

#### The r-mode puzzle



a Stationary reference frame

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# Pulsar periods and Keplerian mass-shedding limit

- fastest observed spin to date: 716 Hz for PSR J1748-2446ad
- Accretion can spin up a pulsar, theoretically up to the Keplerian mass-shedding limit:

$$
f_K \sim \frac{1}{2\pi} \sqrt{\frac{2GM}{R^3}} \sim 2 \text{ kHz or } P_K \sim 500 \,\mu\text{s}
$$





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[Andrew Steiner] What could be limiting the spins of accreting pulsars? Magnetic torques could limit the spin up due to accretion (also r-modes?) [Andersson, Kokkotas & Schutz, 1999, Bildsten, 1998]

# r-modes and the CFS instability

[C. Hanna & B. Owen]

- r-modes: fluid modes of oscillation, restored by the Coriolis force
- generically unstable due to the emission of gravitational waves (CFS instability) [Chandrasekhar, 1970; Friedman & Schutz, 1978; Andersson 1998; Friedman & Morsink, 1998]

for the dominant  $l = m = 2$  r-mode  $σ<sub>R</sub> = (2/3)Ω$  (corotating frame)  $\sigma = \sigma_R - m\Omega = -(4/3)\Omega$  (inertial frame) Instability condition: a counter-rotating mode appears as corotating for an inertial observer UFABC



solar rossby waves





rossby waves on Earth

# Neutron star r-modes and the instability window

- Linear instability; nonlinear couplings should set the saturation amplitude [Arras et al., 2003; Bondarescu, Teukolsky & Wasserman, 2007]
- \* Shear and bulk viscosity shape the instability window The r-mode instability window will be found by solving for the zeros of

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$$
\frac{1}{t} = -\frac{1}{\tau_{gw}} \left(\frac{1 \text{ ms}}{P}\right)^{p_{gw}} + \frac{1}{\tau_{bv}} \left(\frac{1 \text{ ms}}{P}\right)^{p_{bv}} \left(\frac{T}{10^9 \text{ K}}\right)^6 + \frac{1}{\tau_{sv}} \left(\frac{10^9 \text{ K}}{T}\right)^2
$$

[Lindblom, Owen & Morsink, 1998; ! *P* 1 ms"*<sup>p</sup><sup>d</sup>* Andersson, Kokkotas & Schutz, 1999]



[Kokkotas, 2003]

*typical values:*  $\tau_{gw} \sim 20 \text{ s}, \tau_{bv} \sim 10^{10} \text{ s}, \tau_{sv} \sim 10^8 \text{ s},$ <br>  $p_{gw} \sim 6, p_{bv} \sim 2, p_{sv} \sim 0$  $p_{gw}$  ~ 6,  $p_{bv}$  ~ 2,  $p_{sv}$  ~ 0

#### What happens in the r-mode instability?

- the r-mode grows exponentially until it reaches a saturation amplitude α
- the star heats up due to the unstable mode
- the r-mode spins down the star
- the star cools down

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accretion spins up the star again

[Levin, 1999, Andersson et al. 2000] Equilibrium between spin up and spin down is unstable: neutron star follows a limit cycle



### Neutron stars observed in the instability window

- Neutron stars in LMXBs and MSRPs are found in the r-mode instability window
- Should they be observed there? Some stars are so deep into the window that they should have large  $\alpha$  ( $\alpha > 10^{-3}$ )
- For large (small) values of  $\alpha$ , the star should spend less than  $1\,\%$  (30  $\%$ ) of the time inside the window
- Number of sources that should be on:

$$
N_{\rm on} = \frac{\tau_{\rm on}}{\tau_{\rm cycle}} \sim 5 \times 10^{-8} \alpha^{-1.6} N_{\rm sample}
$$

[Heyl, 2002]

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temperature information is modeldependent and uncertain…

## Duration of the GW emission in the r-mode cycle

The time for the system to complete the circuit depends only weakly on A (saturation amplitude), here 3.7 – 4.8 Myr.

Curves are labelled with the duration of the GW emission.

The duty cycle ranges from  $\sim 30\,\%$  to  $\sim 0.0001\,\%$ .





## Uncertainties in the temperature estimates

#### Factors:

- uncertain envelope composition (considered in this plot)
- some systems may still be thermally relaxing to a steady state after an outburst and may have sizeable temperature gradients in the crust [Brown & Cumming, 2009]
- resulting factor of a few uncertainty seems to have no qualitative impact [Haskell, Degenaar and Ho, 2012]





#### r-mode windows? Alternative models and uncertainties



# Corrections to the r-mode frequency

- \* For the  $\ell = m = 2$  r-mode, the observed gravitational wave frequency  $\sigma$  is related to the rotating frame frequency  $\sigma_R$  by  $\sigma = \sigma_R - 2\Omega = (\kappa - 2)\Omega, \quad \kappa \equiv$ *<sup>R</sup>*  $\Omega$
- For a slowly rotating Newtonian star,  $\sigma_R =$  $2m\Omega$  $\frac{1}{\ell(\ell+1)} = (2/3)\Omega$
- Relativistic effects are quantified by the compactness *M/R.*
- The effect of fast rotation can be expressed as  $\kappa = \kappa_0 + \kappa_2$  $\Omega^2$  $\pi G \bar{\rho}_0$
- If we know  $\kappa_0$  and  $\kappa_2$  as functions of the compactness and observe an r-mode, we can solve for *M/R*!



# Corrections to the r-mode frequency





#### Observations of r-modes in LMXBs? XTE J1751-305 and 4U 1636-536

- Strohmayer & Mahmoodifar reported in 2014 oscillations consistent with r-modes in RXTE data:
	- \* XTE J1751-305:  $\kappa$  = 0.5727597 and  $f$  = 435 Hz (found in the 2002 discovery outburst)
	- $*$  4U 1636-536:  $\kappa$  = 0.5645388 and  $f$  = 582 Hz (found in the 2001 thermonuclear superburst)
- More observations? NICER should be able to see these oscillations, if they are there! (But see also "Where are the r-modes?" by Mahmoodifar and Strohmayer, 2017).



#### Caveats: constraints from measured spin-down and alternatives

- Problem: the r-mode amplitude derived from the observed modulation amplitude leads to a spin down rate too large to be consistent with the observations [Strohmayer & Mahmoodifar 2014]
- Alternatives: rotationally modified g-modes, r-modes modified by a solid crust, surface g-modes or r-modes, unstable toroidal crustal modes or amplification at the surface of a core r-mode with lower amplitude [Andersson, Jones & Ho, 2014; Lee, 2014]

#### Questions:

shape of the instability window?  $r$ -mode amplitude  $\times$  signal amplitude?





## An application: results for XTE J1751-305

#### Using

 for fully relativistic uniform *κ*0 density stars [Lokitch, Friedman & Andersson, 2003]

+  $κ$ <sub>2</sub> for fast rotating Newtonian polytropes [Lindblom, Mendell & Owen, 1999]



[Anderson, Jones & Ho, 2014]



#### More *M-R* constraints from the (possible) r-mode frequencies





[Chirenti & Jasiulek, 2018]

For a radius range of  $11 - 13$  km, with the assumed r-mode detections and APR4 + SLY we find the mass estimates:

 $XTE$  J1751-305: 1.48 – 1.56  $M_{\odot}$  4  $U$  1636-536: 1.60 – 1.68  $M_{\odot}$ **JFABC** 

#### Conclusions Puzzles

Puzzle 1: What limits the spin of accreting neutron stars? Effect of magnetic torques in accretion process? r-modes?

Puzzle 2: Why are there (so many) neutron stars in LMXBs and MSRPs that seem to be in the r-mode instability window? measurements? Modeling of the window? Physical assumptions? *T*

Puzzle 3: r-modes in XTE J1751-305 and 4U 1636-536? Modified r-modes? g-modes? mode amplitude  $\times$  amplitude of X-ray modulation?

Puzzle 4: What is the r-mode saturation amplitude? Will r-modes be detectable with LIGO? With 3G detectors?

Recommendation for observers: Look for r-modes in stars inside the window!

