You cannot extract neutron-skin thickness from coherent π^0 photoproduction off nuclei

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Introduction : Equation of State for nuclear matter

- Symmetry energy
- Neutron skin

Impulse Approximation in a Nutshell

- PWIA
- DWIA

3 Coherent π^0 photoproduction on

- ¹²C
 ⁴⁰Ca
- ¹⁰Ca
 ²⁰⁸Pb



Introduction

Nuclear force governs the structure of nuclei as well as that of neutron stars

However, to understand the formation of neutron stars, there is no need for microscopic calculations but we have to understand nuclear matter \Rightarrow (nuclear) Equation of State (EoS)

For nuclear matter, the state variables are Z : proton number, N : neutron number or in infinite matter $\alpha = (N - Z)/A$, the n-p asymmetry ρ the density

$$\epsilon(\rho, \alpha) = \epsilon(\rho, \alpha = 0) + S(\rho) \alpha^2 + \dots$$

where *S* is the symmetry energy *S* characterises the increase in energy from N = Z to neutron matter Recent reviews : [Horowitz *et al.* JPG 41, 093001 (2014)] [Thiel *et al.* JPG 46, 093003 (2019)]

Symmetry energy

Taylor expanded around $\rho = \rho_0$:

$$S(\rho) = S_v + \frac{L}{3} \left(\frac{\rho - \rho_0}{\rho_0} \right) + \frac{1}{18} K_{\text{sym}} \left(\frac{\rho - \rho_0}{\rho_0} \right)^2 + \dots$$

S can be constrained from nuclear experiments (laboratory) Idea : measuring the neutron-skin thickness of ²⁰⁸Pb R_{skin} (or Δr_{np})



[Roca-Maza et al. PRL 106, 252501 (2011)]

*R*_{skin} : balance between surface tension and symmetry term

Neutron-skin thickness

• Coherent π^0 photoproduction

[Tarbert et al. PRL 112, 242502 (2014)]

 γ + ²⁰⁸Pb \rightarrow ²⁰⁸Pb + π^0

- Measurement of the electric dipole polarizability [Tamii et al. PRL 107, 062502 (2011)]
- Parity-violating electron scattering PREX [Abrahamyan *et al.* PRL 108, 112502 (2012)] PREX-II [Adhikari *et al.* PRL 126, 172502 (2021)]

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 PREX-II [Adhikari *et al.* PRL 126, 172502 (2021)]

• PREX :

thick $R_{\rm skin}$ with significant uncertainty

 Coherent π⁰ photoproduction : thin R_{skin} with little uncertainty



0.1 0.2 0.3 0.4 0.5 0.6

[J. Piekarewicz]

$\gamma + {}^{208}{ m Pb} \rightarrow {}^{208}{ m Pb} + \pi^0$ [Tarbert *et al.* PRL 112, 242502 (2014)]

- Measurement done at MaMi (Mainz Microtron)
- γ produced by electron beam
- π⁰ decay into 2γ detected in Crystal Ball and TAPS



$^{208}\text{Pb} \rightarrow ^{208}\text{Pb} + \pi^0$ [Tarbert *et al.* PRL 112, 242502 (2014)]

- Measurement done at MaMi (Mainz Microtron)
- γ produced by electron beam
- π⁰ decay into 2γ detected in Crystal Ball and TAPS
- Precise measurement
- Analysis :
 - Impulse Approximation
 - with simple Fermi density
 - FSI plays a role
- Deduce very thin $R_{skin}(^{208}Pb)$ = $0.15 \pm 0.03(stat)^{+0.01}_{-0.03}(syst)$ fm also very precise
- How reliable is this result?



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Summary

Plane Wave Impulse Approximantion

Coherent π^0 photoproduction

 $\gamma + A \rightarrow A + \pi^0$

Plane Wave : No FSI in exit chanel π^0 -A

At the Impulse Approximation : production of π^0 on one single nucleon \Rightarrow coherent sum on each nucleon

$$\frac{d\sigma}{d\Omega} \propto \left| f_2(\vec{k}_{\pi},\vec{k}_{\gamma}) \, \rho_{\rm A}(q) \right|^2$$



[Colomer, PhD (2020)]

- f2 : CGLN amplitudes from MAID [Drechsel et al. EPJA 34, 69 (2007)]
- $\rho_{\rm A}$: nucleus form factor
- \Rightarrow Should give access to nuclear density, but
 - FSI
 - Higher-order effects [Miller PRC 100, 044608 (2019)]



[Colomer, PhD (2020)]

$$\frac{d\sigma}{d\Omega} \propto \left| f_2(\vec{k}_{\pi}, \vec{k}_{\gamma}) \rho_{\rm A}(q) + \int \frac{d\vec{k}_{\pi}}{2\mathcal{M}} \frac{T_{\pi\rm A}(\vec{k}_{\pi}, \vec{k}_{\pi}) f_2(\vec{k}_{\pi}, \vec{k}_{\gamma}) \rho_{\rm A}(q)}{E(k_{\pi}) - E(k_{\pi}') + i\epsilon} \right|^2$$

 $T_{\pi A}$: computed following [Carr *et al.* PRC 25, 952 (1982)]

 $\Rightarrow d\sigma/d\Omega$ no longer exactly $\propto |
ho_{
m A}|^2$

We test the model

- on different targets (¹²C, ⁴⁰Ca, ²⁰⁸Pb)
- considering different densities with different R_{skin}
- compare to data [Krusche et al. PLB 526, 287 (2002)]

[F. Colomer et al. PRC 106, 044318 (2022)]

¹²C target



[F. Colomer, PhD (2020) & PRC 106, 044318 (2022)]

- Use different nucleonic densities
 - São Paulo [Chamon et al. PRC 66, 014610 (2002)]

Electron scattering (using $\rho_n = \rho_p$) [Dreher *et al.* NPA 235, 219 (1974)] with different neutron skins

- São Paulo : $R_{skin} < 0 (!!)$
- Electron scattering : $\rho_n = \rho_p \Rightarrow R_{skin} = 0$ fm

¹²C target

Exp : [Krusche et al. PLB 526, 287 ('02)]



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- DWIA : distortion increases cross section

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 - FSU calculations [Todd-Rutel & Piekarewicz, PRL 95, 122501 (2005)] with different neutron skins
 - São Paulo : $R_{skin} = -0.30$ fm
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²⁰⁸Pb target



[F. Colomer, PhD (2020) & PRC 106, 044318 (2022)]

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- São Paulo : $R_{skin} = 0.101$ fm
- FSU : $R_{skin} \in [0.176, 0.286]$ fm

²⁰⁸Pb

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- DWIA : cross sections increase ⇒ better agreement with data sensitivity to density is reduced (effect of FSI)

Summary

- Nuclear EoS is key to understand the structure of neutron stars
- Constrain Symmetry term from the neutron-skin thickness R_{skin}
 - PREX I & II : $R_{skin} = 0.29 \pm 0.07$ fm
 - [Adhikari et al. PRL 126, 172502 (2021)]
 - π^0 photoproduction : $R_{skin}(^{208}\text{Pb}) = 0.15 \pm 0.03(\text{stat})^{+0.01}_{-0.03}(\text{syst})$ fm

[Tarbert et al. PRL 112, 242502 (2014)]

- We test this hypothesis with a new DWIA code of the reaction [F. Colomer, PhD (2020) & PRC 106, 044318 (2022)]
- We obtain good agreement with data, especially on Pb
- FSI are significant
- $\sigma^{\gamma\pi^0}$ is purely isoscalar : not sensitive to R_{skin} $R_{skin}(^{208}\text{Pb}) \in [0.1, 0.3]$ fm provide same cross section

 \Rightarrow You cannot infer R_{skin} from π^0 photoproduction

Summary in a nutshell



You cannot infer R_{skin} from π^0 photoproduction

Comparing Krusche and Tarbert data

 $E_{\gamma} = 200 MeV$



[Tarbert et al. PRL 112, 242502 (2014)]