

# From Low to High Densities: an Application of Bayesian Model Mixing to the Dense Matter EOS

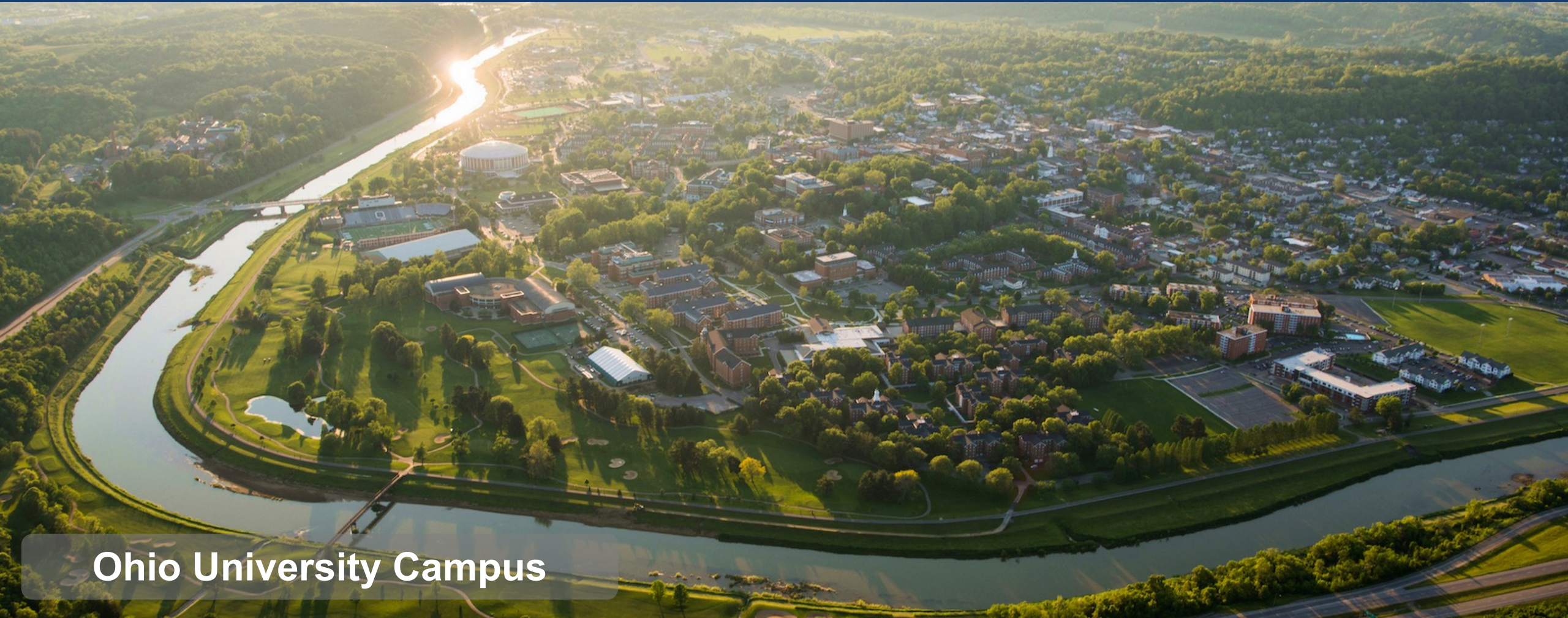
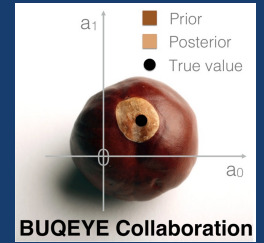
Christian Drischler (drischler@ohio.edu)

INT-24-89W: EOS Measurements with Next-Generation GW Detectors

August 26, 2024 | Institute for Nuclear Theory (INT)

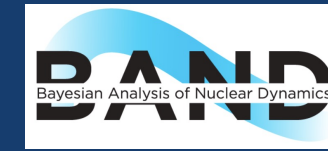


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# Bayesian Model Mixing in SNM



A **Bayesian mixture model** approach to quantifying the *empirical* nuclear **saturation point**

CD, Giuliani, Bezoui, Piekarewicz, and Viens, arXiv:2405.02748

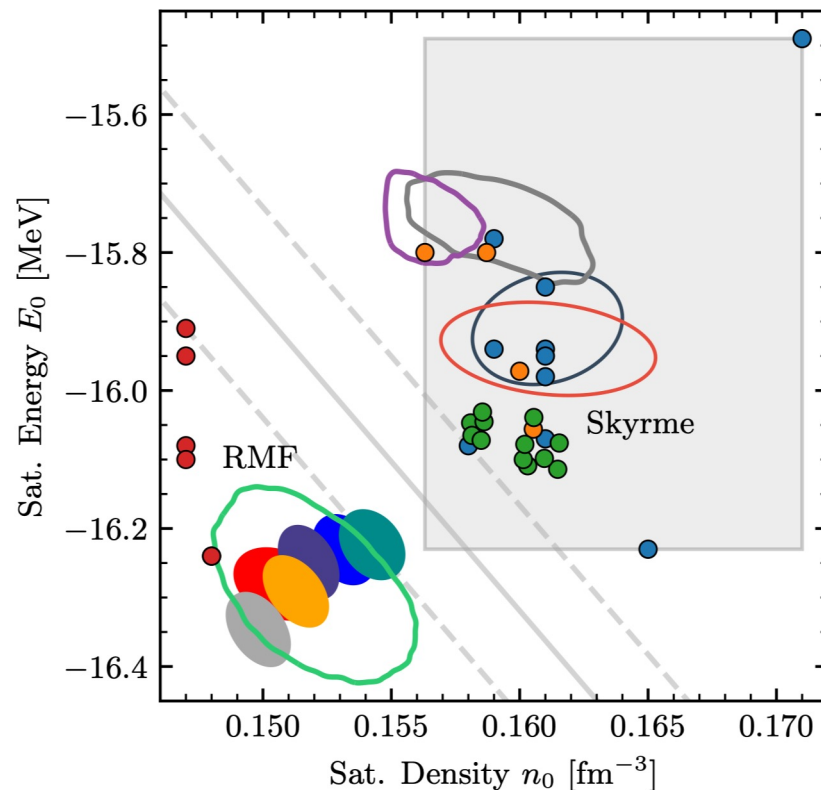
**Goal: rigorous benchmarks of saturation properties of chiral NN+3N interactions (using Skyrme & RMF models)**



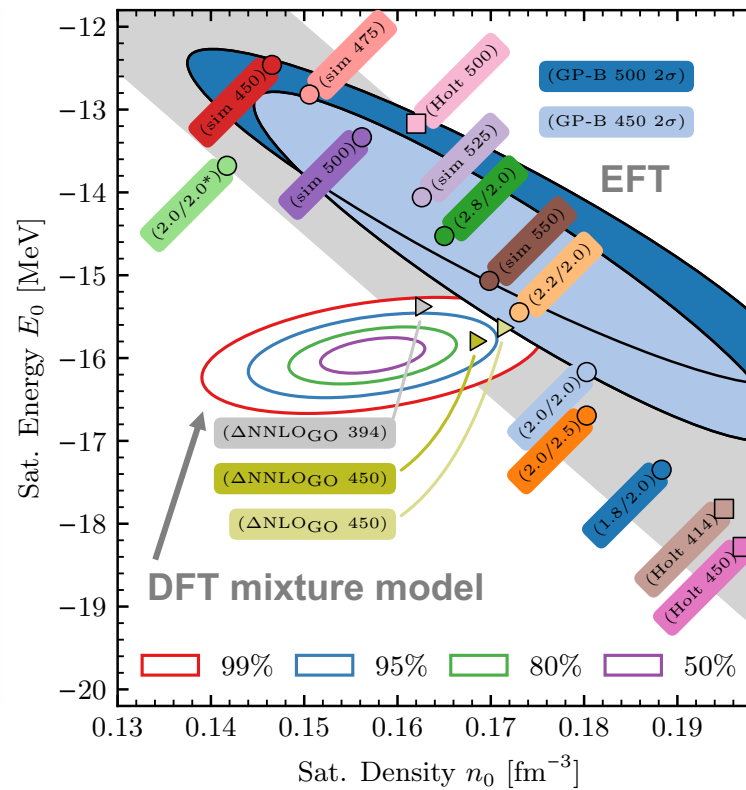
From chiral EFT to perturbative QCD: a **Bayesian model mixing** approach to symmetric matter

Semposki, CD, Furnstahl, Melendez, and Phillips, arXiv:2404.06323

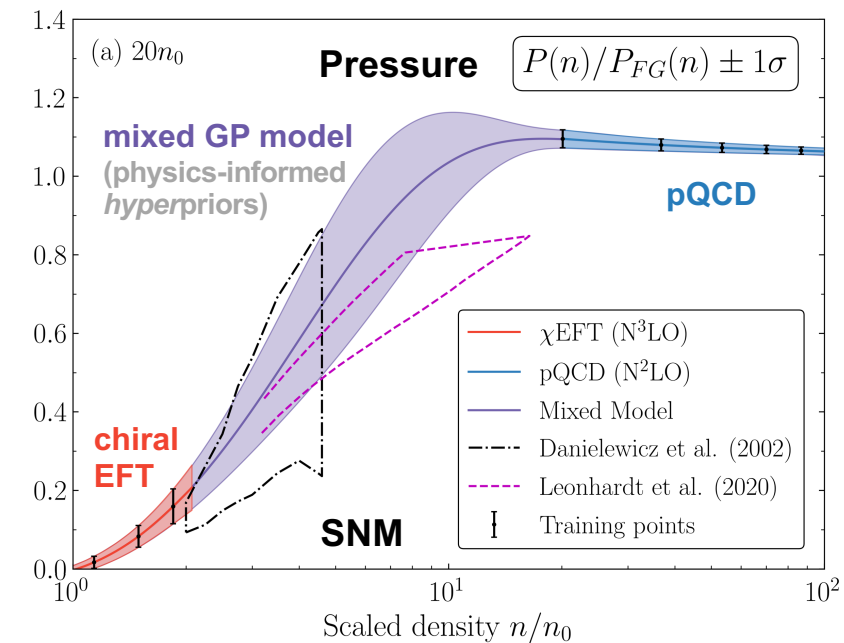
DFT constraints on nuclear saturation



DFT vs EFT: nuclear saturation



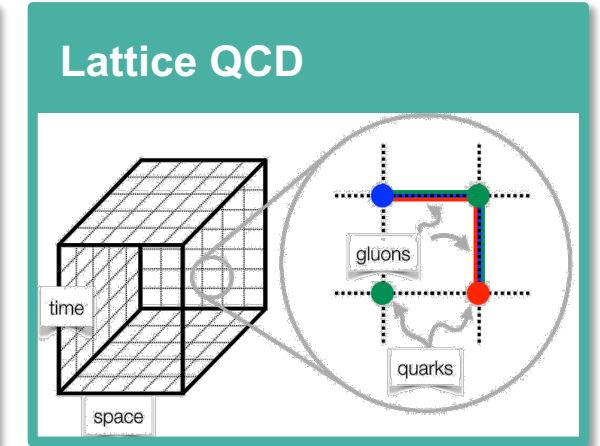
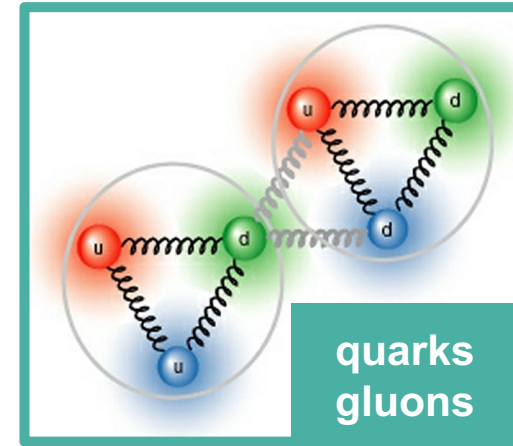
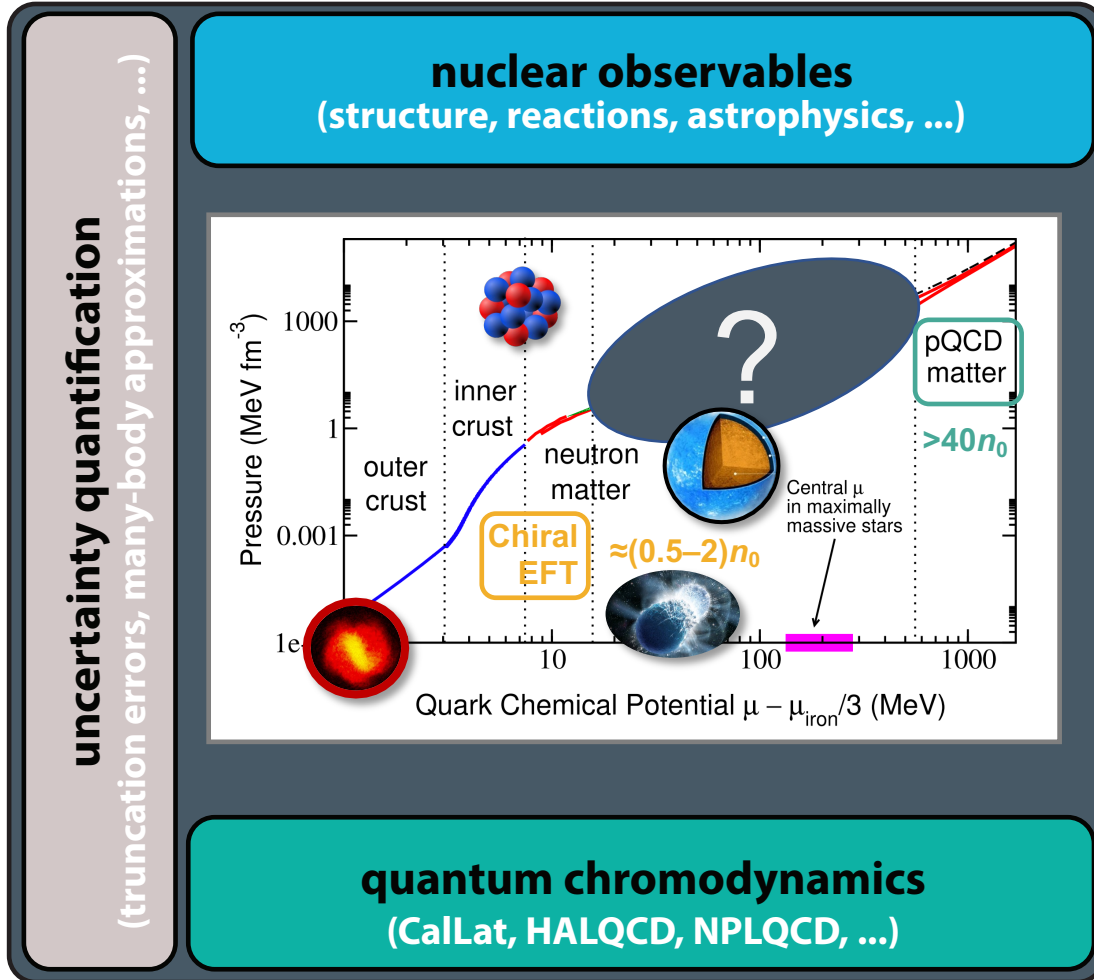
Click to watch **Alexandra's FRIB Theory Seminar** (April 14, 2024)



**Goal: constructing globally predictive, QCD-based EOSs from individual models**

How can we develop *QCD-based* models that are *predictive* across all densities?

Here: nuclear equation of state (EOS)  
Pressure, energy per particle, or sound speed



Fujimoto & Reddy, PRD **109**, 014020

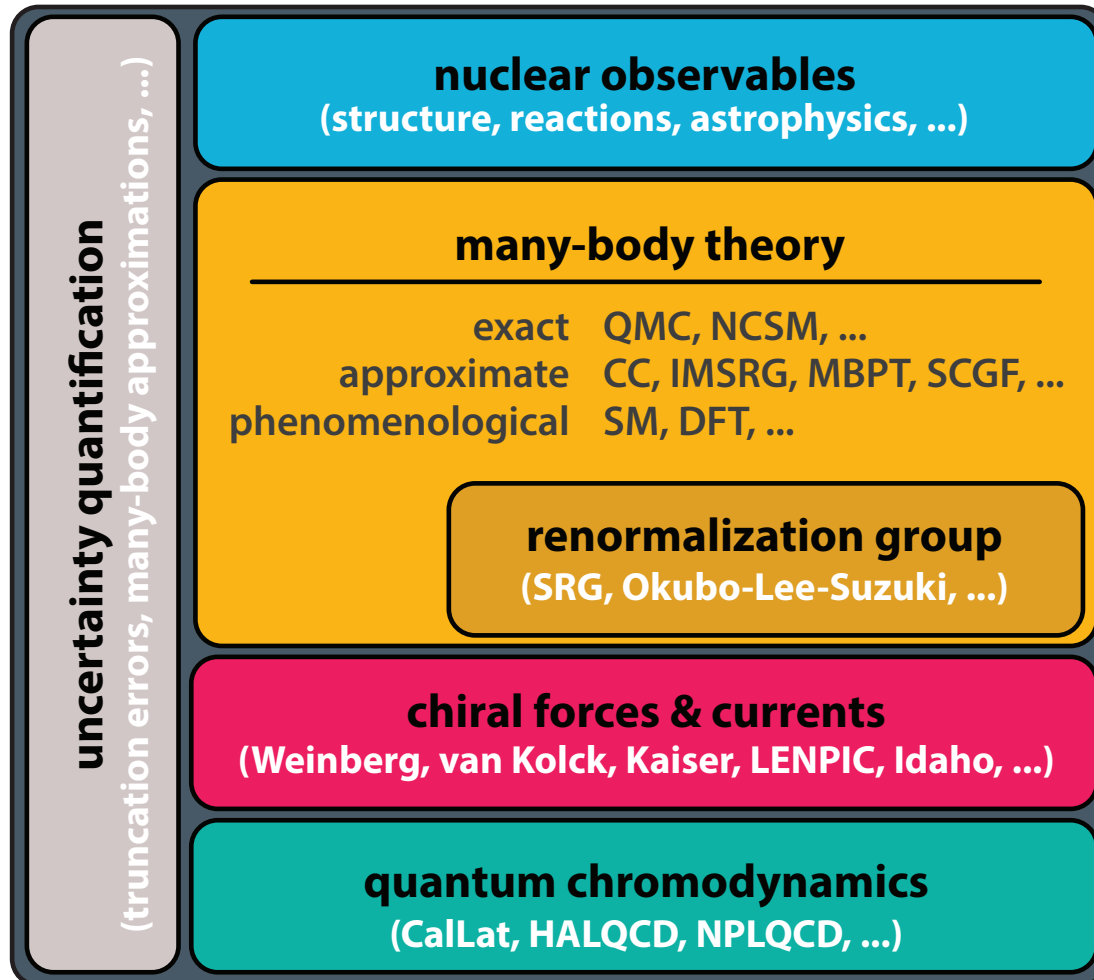
**theory of strong interactions**

QCD is nonperturbative at the low energies  
relevant for nuclear physics (cf. pQCD & LQCD)

CD & Bogner, Few Body Syst. **62**, 109  
e.g., Essick, Tews, Landry, Reddy, Holz, PRC **102**, 055803

CD, Haxton, McElvain, Mereghetti *et al.*, PPNP **121**, 103888

# Low densities: *ab initio* workflow (idealized)



**Here: nuclear equation of state (EOS)**  
Pressure, *energy per particle*, or sound speed

$$\frac{E}{A}(n, \delta, T)$$

baryon density  $n$   
neutron excess  $\delta$   
temperature  $T (= 0)$

## computational framework

solves the (many-body) Schrödinger equation  
requires a nuclear potential as input

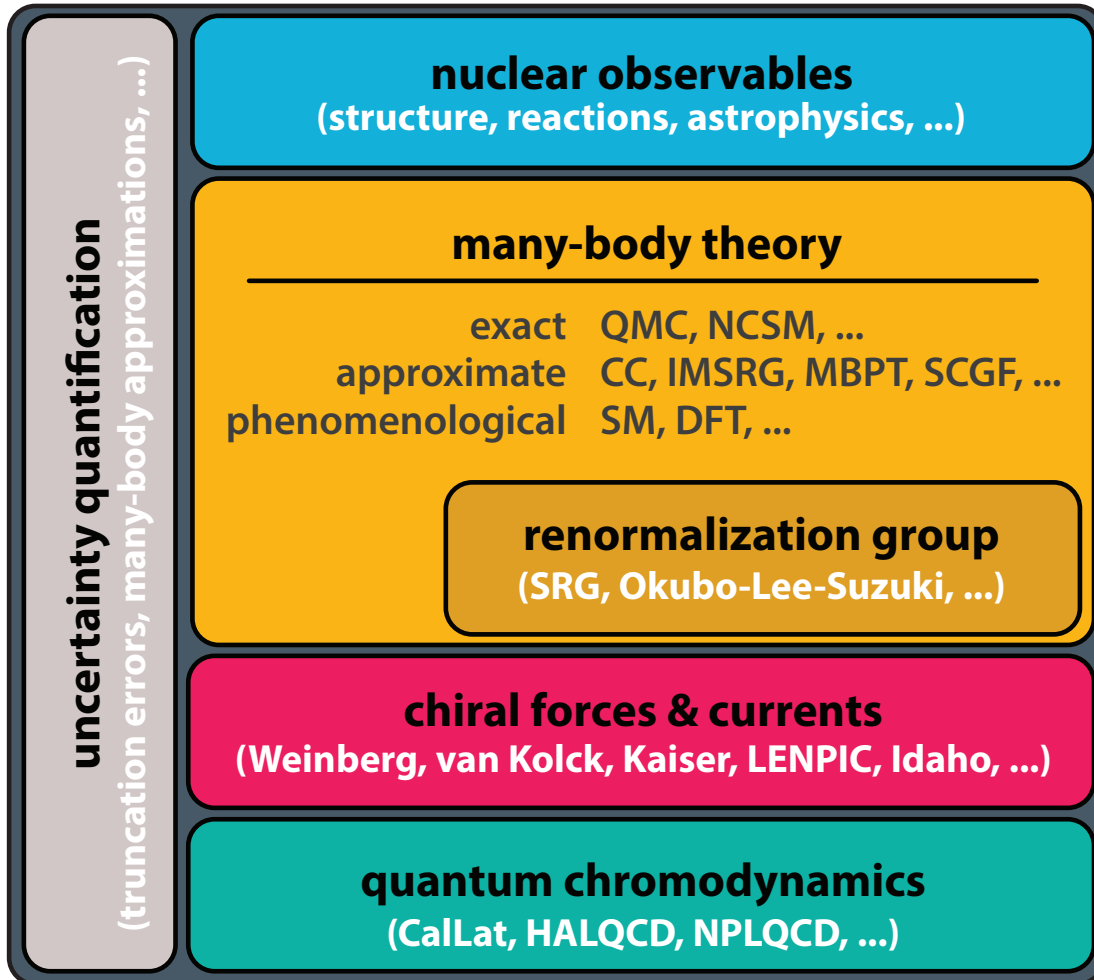
## chiral effective field theory

provides microscopic interactions consistent with  
the symmetries of *low-energy* QCD

## theory of strong interactions

QCD is nonperturbative at the low energies  
relevant for nuclear physics (cf. pQCD & LQCD)

# Low densities: *ab initio* workflow (idealized)



**Here: nuclear equation of state (EOS)**  
Pressure, *energy per particle*, or sound speed

**Exciting developments in *ab initio* many-body theory**      new methods  
*emulators for UQ*  
*new potentials (UQ)*  
e.g., Cook *et al.*, arXiv:2401.11694 (Parametric Matrix Models),  
Somasundaram *et al.*, arXiv:2404.11566

solves the (many-body) Schrodinger equation  
requires a nuclear potential as input

**See also Kang Yu's talk in this session:**  
*Nuclear Matter EOS from the IMSRG*

**theory of strong interactions**  
QCD is nonperturbative at the low energies  
relevant for nuclear physics (cf. pQCD & LQCD)

Interesting new ML/MOR applications: e.g., Fore, Kim *et al.*, PRR 5, 033062

CD & Bogner, Few Body Syst. **62**, 109  
e.g., Essick, Tews, Landry, Reddy, Holz, PRC **102**, 055803

# Modern theory of nuclear forces

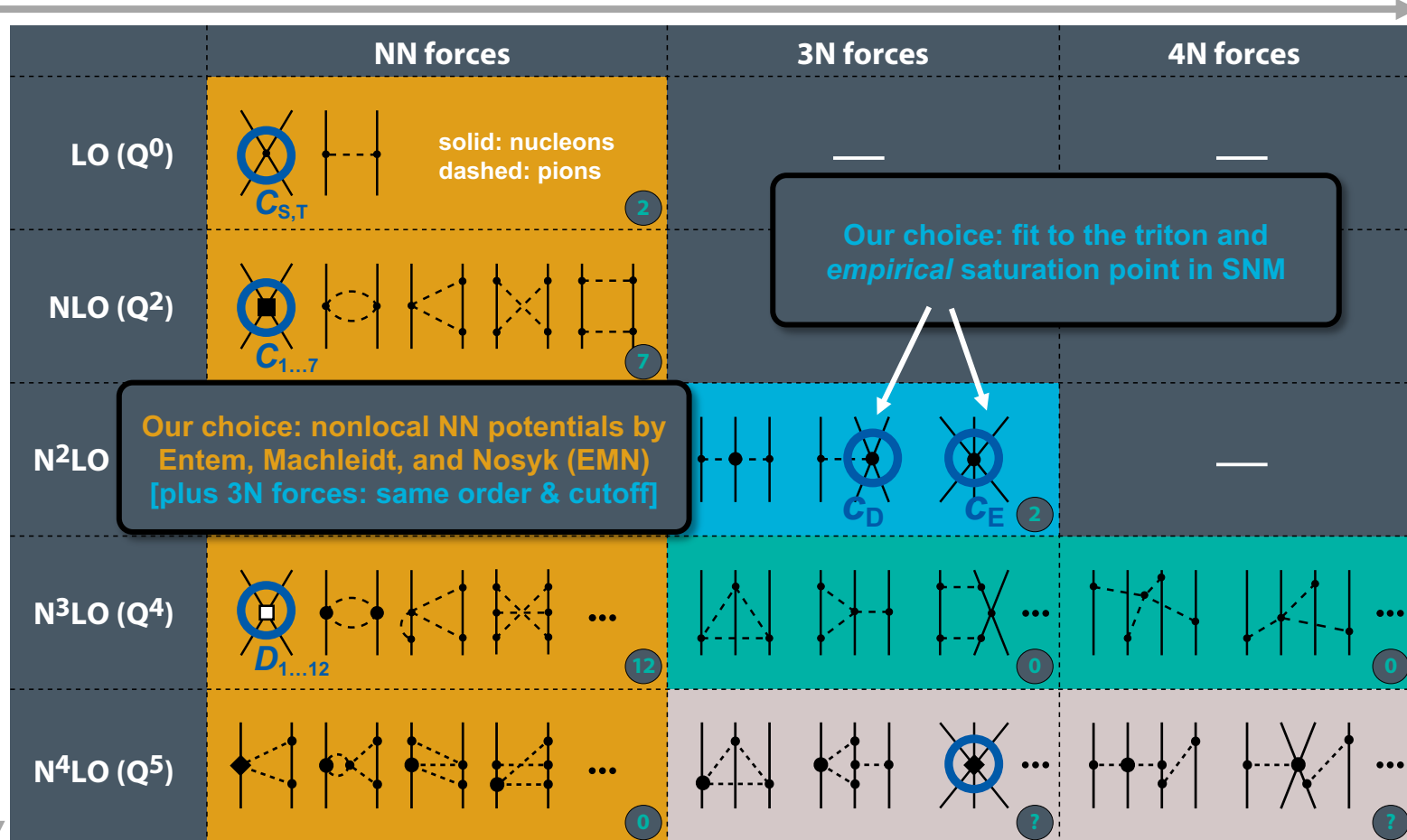


## Hierarchy of chiral nuclear forces up to N<sup>4</sup>LO

Weinberg, van Kolck, Kaplan, Savage, Wise, Epelbaum, Kaiser, Krebs, Machleidt, Meißner, ...

$$Q = \max\left(\frac{p}{\Lambda_b}, \frac{m_\pi}{\Lambda_b}\right) \gtrsim \frac{1}{3}$$

multi-nucleon forces →



### Chiral effective field theory

dominant approach to deriving microscopic interactions consistent with the symmetries of low-energy QCD

degrees of freedom: nucleons & pions

EFT expansion enables uncertainty quantification (EFT truncation errors)

fit the unknown couplings to experimental (or lattice) data

- NN: phase shifts & deuteron
- 3N: binding energies, charge radii, ... (only 2 couplings through N<sup>3</sup>LO)

For recent reviews of delta-full EFT, see, e.g.: Piarulli & Tews, Front. Phys. 7, 245; Piarulli & Schiavilla, Few Body Syst. 62, 10

↑ increasing accuracy

# Correlated EFT truncation error model

Melendez, Furnstahl *et al.*,  
PRC 100, 044001



your  
EFT

$$y_k = y_{\text{ref}} \sum_{n=0}^k c_n Q^n$$

EFT prediction at  $k^{\text{th}}$  order

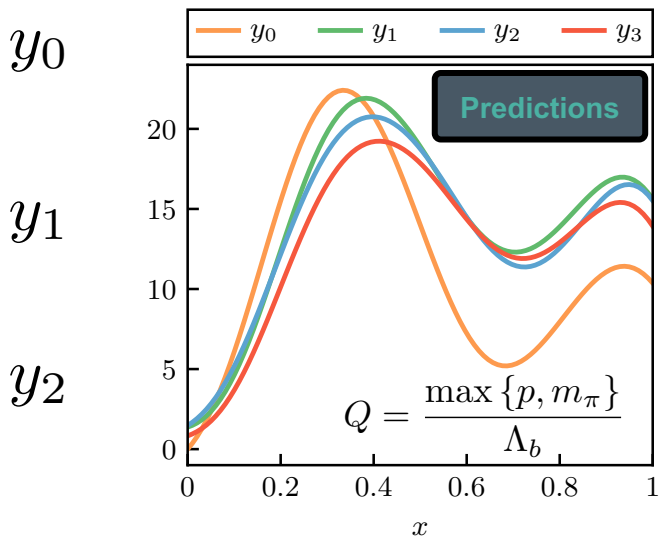
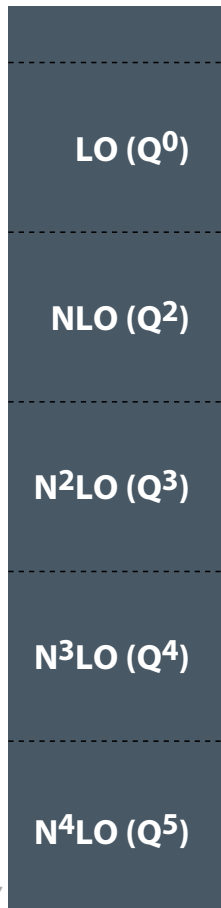
want full prediction:  $y = y_k + \delta y_k$

need to infer theory uncertainty from  
the *computed* EFT orders

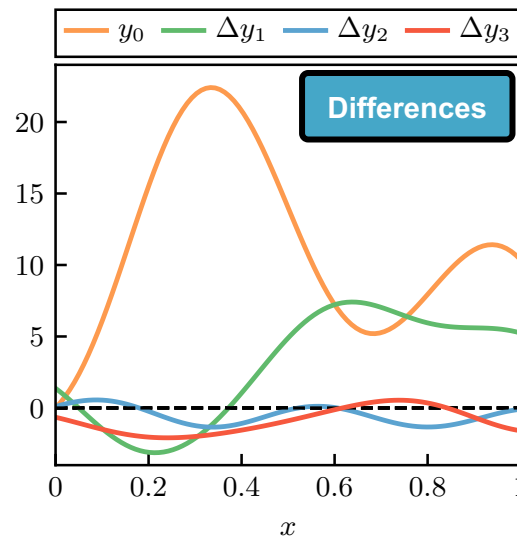
$$\delta y_k = y_{\text{ref}} \sum_{n=k+1}^{\infty} c_n Q^n$$

to-all-orders truncation error

increasing accuracy

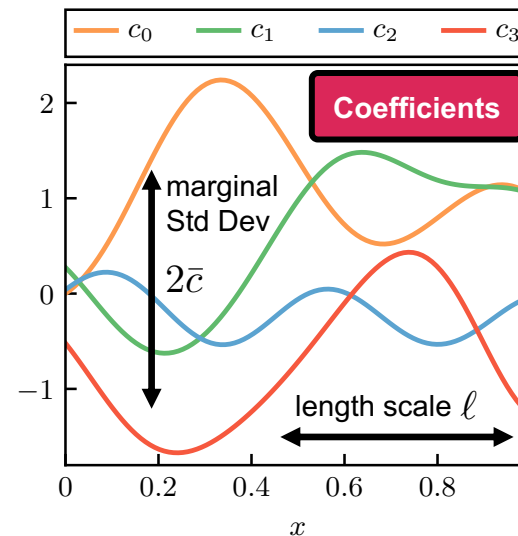


predict observable  $y_k$   
order by order in EFT



$$\Delta y_n = y_n - y_{n-1}$$

differences



model all **coefficients** as  
independent draws from a  
**single Gaussian Process**

$$\mathcal{GP}[0, \bar{c}^2 r(x, x'; \ell)]$$

*natural* coefficients  
(Bayesian) estimation  
of the **kernel** hyper-  
parameters (guided  
by prior information)  
Here: **RBF** (stationary)

Accounts for **correlations** in the observable  $y(x)$  **EFT breakdown scale estimation**  
many applications of this  
truncation error model



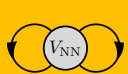
Note:  $c_n$  are not the EFT's LEC

Model checking diagnostics (e.g., Mahalanobis distance)

<https://github.com/buqeye/gsum>

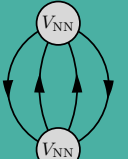
# MBPT in a nutshell

CD, Hebeler, Schwenk, PRL 122, 042501  
 CD, Holt, and Wellenhofer, ARNPP 71, 403  
 Arhuis *et al.*, Comput. Phys. 240, 202




$$\frac{E^{(0)}}{V} = +\frac{1}{2} \sum_{ij} \langle ij | \bar{V}_{NN} | ij \rangle$$

Hartree-Fock



$$\frac{E^{(2)}}{V} = \frac{1}{4} \sum_{ijab} \frac{|\langle ij | \bar{V}_{NN} | ab \rangle|^2}{\varepsilon_i + \varepsilon_j - \varepsilon_a - \varepsilon_b}$$

second order

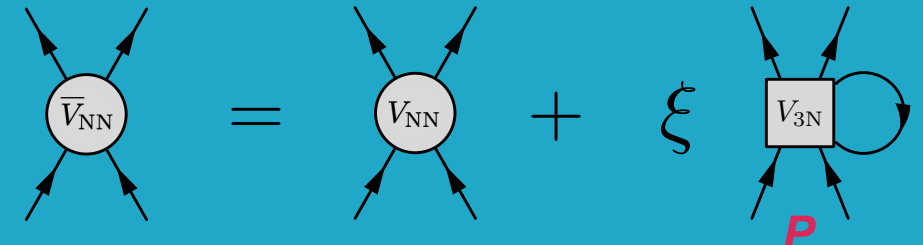


holes:  
 $i, j, k, \dots$

particles:  
 $a, b, c, \dots$

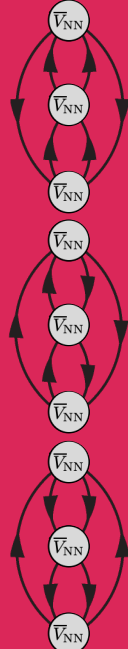
$$|i\rangle = |\mathbf{k}_i \sigma_i \tau_i\rangle$$

effective potential      genuine NN forces      normal-ordered 3N forces



$$\bar{V}_{NN} = V_{NN} + \mathcal{P} V_{3N}$$

Normal ordering 3N forces results in effective two-body potentials



$$\frac{E_{hh}^{(3)}}{V} = +\frac{1}{8} \sum_{ijkl} \frac{\langle ij | \bar{V}_{NN} | ab \rangle \langle kl | \bar{V}_{NN} | ij \rangle \langle ab | \bar{V}_{NN} | kl \rangle}{D_{ijab} D_{klab}}$$

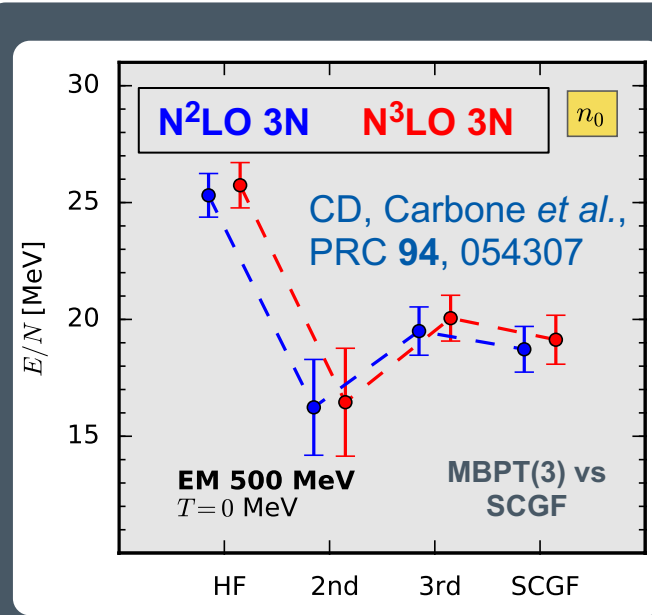
number of MBPT diagrams increases rapidly

$$\frac{E_{ph}^{(3)}}{V} = + \sum_{abcijk} \frac{\langle ij | \bar{V}_{NN} | ab \rangle \langle ak | \bar{V}_{NN} | ic \rangle \langle bc | \bar{V}_{NN} | jk \rangle}{D_{ijab} D_{jkbc}}$$

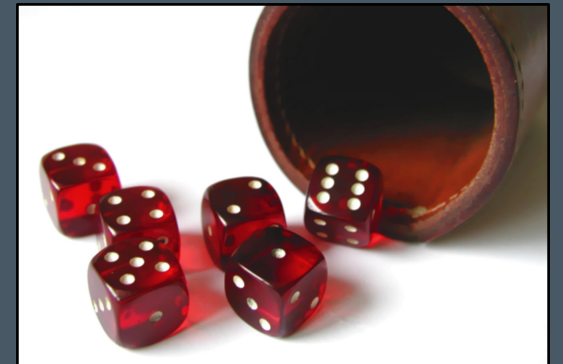
see Coraggio, Holt *et al.*, PRC 89, 044321

$$\frac{E_{pp}^{(3)}}{V} = +\frac{1}{8} \sum_{abcdij} \frac{\langle ij | \bar{V}_{NN} | ab \rangle \langle ab | \bar{V}_{NN} | cd \rangle \langle cd | \bar{V}_{NN} | ij \rangle}{D_{ijab} D_{ijcd}}$$

third order



Controlled evaluation of multi-dimensional momentum integrals:  
MC via improved VEGAS



Automated derivation & evaluation of (thousands of) MBPT diagrams



# Correlated EFT truncation error model (revisited)

your EFT

$$y_k = y_{\text{ref}} \sum_{n=0}^k c_n Q^n$$

EFT prediction at  $k^{\text{th}}$  order

want full prediction:  $y = y_k + \delta y_k$

need to infer theory uncertainty from the *computed* EFT orders

$$\delta y_k = y_{\text{ref}} \sum_{n=k+1}^{\infty} c_n Q^n$$

to-all-orders truncation error

$$Q(k_F) = \frac{k_F}{\Lambda_b}$$

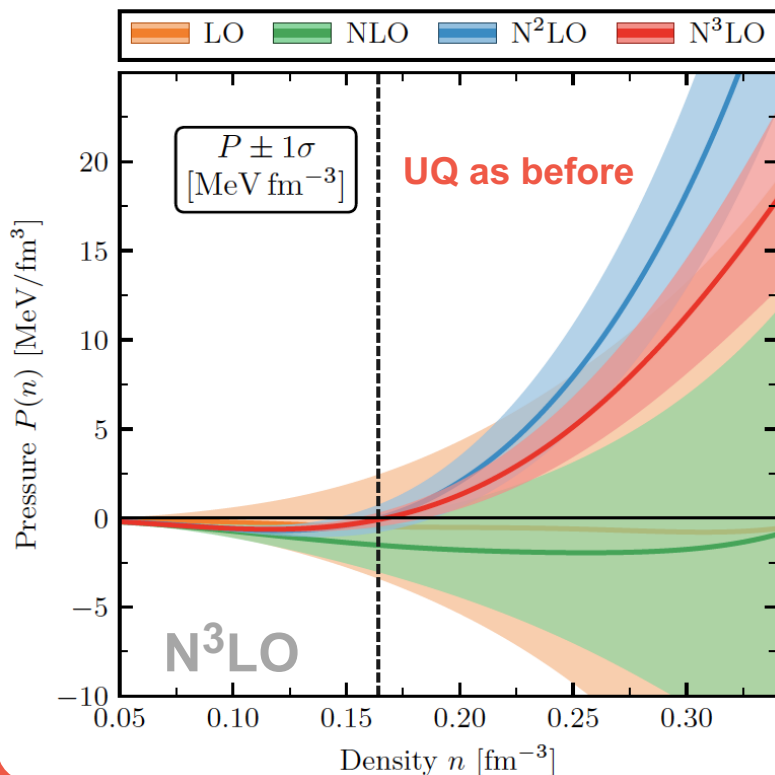
$$y_{\text{ref}}(k_F) = 16 \text{ MeV} \left( \frac{k_F}{k_{F,0}} \right)^2$$

GPs are closed under differentiation:

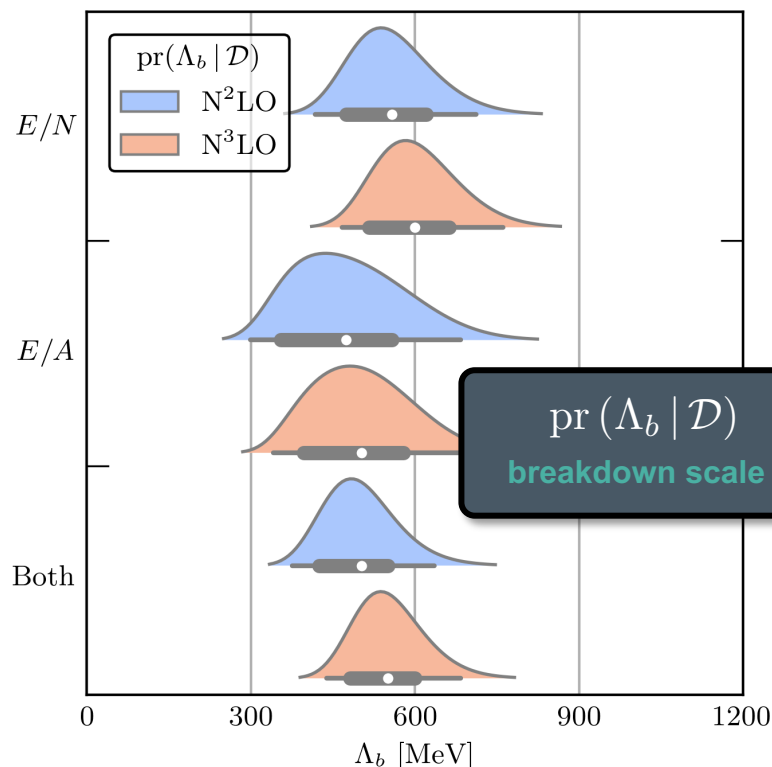
$$P(n) = n^2 \frac{d}{dn} \frac{E}{A}(n)$$

At what *density* does chiral EFT break down, and why?

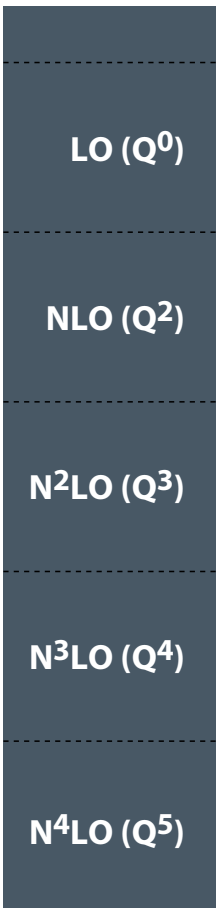
## Chiral EFT



## Bayesian inference of the in-medium breakdown scale



Increasing accuracy



# Correlated EFT truncation errors (for pQCD)

Semposki, CD, Furnstahl,  
Melendez, and Phillips,  
arXiv:2404.06323

your  
EFT

$$y_k = y_{\text{ref}} \sum_{n=0}^k c_n Q^n$$

EFT prediction at  $k^{\text{th}}$  order

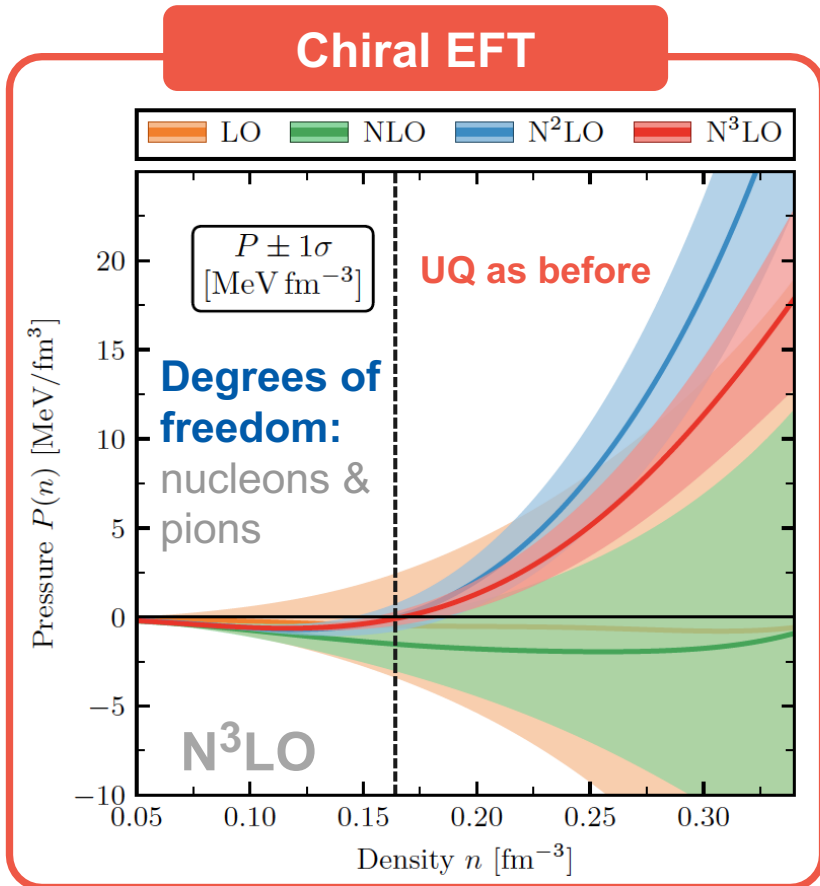
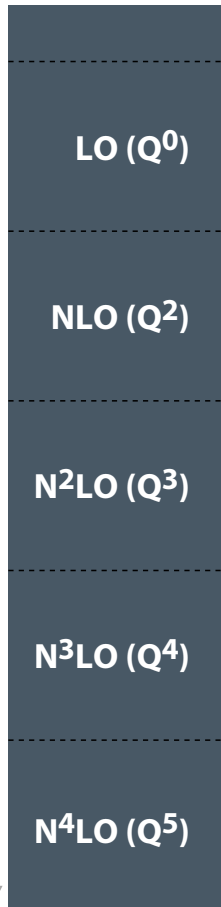
want full prediction:  $y = y_k + \delta y_k$

need to infer theory uncertainty from  
the *computed* EFT orders

$$\delta y_k = y_{\text{ref}} \sum_{n=k+1}^{\infty} c_n Q^n$$

to-all-orders truncation error

Increasing accuracy



**pQCD**

**Perturbative expansion in  
the strong coupling constant**  
(at the two-loop level)

$$\alpha_s(\bar{\Lambda}) = \frac{4\pi}{\beta_0 L} \left[ 1 - \frac{2\beta_1 \ln L}{\beta_0^2 L} \right]$$

$$L = \ln(\bar{\Lambda}^2 / \Lambda_{MS}^2), \quad \bar{\Lambda} = 2X\mu,$$

$$\frac{P(\mu)}{P_{FG}(\mu)} \approx 1 + \alpha_s(\bar{\Lambda}) \left[ \frac{a_{2,2}}{2\mu} + a_{2,3} \right] + \mathcal{O}(\alpha_s^3),$$

See also talks by **Aleksi Vuorinen**  
and **Tyler Gorda** next Monday

**Degrees of freedom:** massless quarks  
(*up* & *down*, with equal chemical potential  $\mu$ ) and **gluons**

cf. Bayesian analysis using the MiHO framework: Gorda *et al.*, PRL **131**, 181902; JHEP **06**, 002 (see also more recent work)

# Correlated EFT truncation errors (for pQCD)

Semposki, CD, Furnstahl,  
Melendez, and Phillips,  
arXiv:2404.06323

your  
EFT

$$y_k = y_{\text{ref}} \sum_{n=0}^k c_n Q^n$$

EFT prediction at  $k^{\text{th}}$  order

want full prediction:  $y = y_k + \delta y_k$

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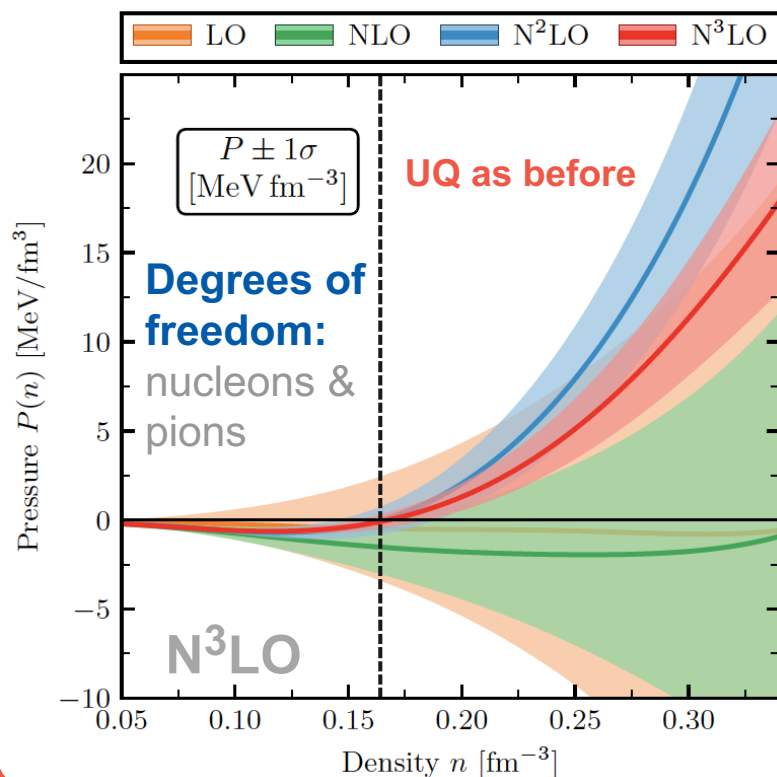
$$\delta y_k = y_{\text{ref}} \sum_{n=k+1}^{\infty} c_n Q^n$$

to-all-orders truncation error

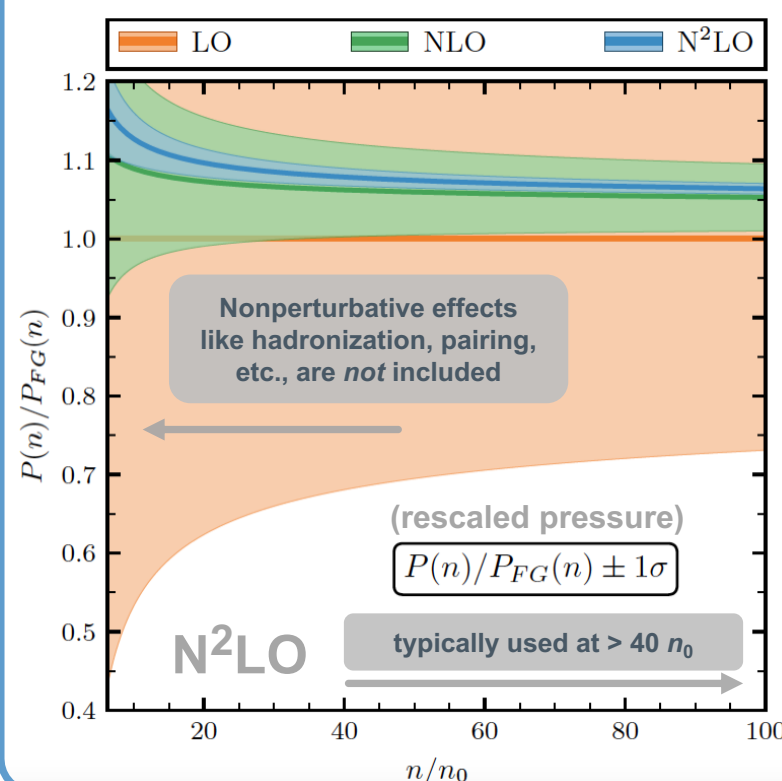
Increasing accuracy



Chiral EFT



pQCD



Truncation error estimation:

$$Q = \frac{N_f}{\pi} \alpha_s (\bar{\Lambda}(\mu_{FG}))$$

$$y_{\text{ref}} = P_{FG}(n) \text{ (two-loop level)}$$

Kohn-Luttinger-Ward inversion:

$$P(\mu) \rightarrow P(n)$$

(consistent up to the desired order in pQCD)

pQCD prediction:

$$P(\mu)$$

workflow

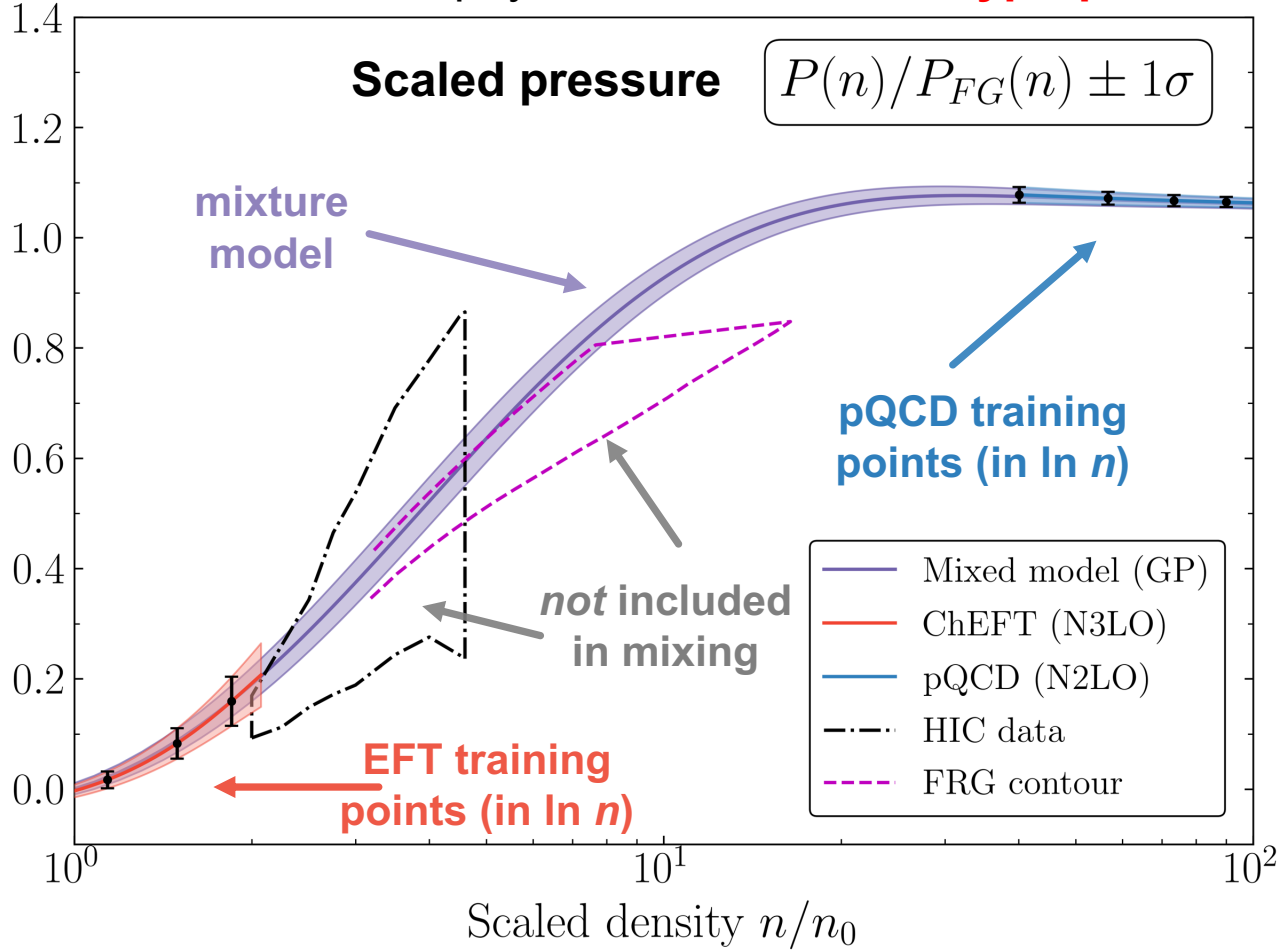
# Curvevise mixing of random variables



Semposki, CD, Furnstahl,  
Melendez, and Phillips,  
arXiv:2404.06323

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One can include physics constraints via **hyperpriors**



Here, only 2 models:

- 1: Chiral EFT
- 2: pQCD



random variable (at a given density) corresponding to the predictions of model  $i$

$$Y^{(i)} = F + \delta Y^{(i)}$$

(assumes common mean)

**QCD, with prior**

$$\sim \text{GP}[0, \kappa_f(x, x')]$$

kernel choice: here, RBF  
(hyperparameters estimated from data)

$$i \in [1, M]$$

**BUQEYE truncation error**

$$\sim \text{GP}[0, \kappa_y^{(i)}(x, x')]$$

full, block-diagonal covariance matrix

**We found for the BMM:**

$$\vec{F} \mid \vec{y}, K_y, K_f \sim \mathcal{N}[\vec{\mu}, \Sigma]$$

**Assumptions (not necessarily satisfied, validation needed):**

- $F$  is smooth, precluding *discontinuous* phase transitions
- stationarity: persistence in size & length scale of EOS's variability

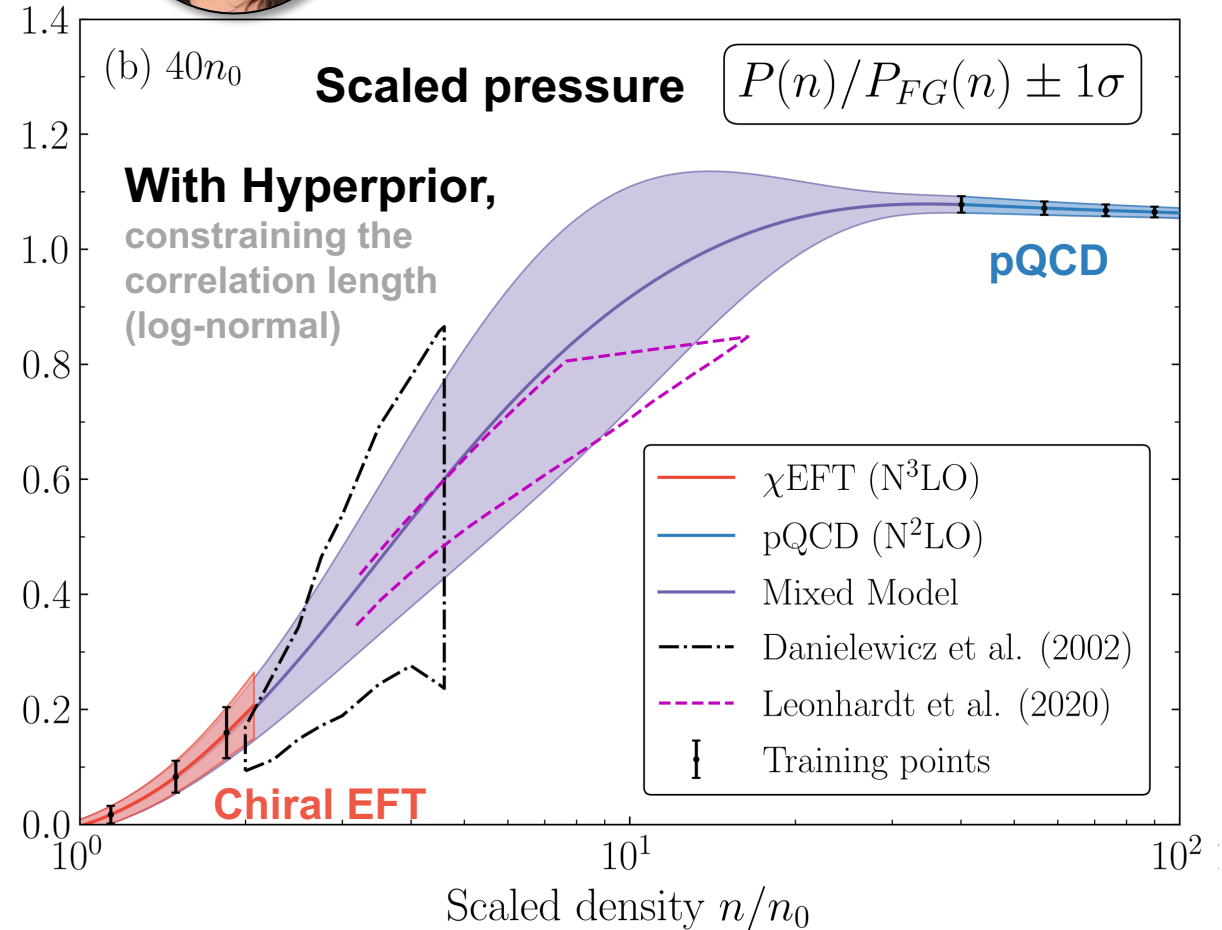
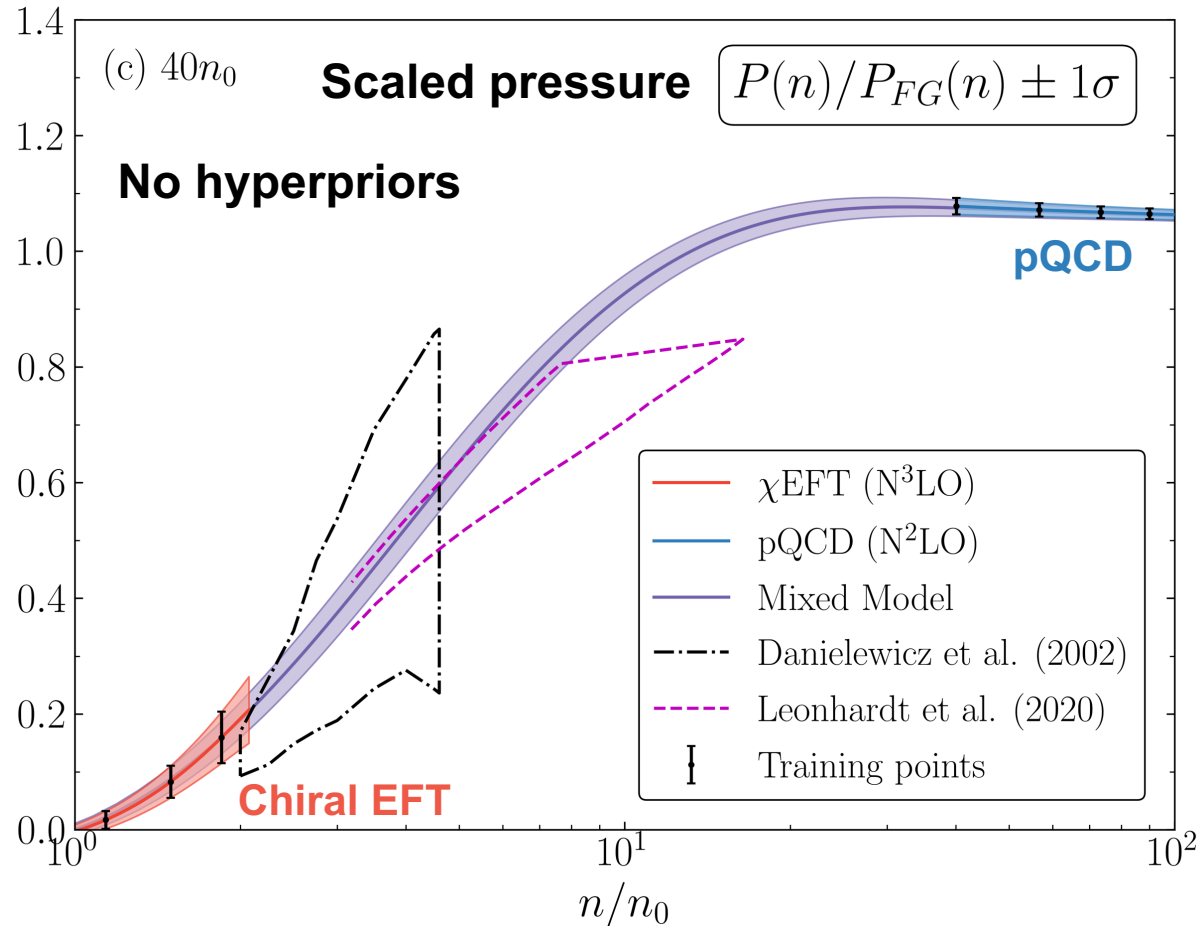
$$\vec{\mu} \equiv \Sigma B_t^T K_y^{-1} \vec{y}, \quad \Sigma \equiv (K_f^{-1} + B_t^T K_y^{-1} B_t)^{-1}$$

# Sensitivity on physics-informed priors



Semposki, CD, Furnstahl,  
Melendez, and Phillips,  
arXiv:2404.06323

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Inferred **long correlation lengths** render uncertainty on the mixed EOS very small due, even smaller than each model

**Unrealistic, large impact of pQCD on chiral EFT region**

We placed a *hyperprior* on the correlation length to **enforce small covariances between EFT & pQCD**

Smaller length scales result in larger uncertainty bands

# Training points in pQCD region



Semposki, CD, Furnstahl,  
Melendez, and Phillips,  
arXiv:2404.06323

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The pQCD uncertainties do *not* account for **nonperturbative effects** (such as hadronization and pairing)

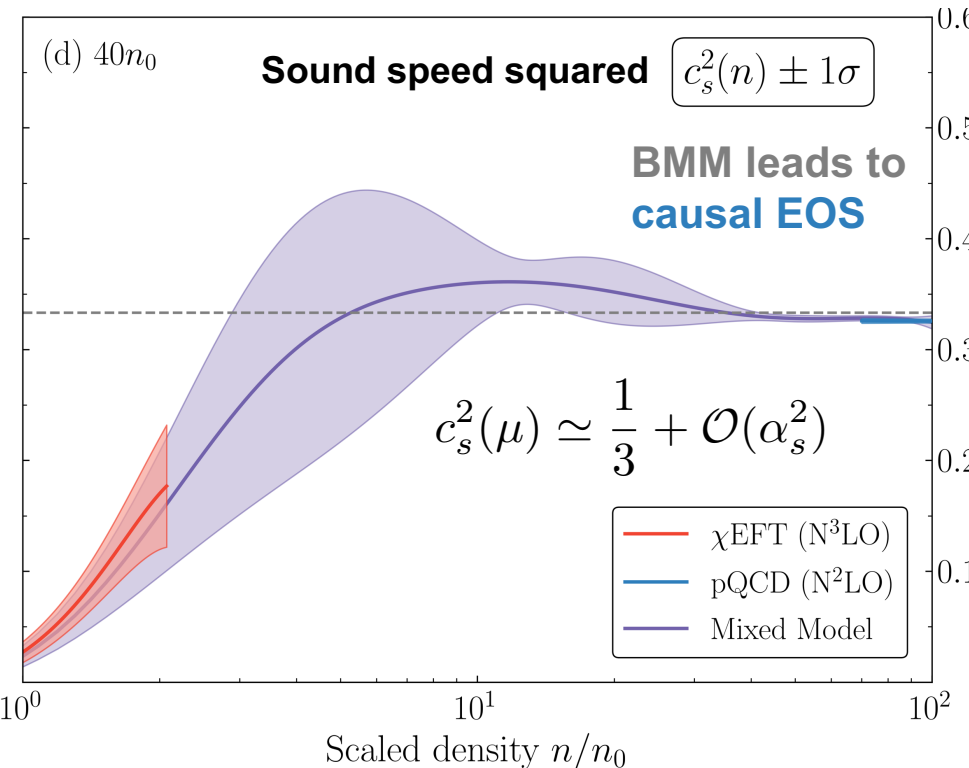
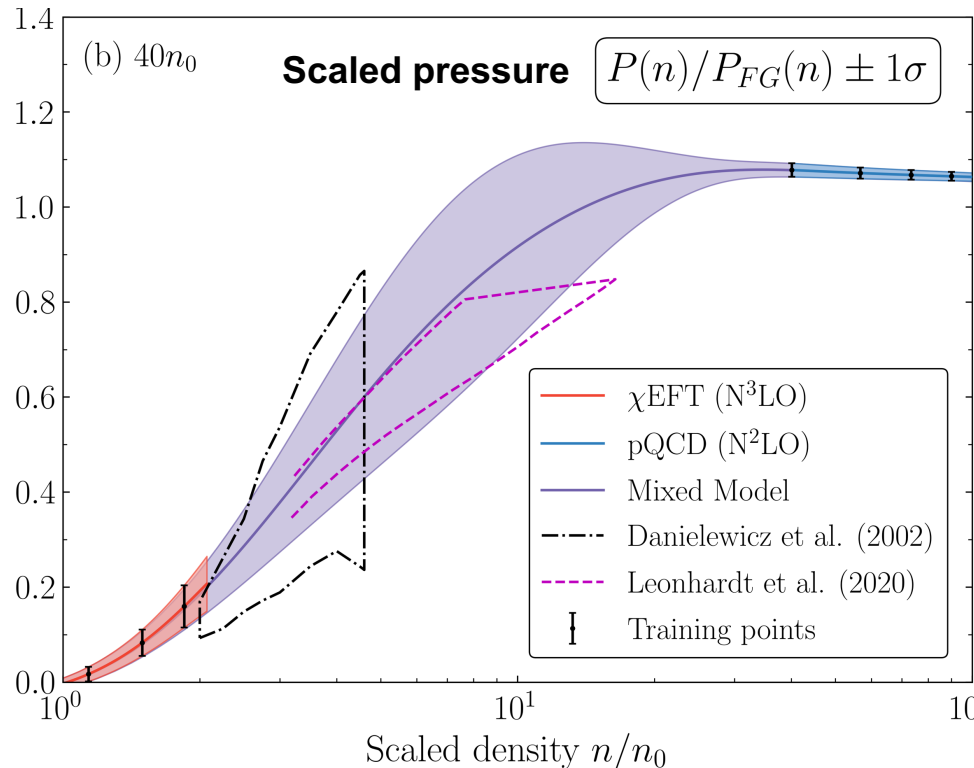
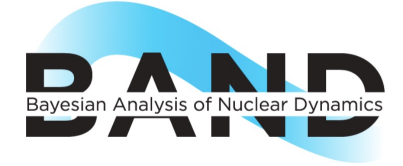
These effects become more important as the density is lowered

What is the lowest density for including pQCD constraints?

$$n \geq 40n_0$$



Open-Source Software:  
**Taweret** (BAND framework)



The **mixed model** approaches the **conformal limit** from below, as expected

$$\longrightarrow c_s^2 = \frac{1}{3}$$

pQCD:  
two massless  
quark flavors

$$c_s^2(n) = \frac{\partial P}{\partial \varepsilon}$$

The FRG & HIC constraints are only shown as references as they do not provide a C.L.

# Training points in pQCD region



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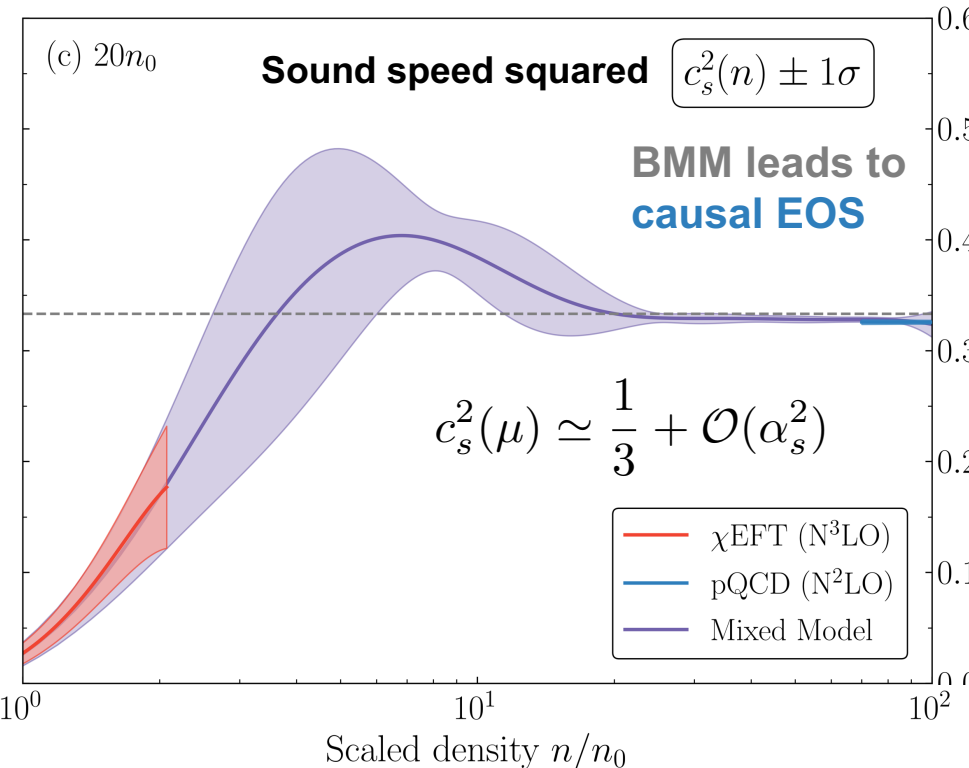
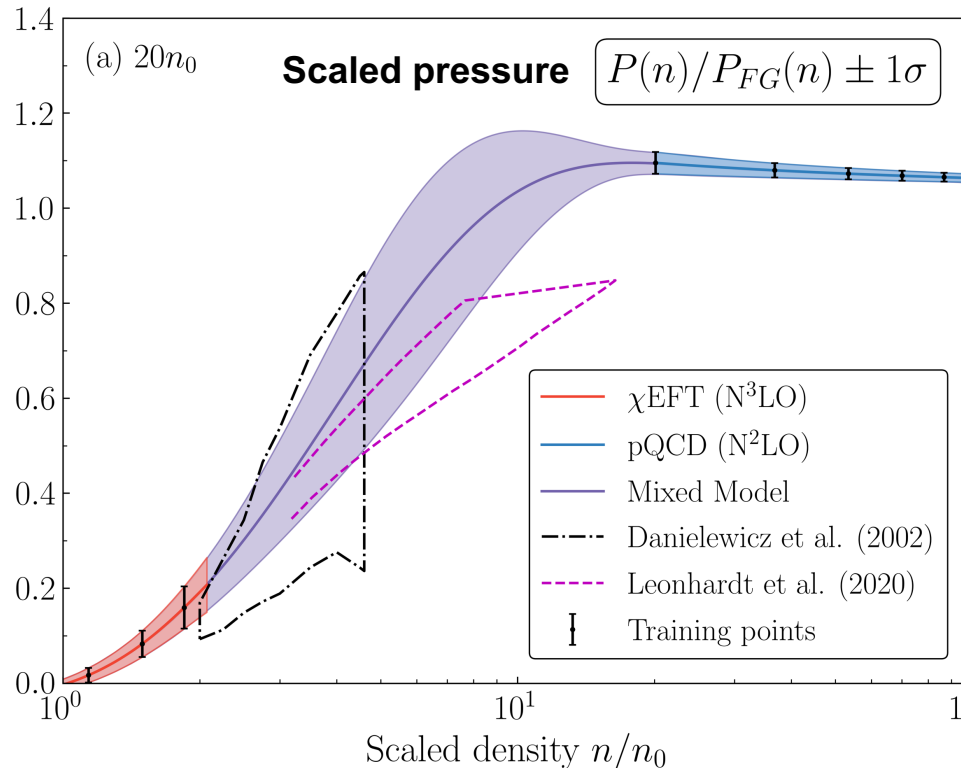
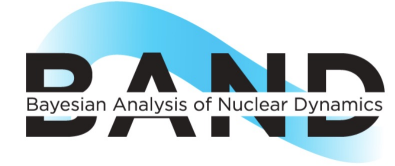
These effects become more important as the density is lowered

What is the lowest density for including pQCD constraints?

$$n \geq 20n_0$$



Open-Source Software:  
**Taweret** (BAND framework)



The **mixed model** approaches the **conformal limit** from below, as expected

$$\longrightarrow c_s^2 = \frac{1}{3}$$

pQCD:  
two massless  
quark flavors

$$c_s^2(n) = \frac{\partial P}{\partial \varepsilon}$$

The FRG & HIC constraints are only shown as references as they do *not* provide a C.L.

- 1 Chiral EFT enables **microscopic calculations** of nuclei and infinite matter at  $n \lesssim 2n_0$  (and finite temperature) **with quantified uncertainties**
- 2 BMM combines multiple predictive models in different regions into one **overall predictive composite model**. Not limited to the EOS, MBPT, or EFT!
- 3 Promising method for constructing **globally predictive, QCD-based EOSs with rigorous UQ** to study the structure & evolution of neutron stars  
*Uncertainties in the mixed region depend significantly on **physics-informed priors**. Guidance needed.*
- 4 Requires **extension to neutron star matter** (and finite temperatures) and inclusion of recent **neutron star observations & nuclear experiments**



$a_1$  Prior  
Posterior  
True value



Many thanks to: **R. Furnstahl** **J. Melendez** **D. R. Phillips** **A. Semposki**