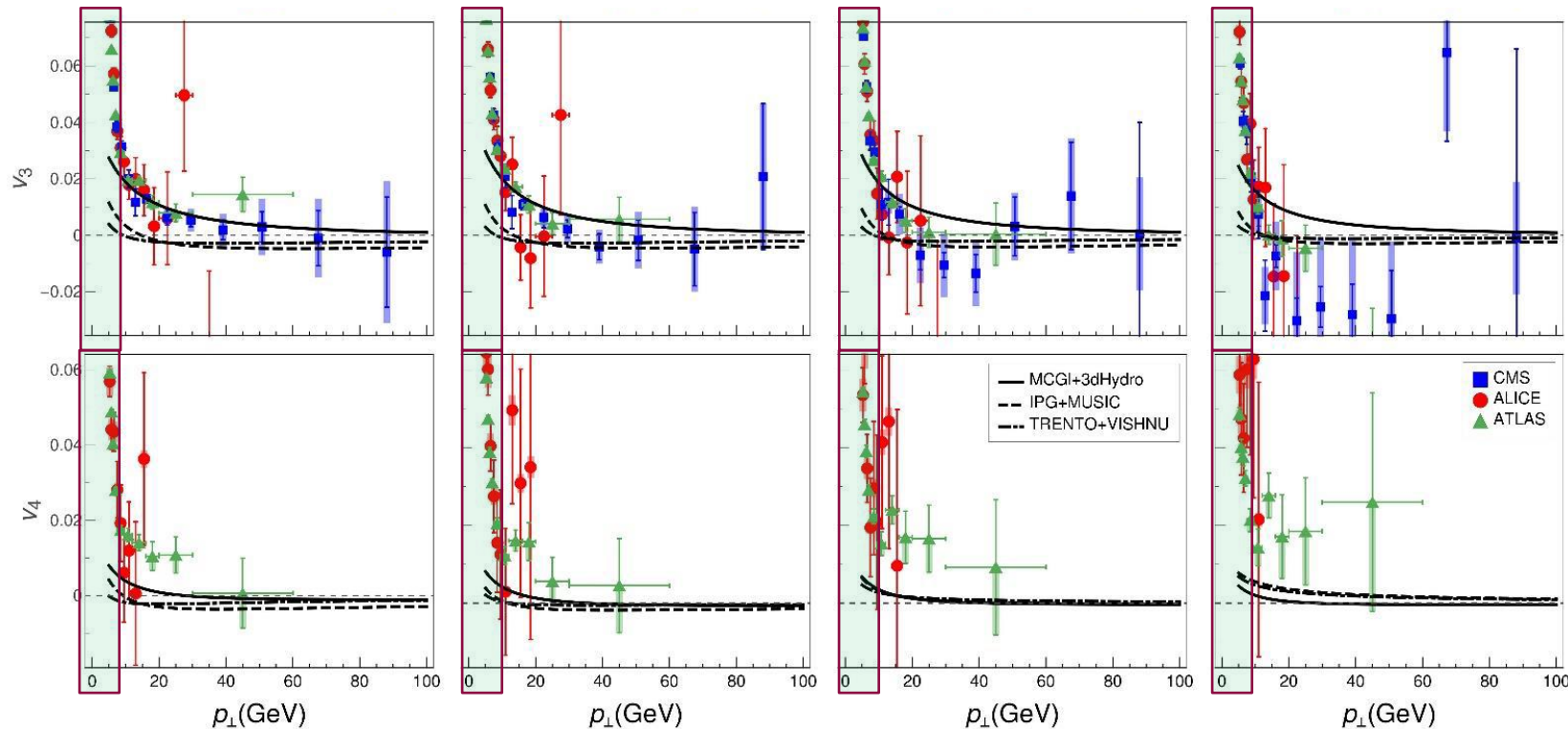


DREENA-A Ch. Had.  
LHC PbPb  
 $\sqrt{s} = 5.02$  TeV  
Perturbative Energy Loss  
Stojku et al.  
[iNSPIRE:2640923](https://arxiv.org/abs/1602.03788)





# Modulation not Limited to Elliptic Harmonics



([arXiv:2208.09886](https://arxiv.org/abs/2208.09886))

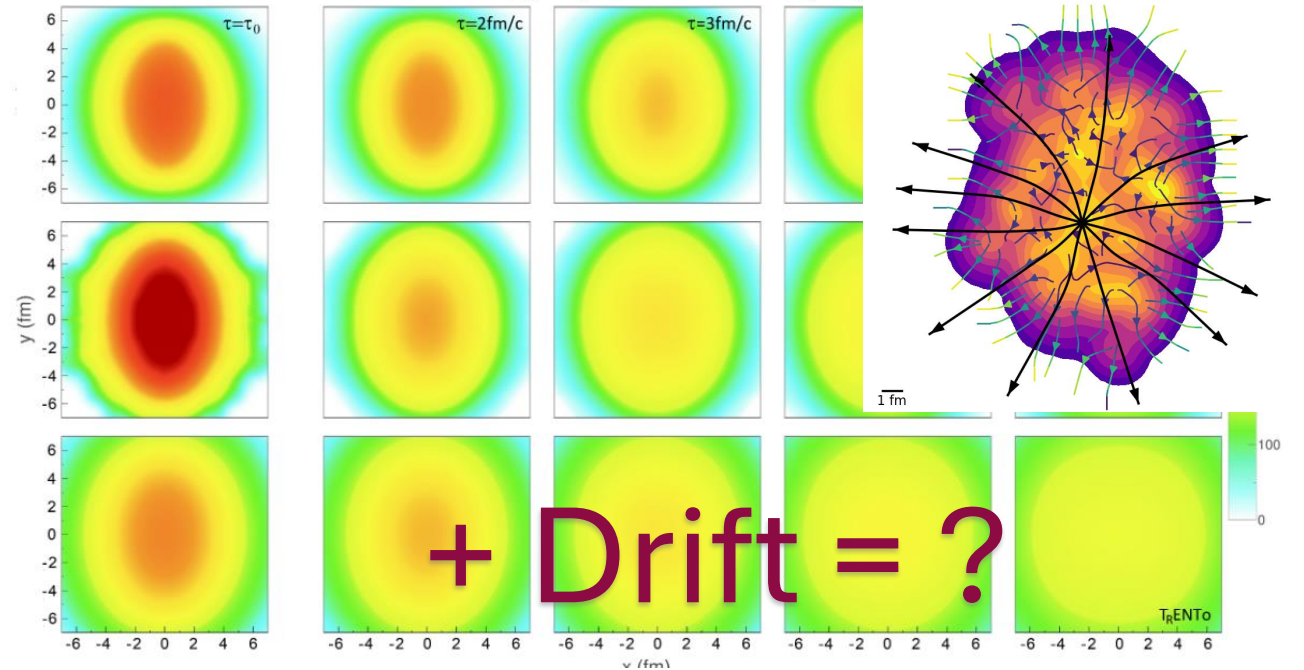
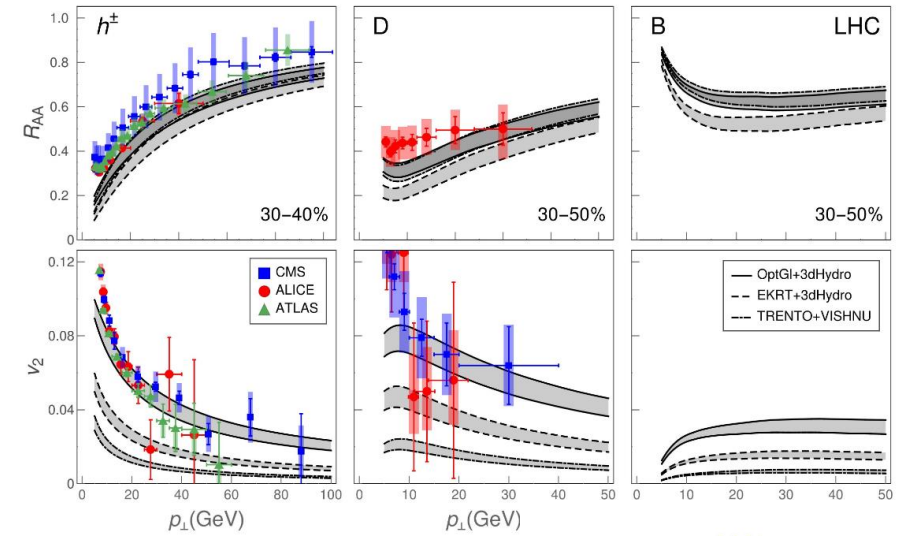
- Higher harmonic coupling works, too
- Possible importance to  $v_3$  &  $v_4$
- Difficult to couple to small anisotropies with energy loss alone
- Drift produces no substantial  $v_3$  or  $v_4$  enhancement in our model
  - Possibly needs better sampling
  - Large in later-time eccentricities



# Drift $v_n$ sensitive to pathlength ordered internal medium props.

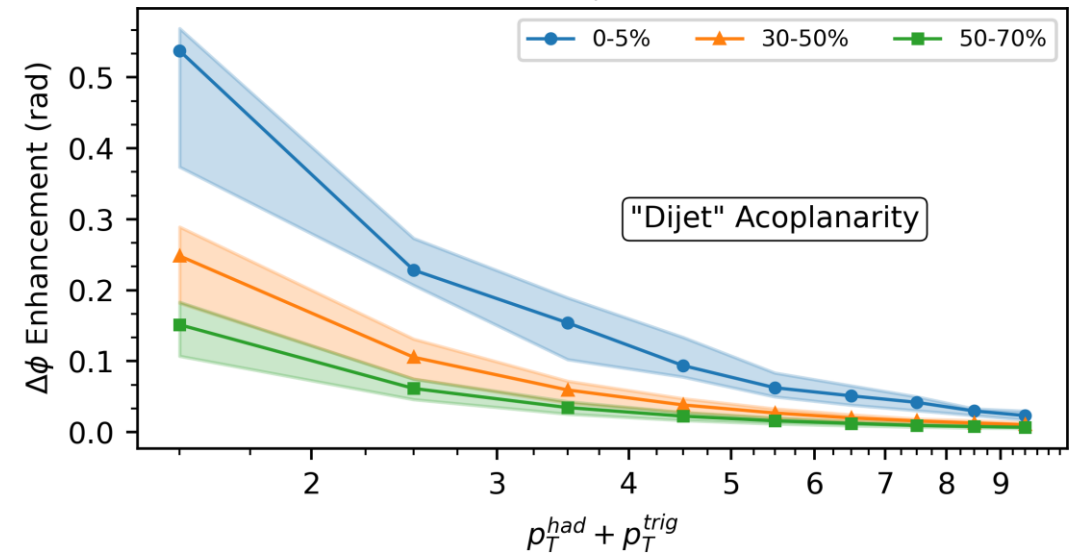
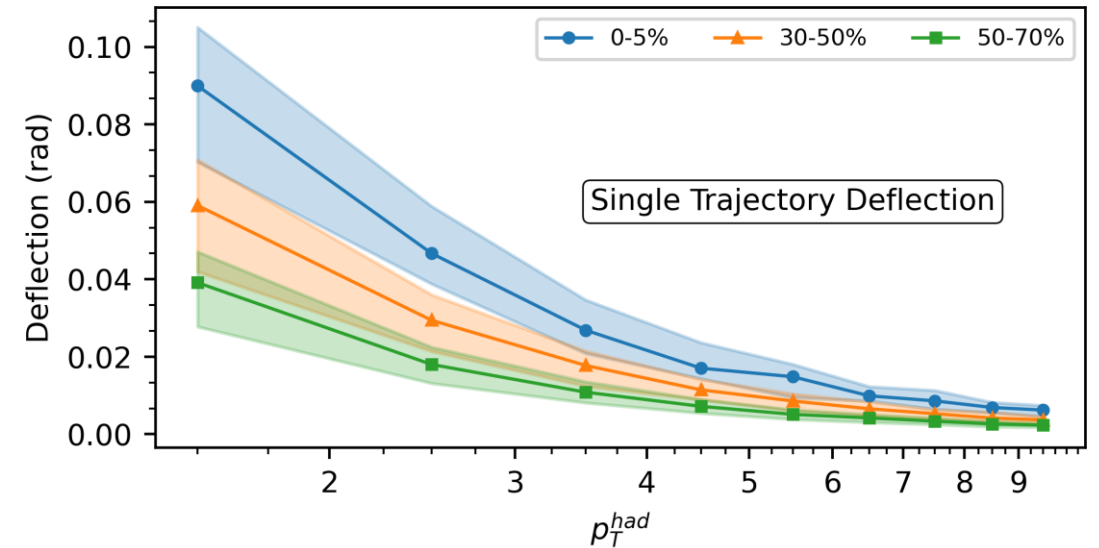
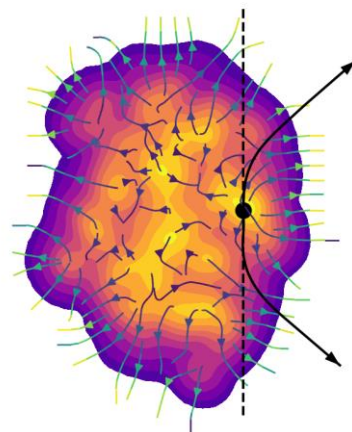
(iNSPIRE:2606181)

- Drift distinguishes between high anisotropy at early times vs at late times via interplay with energy loss
  - Powerful additional constraint on evolution dynamics!!!
- How would a Bayesian parameter extraction like Magdalena's differ with the inclusion of drift?
  - Possibly selects on slightly smaller anisotropy, likely changes story of free streaming parameters



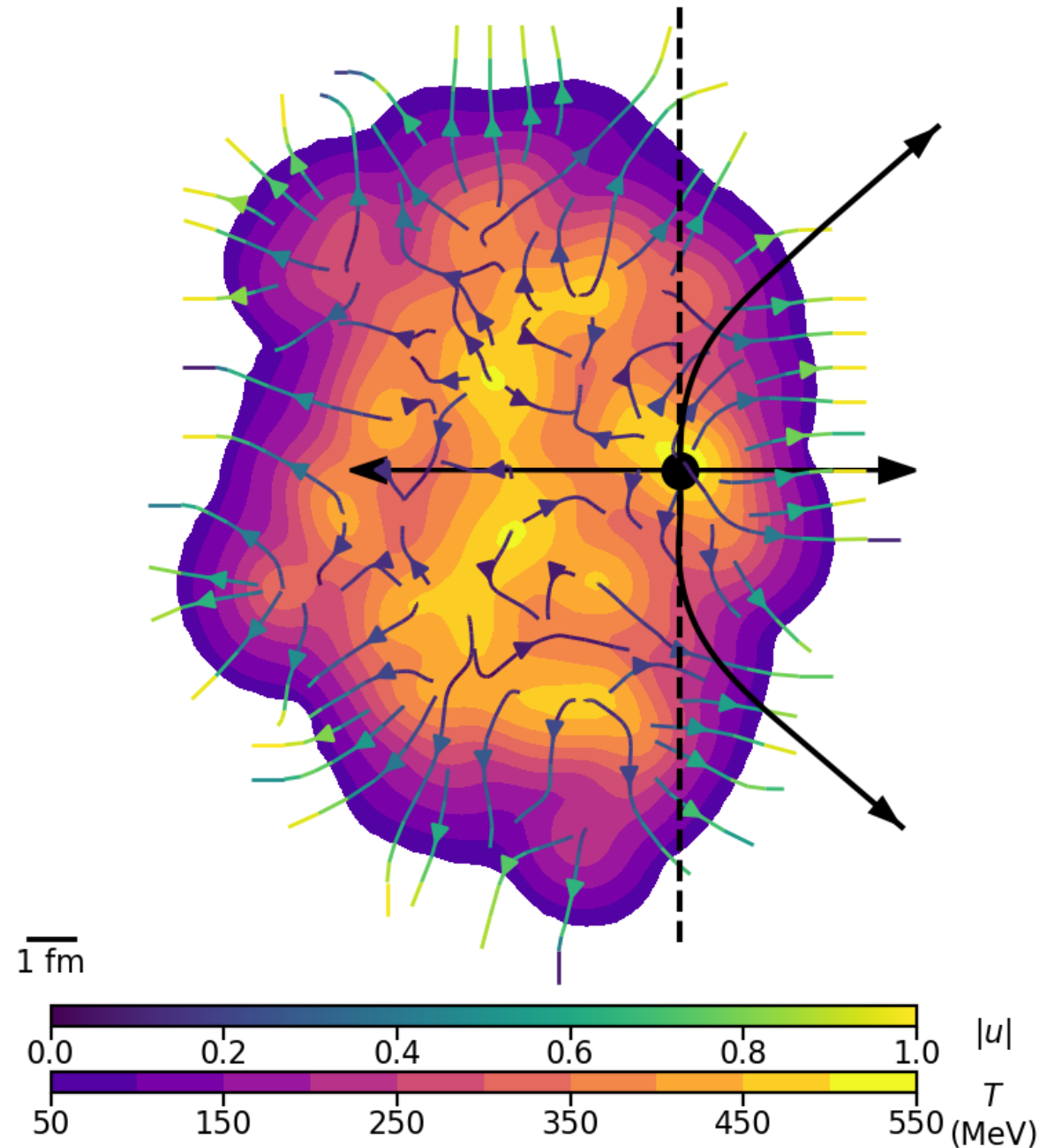
# Acoplanarity Enhancement RAD + CEL v. RAD

- Note centrality reversal
  - $V_2$  correlated to event plane
  - Acoplanarities access absolute deflections
- Initial acoplanarity fluctuation will change magnitude
  - Currently tree-level scattering: coplanar “dijets”



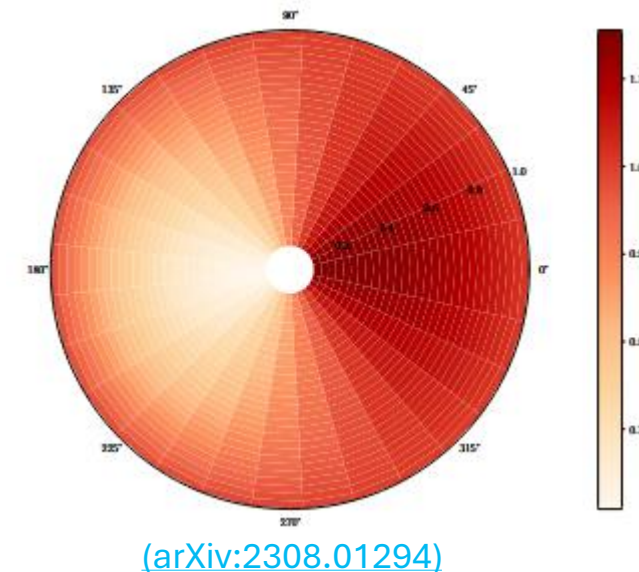
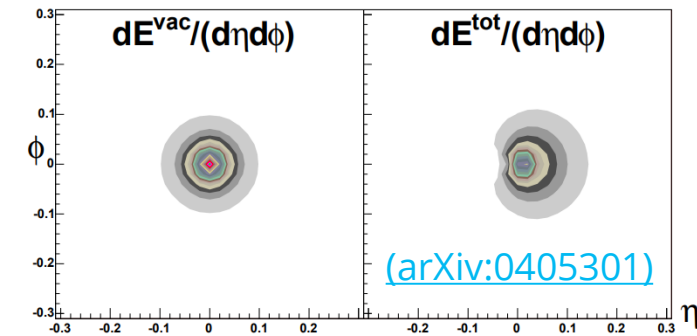
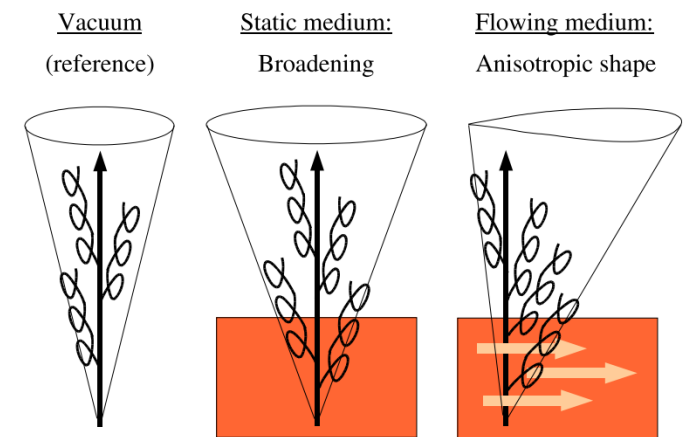
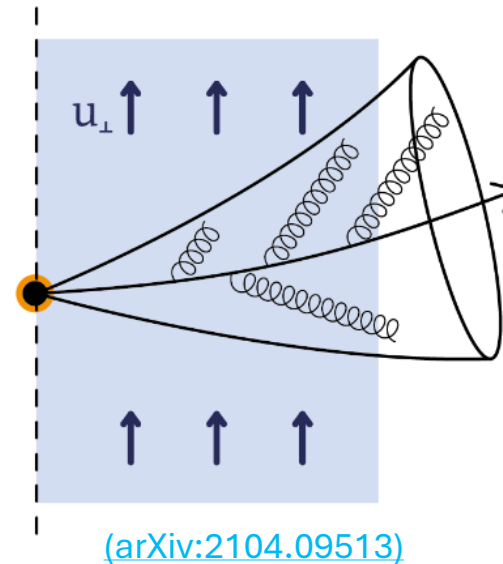
# Acoplanarity Cuts

- Acoplanarity modulation still tied to event plane
  - Possible selection cuts for velocity tomography
  - Difficult to measure independent of pathlength effects
- Collimation effect possible along event plane



# Anisotropic Jet Substructure

- People have been thinking about anisotropic radiation distribution for some time
  - Armesto, Salgado, & Wiedemann (2004)  
Flow anisotropy as Lorentz boost:  
<https://inspirehep.net/literature/651342>
  - Sievert, Sadofyev, Vitev (2021)  
Perturbative scalar calculation of radiation:  
<https://inspirehep.net/literature/1859289>
  - Barata, Milhano, & Sadofyev:  
Jet substructure harmonics:  
<https://inspirehep.net/literature/2684595>
  - Kuzmin & López  
Perturbative real gluon calculation of radiation:  
<https://inspirehep.net/literature/2801226>
- Drift of jet particles is also likely important!



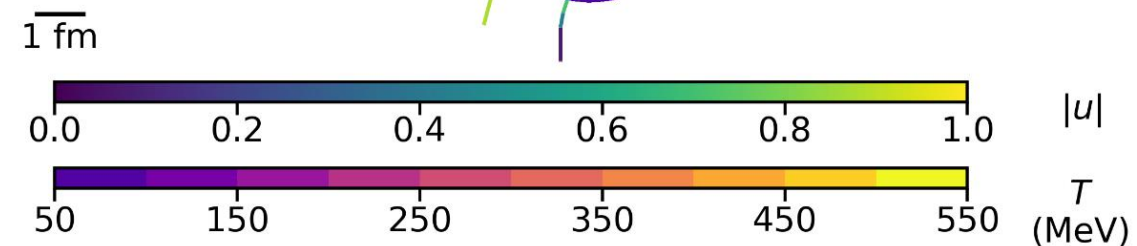
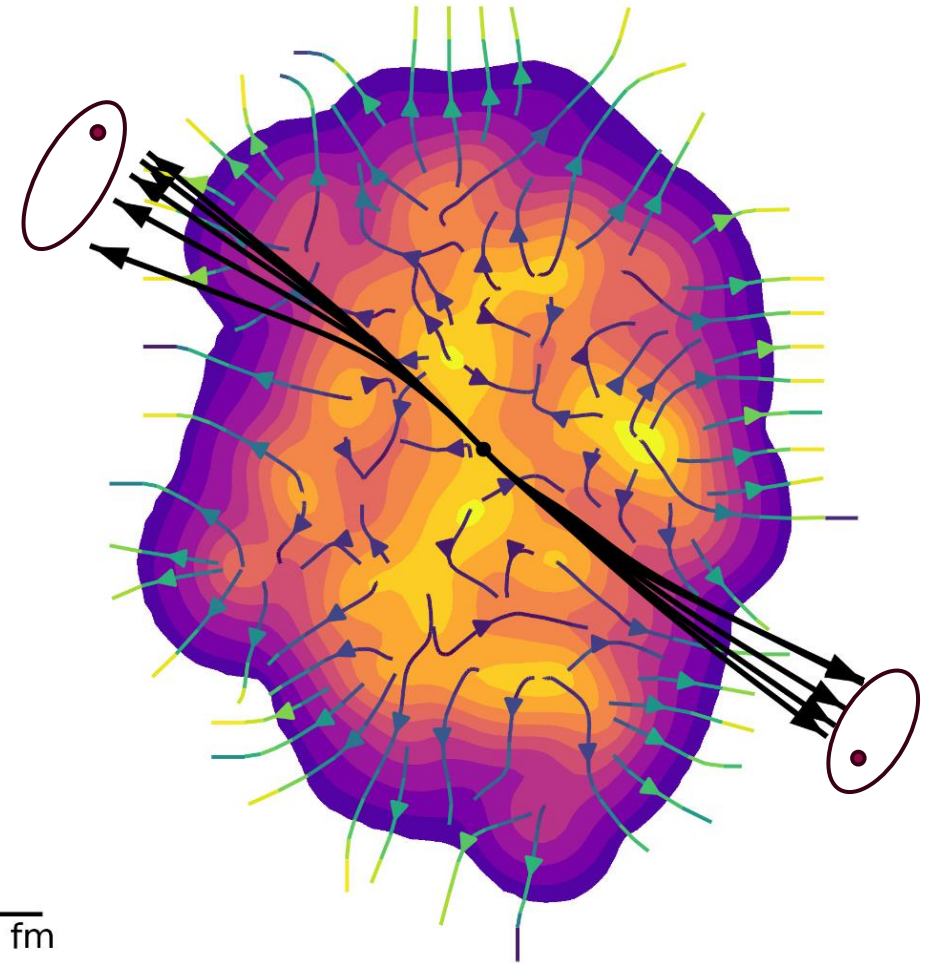


# Drift of Ensemble of Particles

- Energy suppression naturally produces dispersion of hard and soft particles within jet
- Sub-eikonal property a detriment for inclusive measurements, but well suited for jet substructure modification

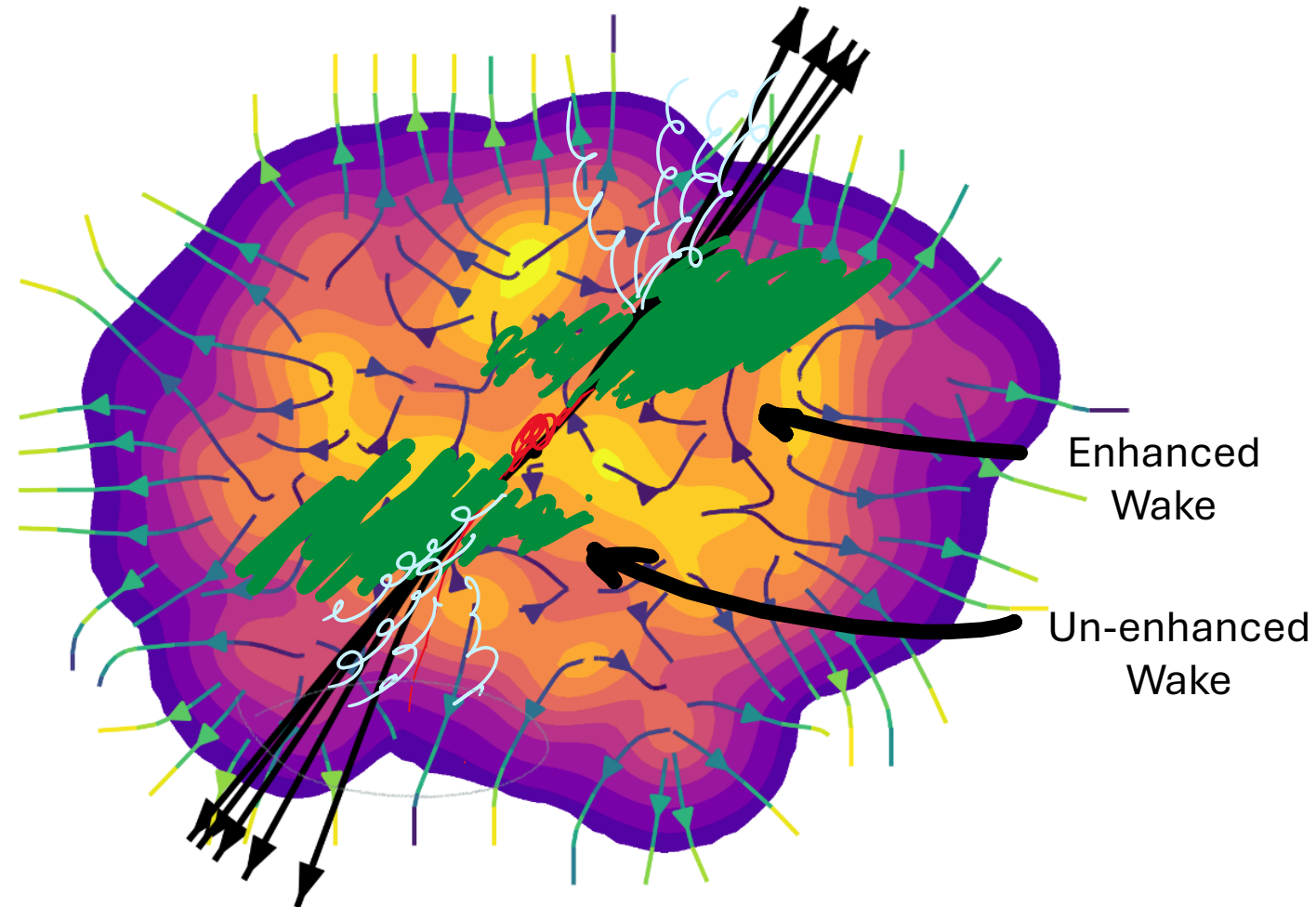
$$\langle \vec{q}_{drift} \rangle = \hat{e}_\perp \int d\tau \frac{3}{E(\tau)} \frac{\mu^2(\tau)}{\lambda(\tau)} \ln \frac{E(\tau)}{\mu(\tau)} \frac{u_\perp(\tau)}{1 - u_\parallel(\tau)}$$

Energy suppressed



# Anisotropic Jet Wake

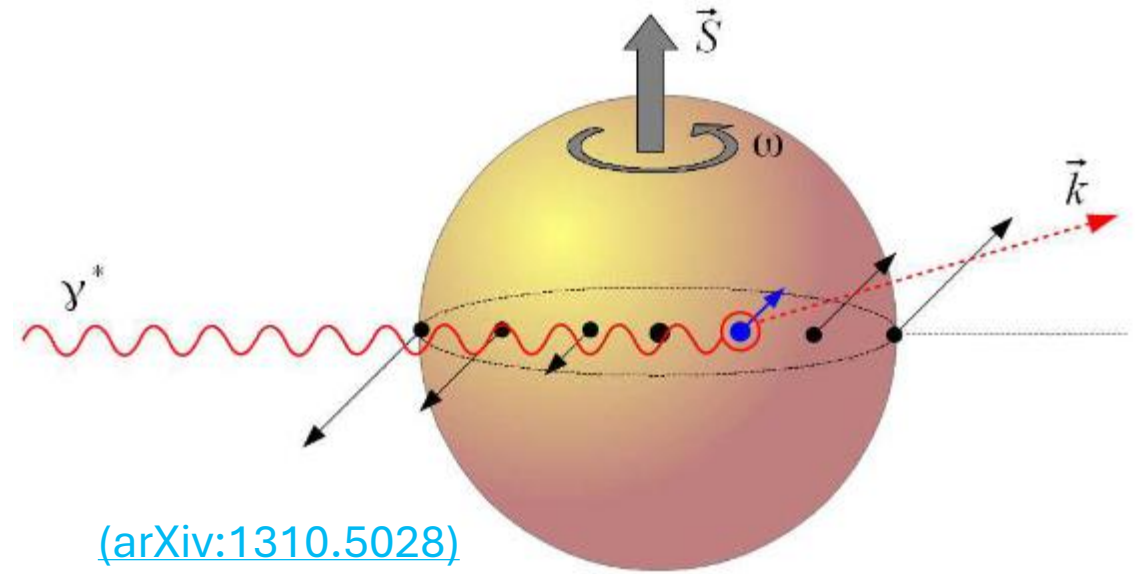
- Conservation of momentum – should be anisotropic wake
  - Constructive with induced radiation anisotropy
- How does this interplay with jet substructure anisotropies?
  - Possible anisotropic background contamination of jet substructure observables





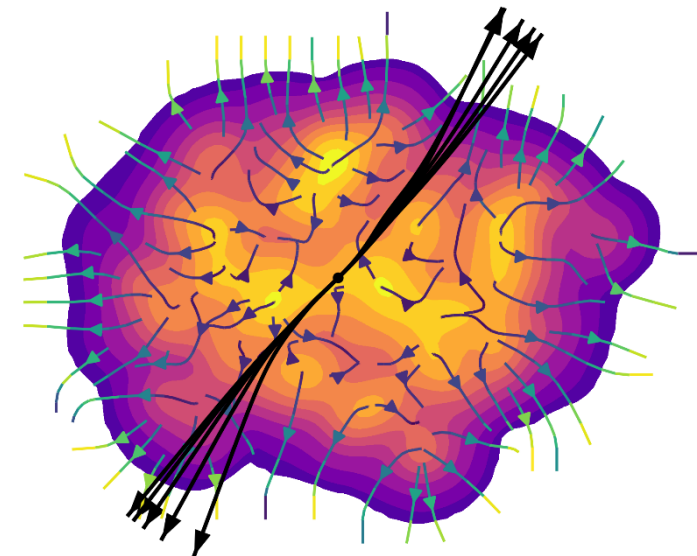
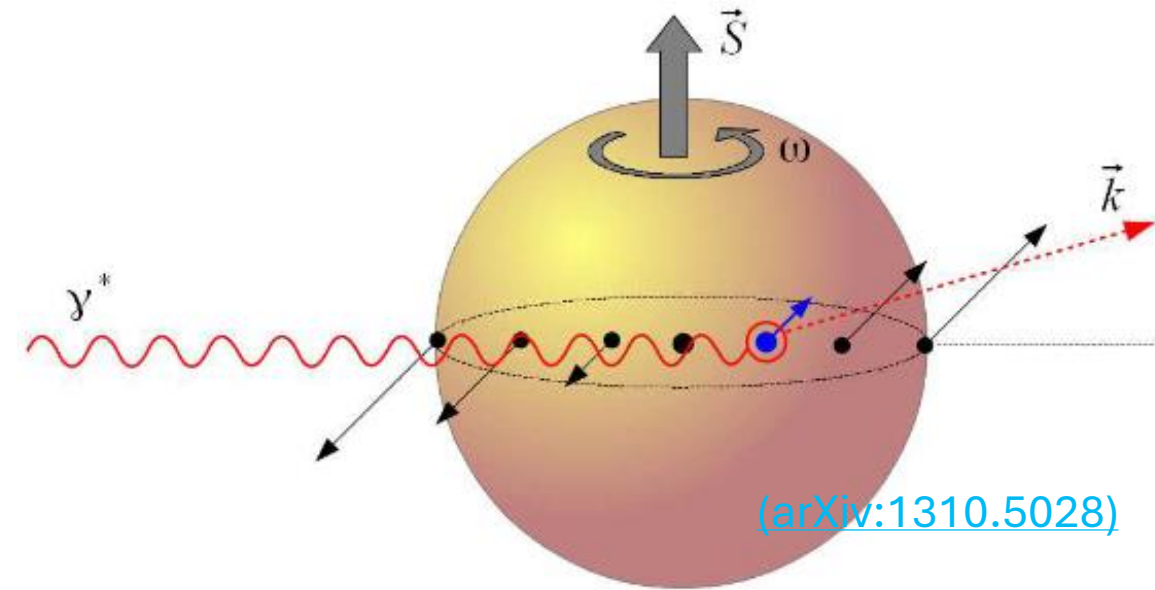
# Drift at the EIC

- Cold nuclear matter anisotropies can couple similarly to QGP flow
  - “Spin flow”, gradients, etc.
- Possible distinction between pre-equilibrium and equilibrium anisotropy
  - Could provide constraints on pre-equilibrium qhat
- Possible large impact on tomographic parameter extraction via hard probes
- Comparative laboratory for jet substructure dispersion



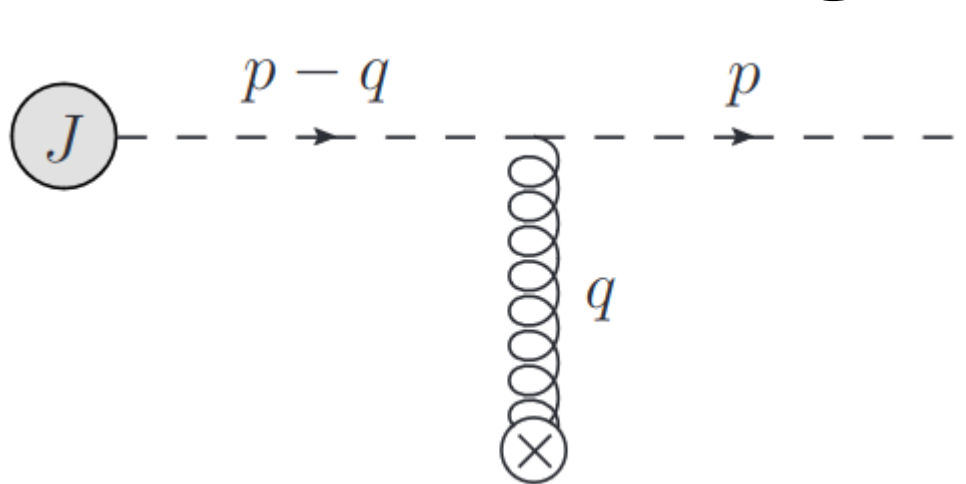
# Next Steps for Drift investigations

- CNM drift investigations
  - Drift due to polarized nuclei
  - Nicholas Baldonado, Alex Garcia, & Matt Sievert
- Jet Substructure
  - Drift-dispersion + flow-induced radiation => jet substructure anisotropy
  - Jo Bahder, Hasan Rahman, Matt Sievert, & Ivan Vitev

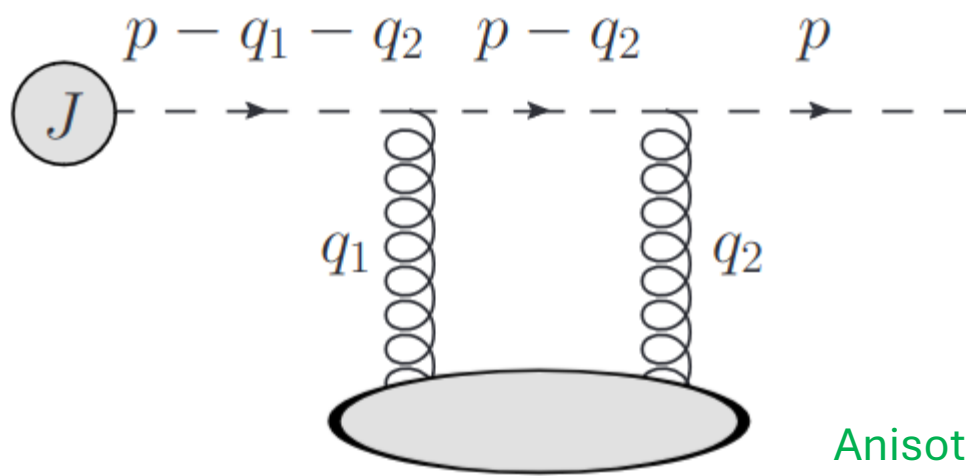


**Discussion!**

# Jet Broadening – Isotropic vs Anisotropic



A. V. Sadofyev, I. Vitev,  
& M. D. Sievert  
Phys.Rev.D 104 (2021)  
([arXiv:2104.09513](https://arxiv.org/abs/2104.09513))



Anisotropic

$$g a_i^{\mu a}(q) = t_i^a u_i^\mu v_i(q) (2\pi) \delta(q^0 - \mathbf{u}_i \cdot \mathbf{q})$$

Preferred direction!

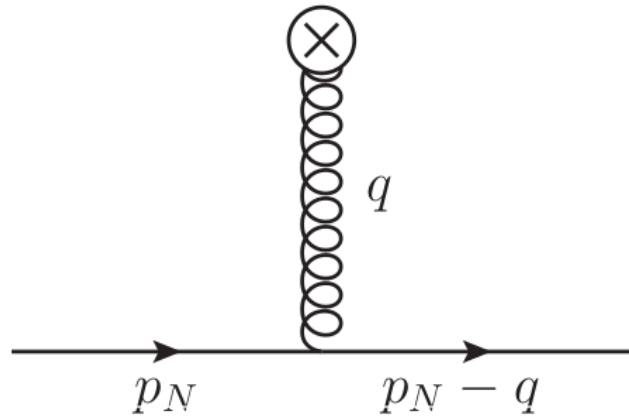
Isotropic

$$a_i^{\mu a}(q) = g^{\mu+} (t^a)_i \left[ 2\pi \delta(q^+) \right] \left[ \frac{-g_{eff}}{q_T^2} \right]$$

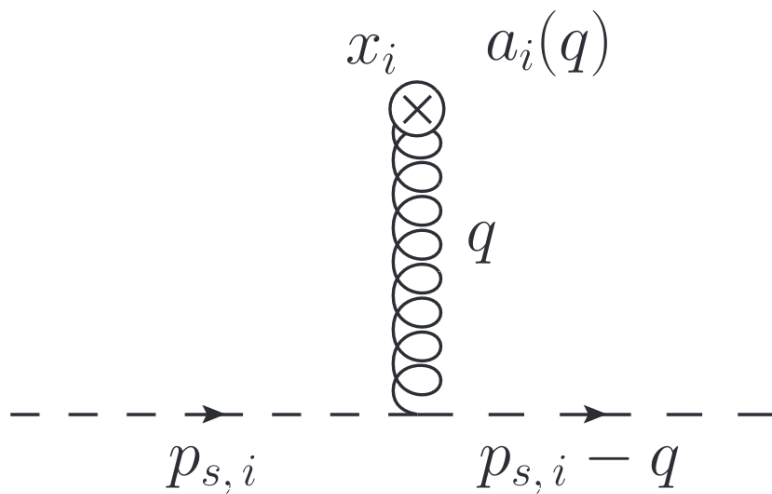
No vector info

- Difference in setup – constraints on external gluon field
- Isotropic
  - Jet quark moves lightcone (+), medium quark moves lightcone (-)
- Anisotropic
  - Jet quark moves lightcone (+), medium quark moves with medium flow
- Two vector directions associated with the medium
  - Flow
  - Gradients

# Medium Gluon Field Potentials



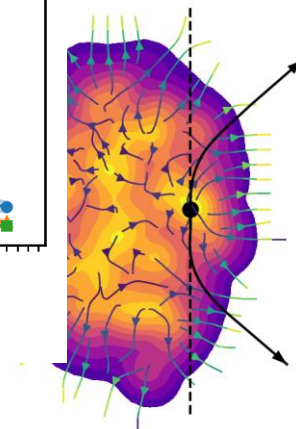
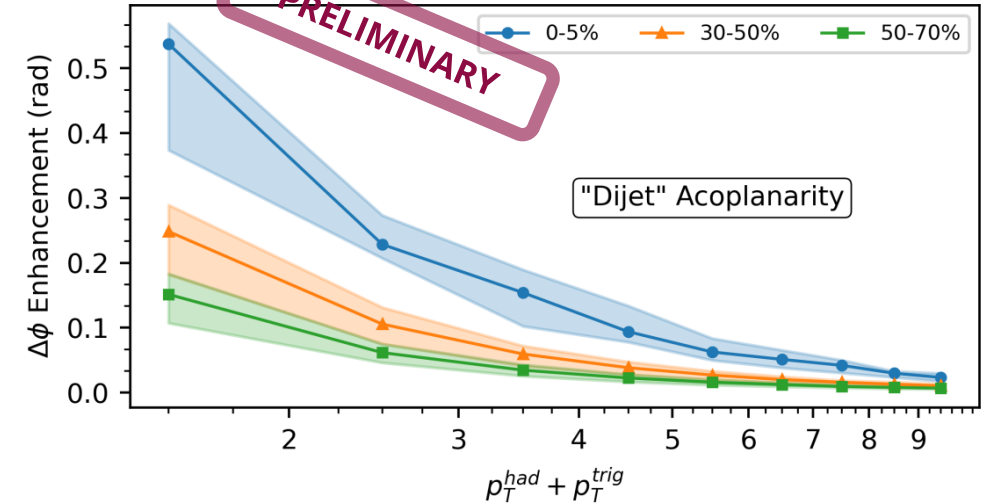
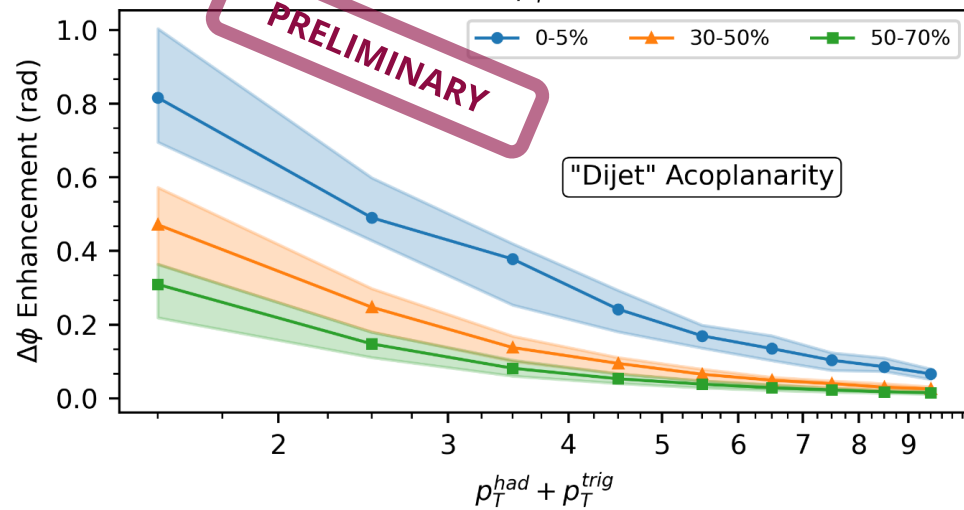
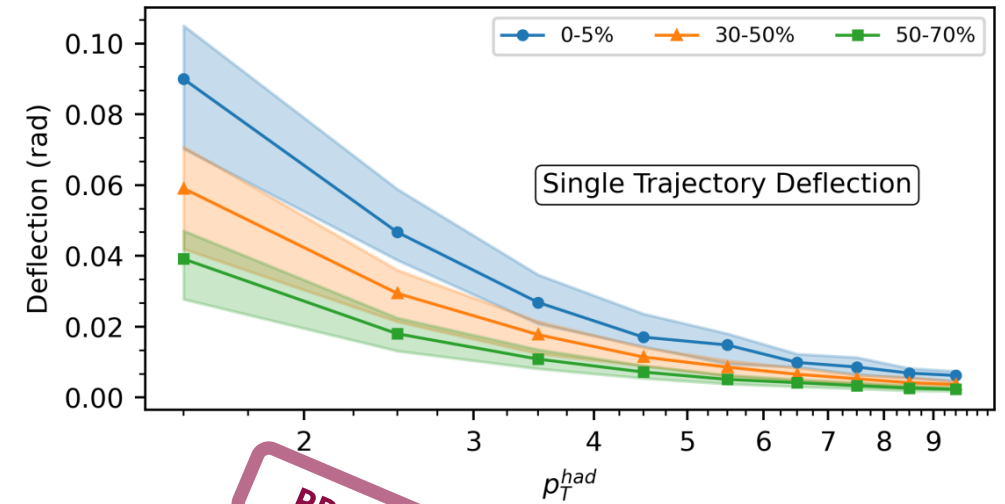
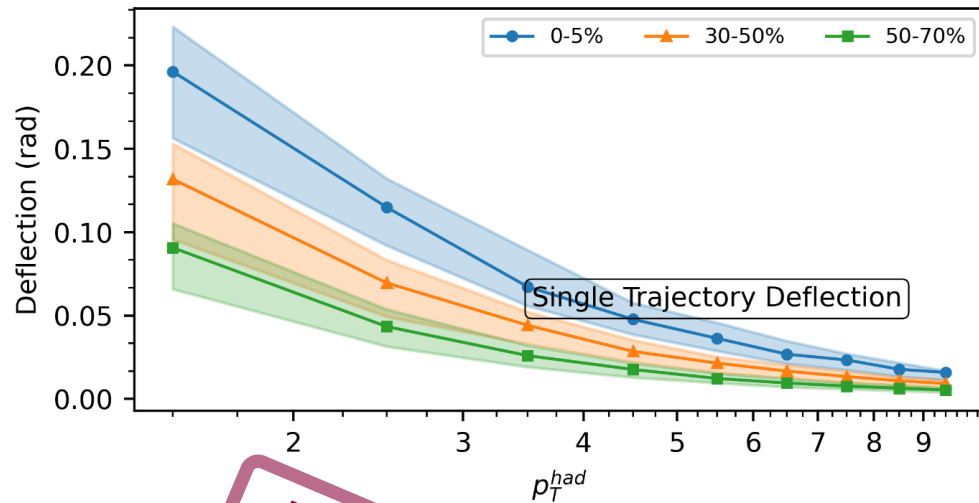
$$a_i^{\mu a}(q) = g^{\mu+} (t^a)_i \underbrace{[2\pi\delta(q^+)]}_{\text{Eikonal delta function (antiparallel)}} \underbrace{\left[ \frac{-g_{eff}}{q_T^2} \right]}_{\text{v(q) Isotropic scattering centers}}$$



$$g a_i^{\mu a}(q) = t_i^a u_i^\mu v_i(q) (2\pi) \delta(q^0 - \mathbf{u}_i \cdot \mathbf{q}) \left. \vphantom{g a_i^{\mu a}(q)} \right\} \text{Quark moves with medium}$$

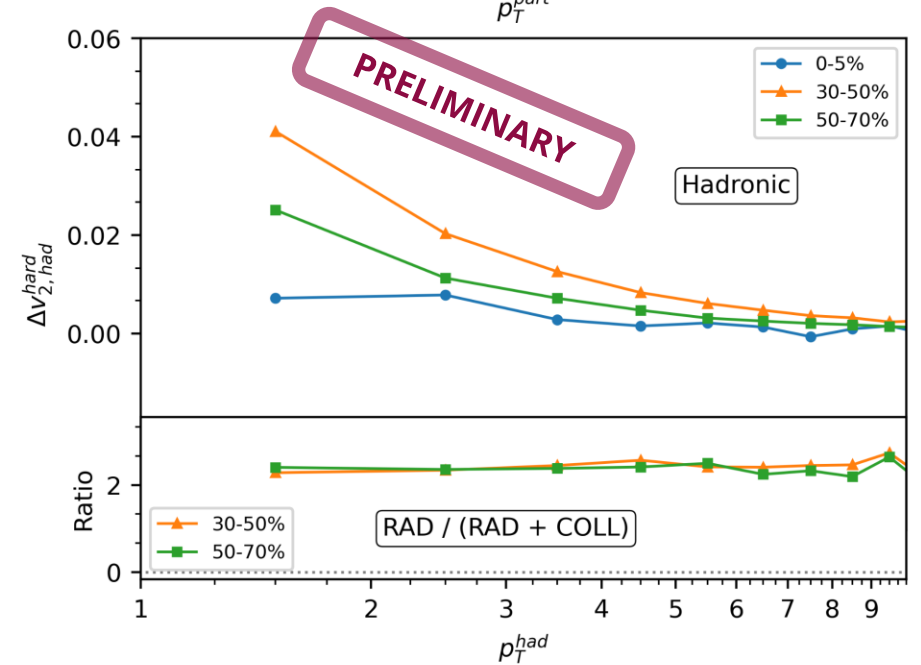
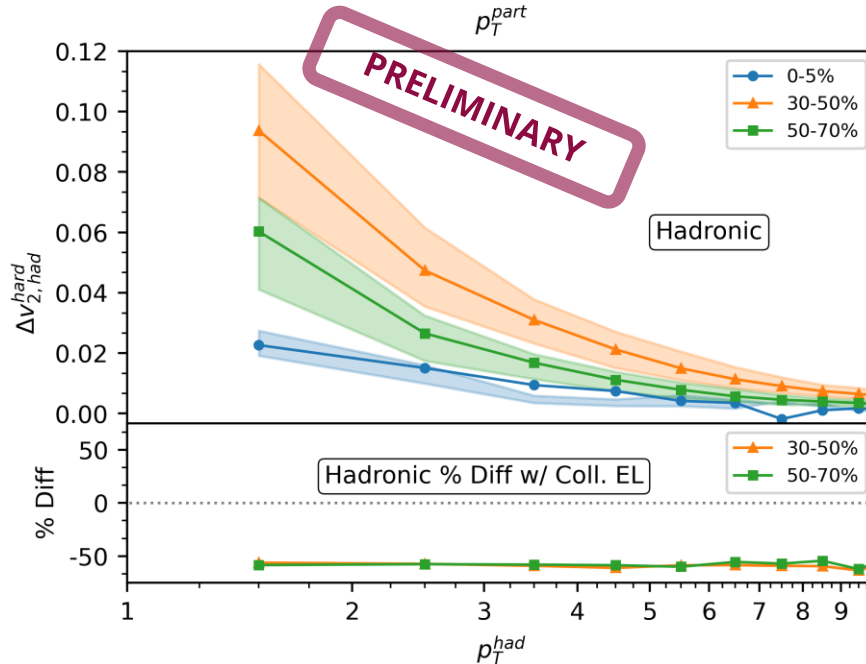
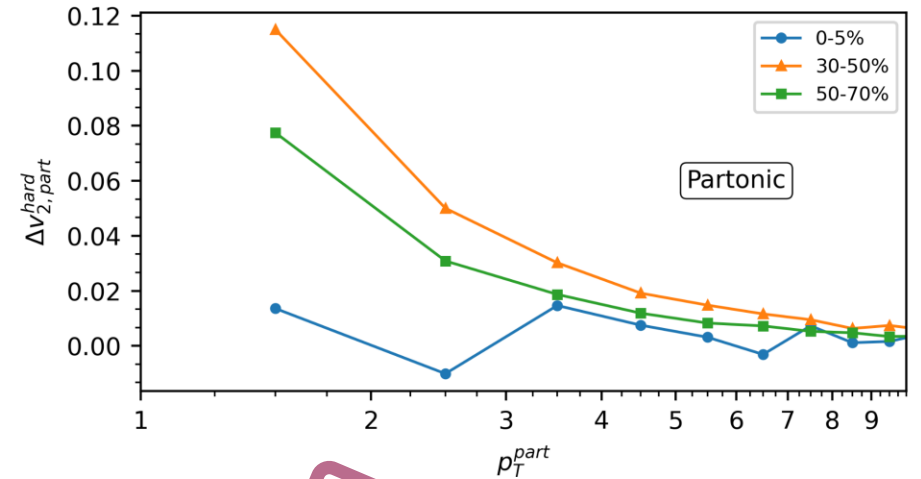
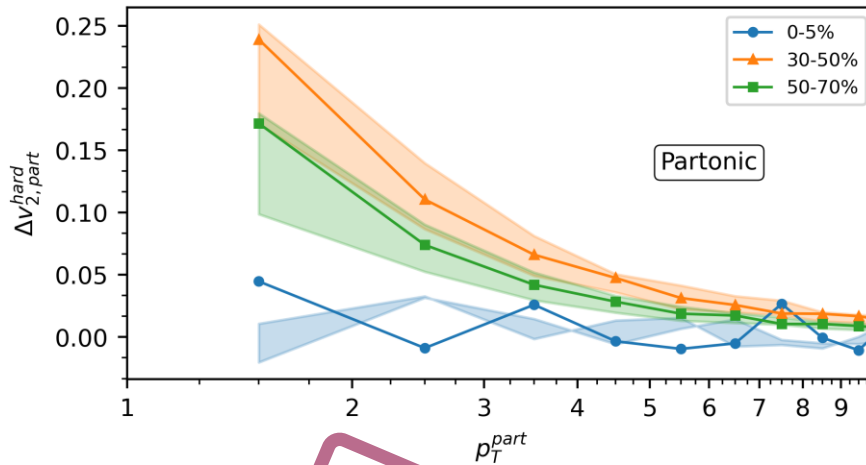
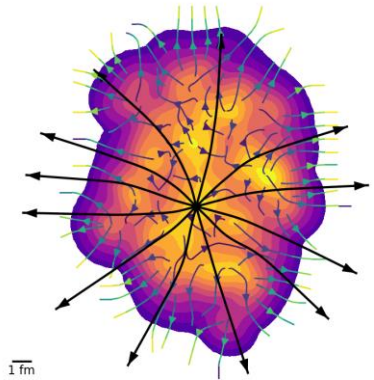
$$v_i(q) \equiv v_i(q^2 - (\mathbf{u}_i \cdot \mathbf{q})^2) \equiv \underbrace{\frac{-g^2}{q^2 + \mu_i^2 - (\mathbf{u}_i \cdot \mathbf{q})^2 - i\epsilon}}_{\text{v(q) Directional scattering centers}}$$

# Acoplanarity Enhancement RAD + CEL v. RAD





# v2 Enhancement RAD + CEL v. RAD



# CNM Effects

Implemented as in [arXiv:2009161](https://arxiv.org/abs/2009.1611)  
 Sans energy loss

$$R_{BA}(p_T) = \begin{cases} \frac{d\sigma^{dA}}{dyd^2p_T} / \frac{2A d\sigma^{pp}}{dyd^2p_T} & \text{in } d + A \\ \frac{dN^{AA}(b)}{dyd^2p_T} / \frac{T_{AA}(b) d\sigma^{pp}}{dyd^2p_T} & \text{in } A + A \end{cases}, \quad (1)$$

Compute CNM  $R_{AA}$  – Reweight distribution of hard processes

$$\left. \begin{aligned} & \frac{1}{2A} \frac{d\sigma^{dA}}{dyd^2p_T} \\ & \frac{1}{T_{AA}(b)} \frac{dN^{AA}(b)}{dyd^2p_T} \end{aligned} \right\} = K \sum_{abcd} \int dx_a dx_b \int d^2k_a d^2k_b g(\mathbf{k}_a) g(\mathbf{k}_b) \times S_A(x_a, Q_a^2) S_B(x_b, Q_b^2) \times f_{a/A}(x_a, Q_a^2) f_{b/B}(x_b, Q_b^2) \frac{d\sigma^{ab \rightarrow cd}}{d\hat{t}} \times \int_0^1 d\epsilon P(\epsilon) \frac{z_c^*}{z_c} \frac{D_{h/c}(z_c^*, Q_c^2)}{\pi z_c}. \quad (2)$$

PDFs

CNM Energy Loss Prob. (Not Used)

FFs

Shadowing Function