

# Statistical distributions of compact remnants from supernovae and the nuclear equation of state

**Carla Fröhlich**

North Carolina State University



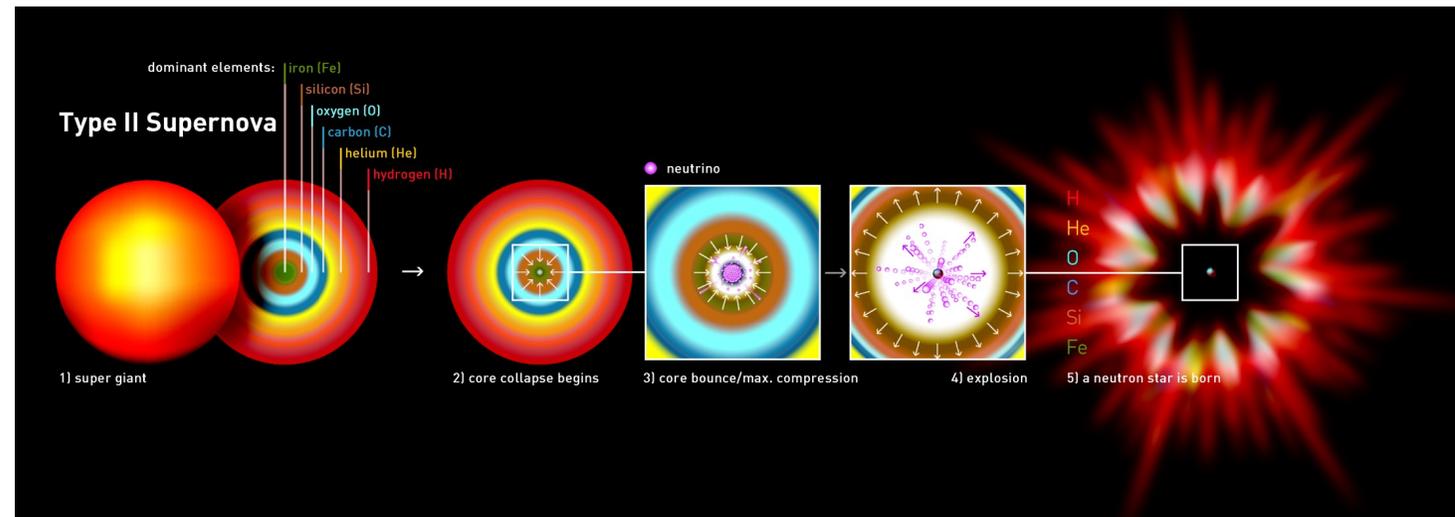
Statistical distributions  
of compact remnants from **supernovae**  
and  
the nuclear equation of state

# Core-Collapse Supernovae (CCSNe)

- Massive stars ( $> \sim 8-10 M_{\text{sun}}$ ) at the end of their lives
  - After Si-burning
  - Onset of collapse (negative velocities in the core)

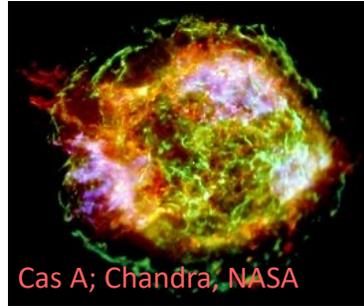
Janka (2012)

- Phases of a CCSN:
  - Collapse
  - Core-bounce
  - Prompt shock
  - Shock stall
  - Revival of shock / no revival
  - Explosion / no explosion



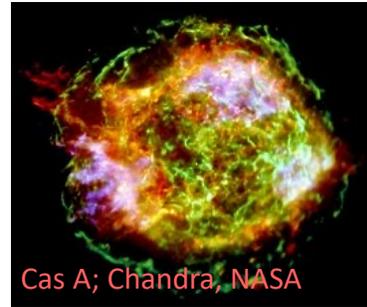
# Core-collapse supernova simulations

- Multi-dimensional problem
- Multi-physics problem:
  - General relativity
  - Nuclear physics of dense matter
  - Neutrino transport  
(trapped, diffusive, free-streaming regimes)
- Multi-scale problem:
  - shock formation at  $\sim 200$  km vs entire star  $10^8$  km
  - collapse and shock formation  $\sim 1$  s vs shock breakout  $\sim 1$  day



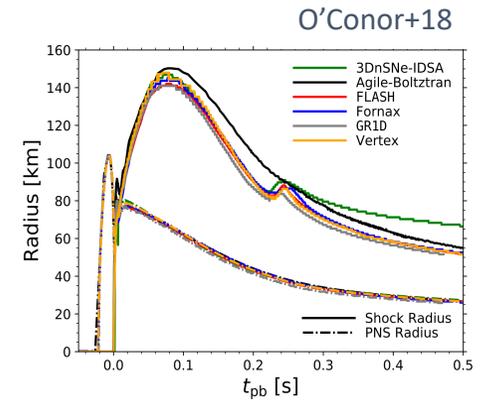
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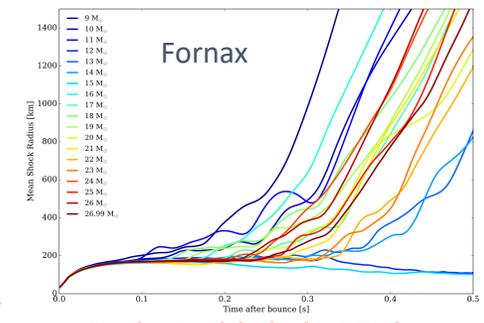
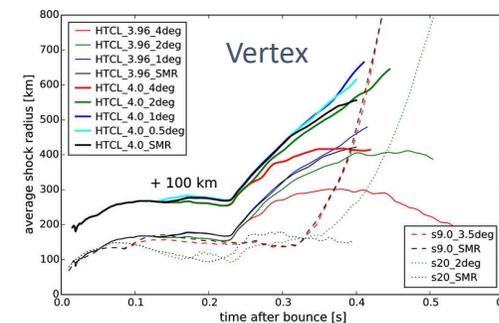
Simulation Status:

1D: in general no self-consistent explosions  
 $\sim 10$  CPUh/model



2D: models have converged

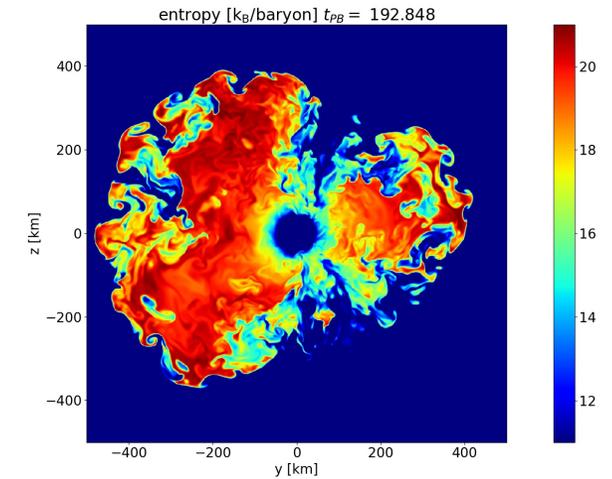
3D: mixed results  
 $\sim$  Mio CPUh/model



Carla Frohlich (NCSU)

# The paths forward

- Self-consistent 3D simulations
  - The ultimate goal
  - Computationally expensive  $\rightarrow$  can do  $O(10)$
- Effective models
  - Simplify part of the problem, but have free parameters
  - Physically reliable
  - Computationally efficient  $\rightarrow$  can do  $O(1000)$
- The two paths are complimentary (3D-1D-3D feedback loop)
- Both paths are needed for current open science questions



# Effective CCSN Models

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- Parametrize a multi-dimensional aspect in 1D simulations
  - Mixing above the PNS, enhanced neutrino heating, etc
- Calibrate parametrization, then apply to many models
  - Eg a suitable model should reproduce observables of SN1987A
  - Predictive within the framework

- **PUSH**: Parametrized neutrino heating

Perego+15, Ebinger+19, Curtis+19, Ebinger+20, Ghosh+23

- **PHOT-B**: Parametrized neutrino heating

Ugliano+12, Ertl+15, Sukhbold+16

- **STIR**: Parametrized mixing above PNS

Couch+20

O'Connor+13; Mueller+15; Pejcha15; Fryer+12,22; ...

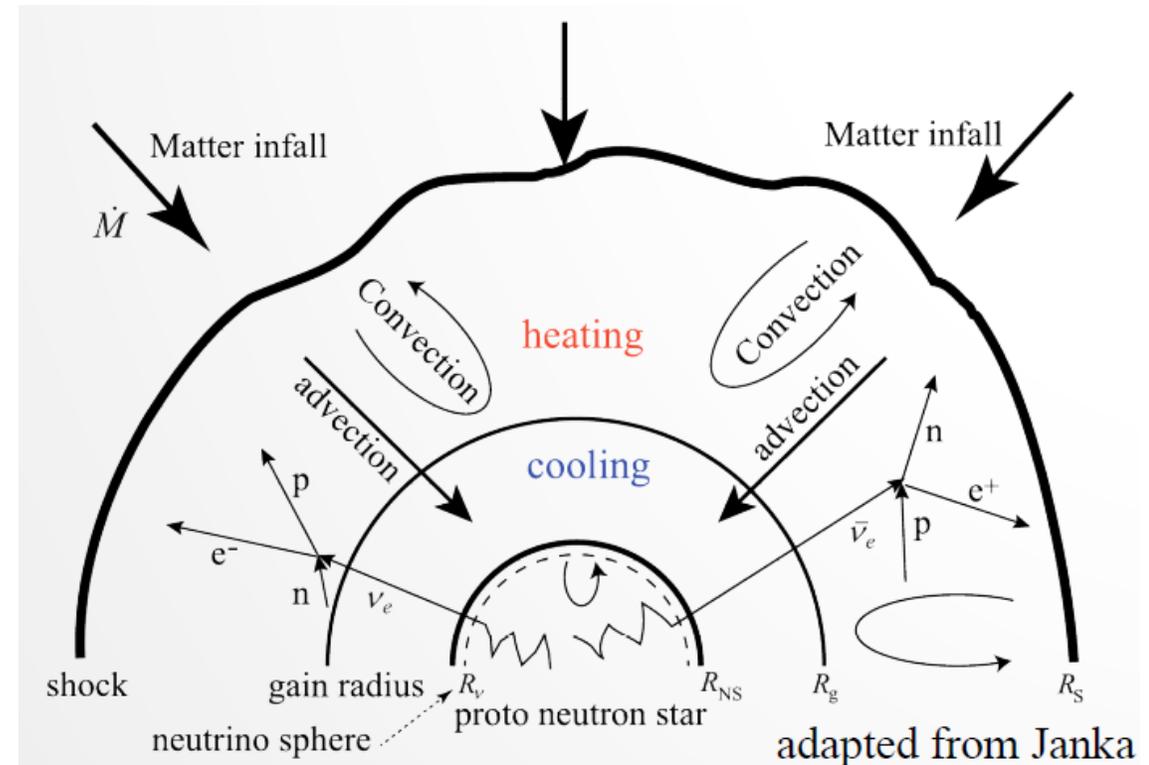
# PUSH: An effective CCSN Model

- Neutrino-driven (convection-aided) mechanism
  - Neutrinos are emitted from hot PNS, deposit energy behind the shock
  - Material behind the shock is unstable to convection → enhanced neutrino heating
- Additional (artificial) heating term:

$$\frac{dE_{\text{tot}}}{dt} = \frac{dE_{\nu_e + \bar{\nu}_e}}{dt} + \frac{dE_{\text{push}}}{dt}$$

Standard heating  
from electron  
(anti-) neutrinos

Additional  
heating in PUSH



# Simulation Setup: PUSH

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- General relativistic hydrodynamics: Agile

Liebendoerfer+02

Simulation time: up to 15sec (typically ~8sec)

- Neutrino transport:
  - IDSA and advanced spectral leakage (ASL)

Lieberdoerfer+09; Perego+16

Electron fraction is evolved during collapse and explosion

- Nuclear EOS: 6 different nuclear EOSs
  - DD2, SFHo, SFHx, BHB $\lambda\phi$ , TM1, NL3

Hempel+02; Typel+10

Mass cut emerges from the simulation  
→ ejecta and explosion energy are not independent “knobs to turn”

→ Predictive (within the framework) for outcome (NS or BH), explosion energy, etc

# Statistical distributions of compact remnants from supernovae and the nuclear equation of state

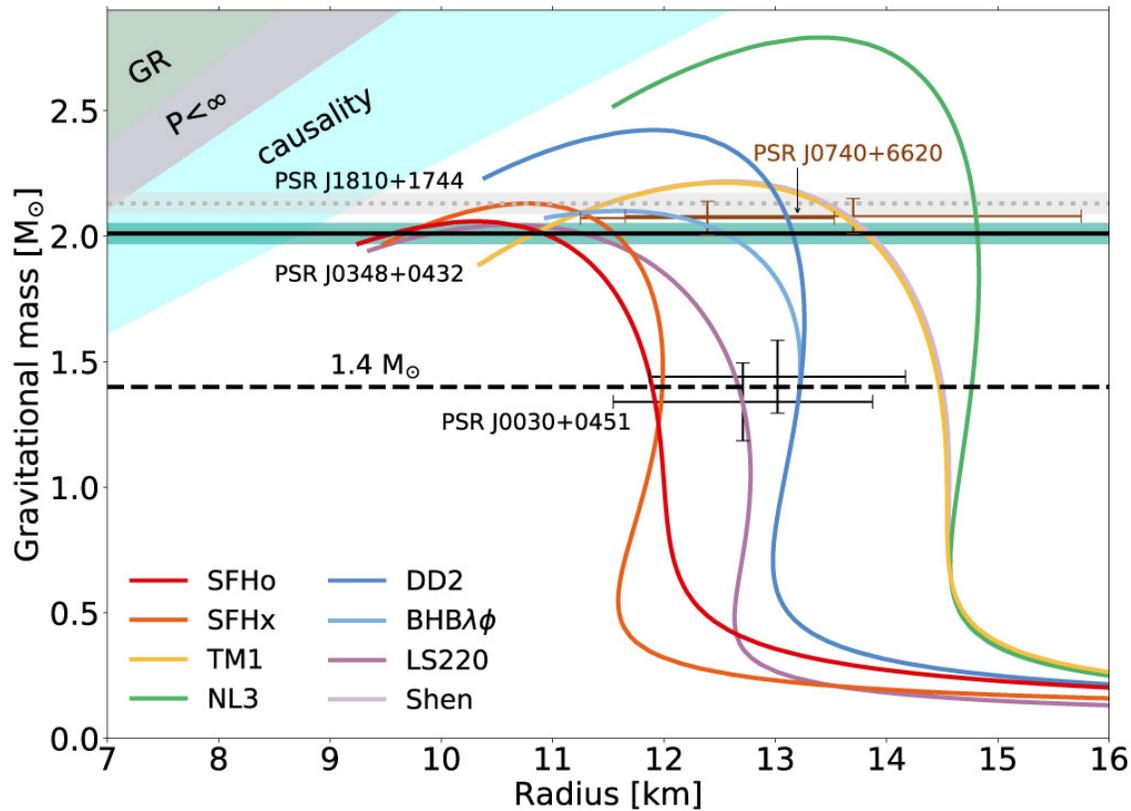
(as relevant for supernova simulations)

# Supernova nuclear EOS

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- Nuclear physics input to astrophysical simulations
  - Thermodynamic quantities
  - Nuclear composition
- Challenges:
  - Finite temperature:  $T = 0 - 100 \text{ MeV}$
  - No weak equilibrium:  $Y_e = 0 - 0.6$
  - Wide density range:  $\rho = 10^4 - 10^{15} \text{ g/cm}^3$
  - **In tabular form:**  $\sim 1$  million points in  $(T, Y_e, \rho)$

# Tabulated EOS



EOS	$K$ (MeV)	$m_n^*/m_n$	$m_p^*/m_p$	$M_{\max}$ ( $M_{\odot}$ )	$R_{1.4M_{\odot}}$ (km)
DD2	242.7	0.5628	0.5622	2.42	13.2
SFHo	245.4	0.7609	0.7606	2.06	11.9
SFHx	238.8	0.7179	0.7174	2.13	12.0
BHB $\lambda\phi$	242.7	0.5628	0.5622	2.10	13.2
TM1	281.6	0.6343	0.6338	2.21	14.5
NL3	271.5	0.5954	0.5949	2.79	14.8

# Statistical distributions of compact remnants from supernovae and the nuclear equation of state

Or, what can we do with >1500 supernova simulations?

# How to compare simulated and observed data

- Typical setup:

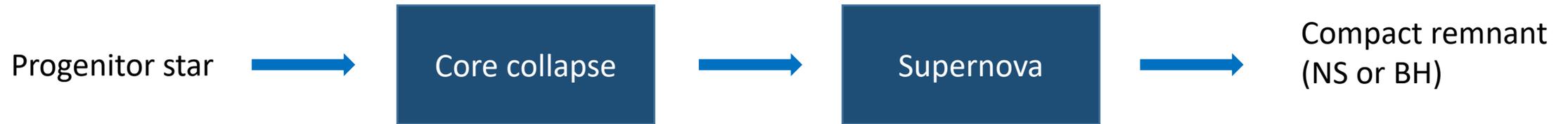
- Collect input-output pairs from Nature  $\{((x)_i, (\tilde{M}_{\text{BH}})_i)\}_{i=1}^n$  x ... initial mass & metallicity

- Calibrate the simulator  $\arg \min_s \frac{1}{n} \sum_{i=1}^n \left\{ (\tilde{M}_{\text{BH}})_i - \eta((x)_i; s) \right\}^2$  S ... EOS
- observed                      simulated
-

# From progenitor to compact remnant

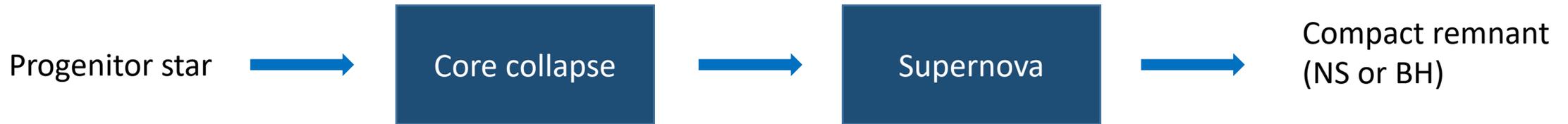
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Simulation:

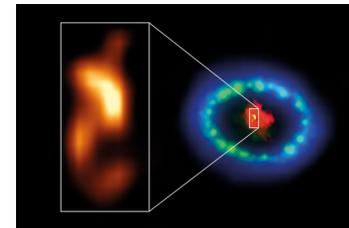
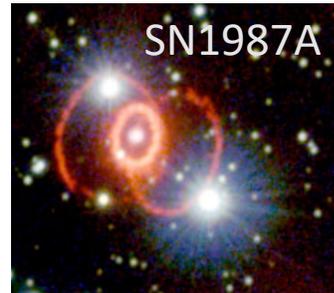
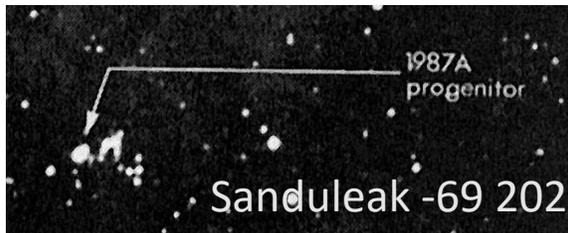


# From progenitor to compact remnant

Simulation:

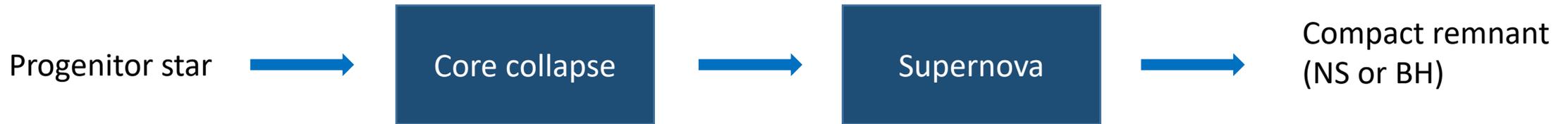


SN1987A:



# From progenitor to compact remnant

Simulation:

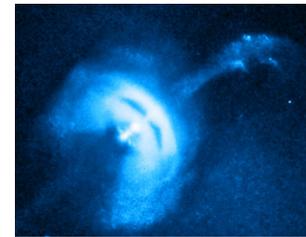
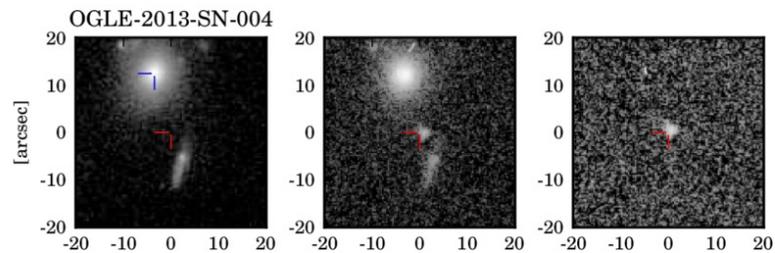


Nature:

Stars

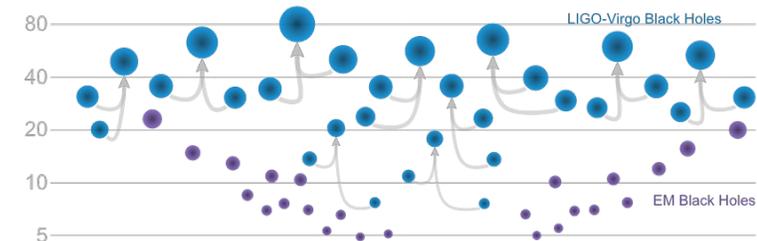


Supernovae



Neutron stars

Black holes



# How to compare simulated and observed data

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- observed                      simulated

- Here:

- We cannot observe such input-output pairs from Nature
- We cannot calibrate on the joint distribution  $p(x, \tilde{M}_{\text{BH}})$   
nor the conditional distribution  $p(\tilde{M}_{\text{BH}} | x)$
- Instead obtain **marginal distributions**:  $p(x)$  and  $p(\tilde{M}_{\text{BH}})$

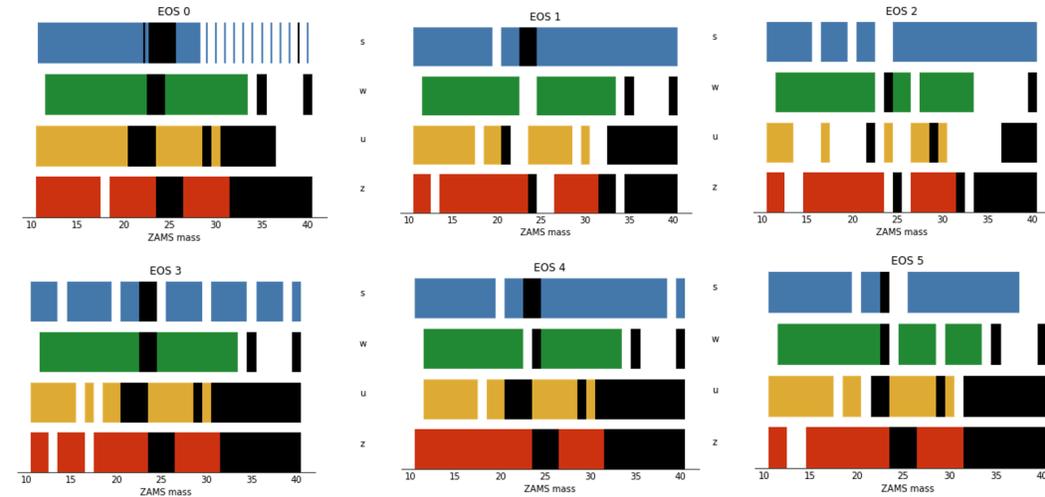
# Simulated data

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- Simulations of core-collapse supernovae using PUSH
- Each simulation predicts as output as remnant mass  $\eta(x,s)$ 
  - Either a neutron star (NS) of a given mass
  - Or a black hole (BH) of a given mass
- Simulation is a mapping  $\zeta(x): x \rightarrow \eta(x,s)$
- Assume: mapping mimics physical reality for ideal EOS  $s_0$   
$$\zeta(x) \approx \eta(x,s_0)$$

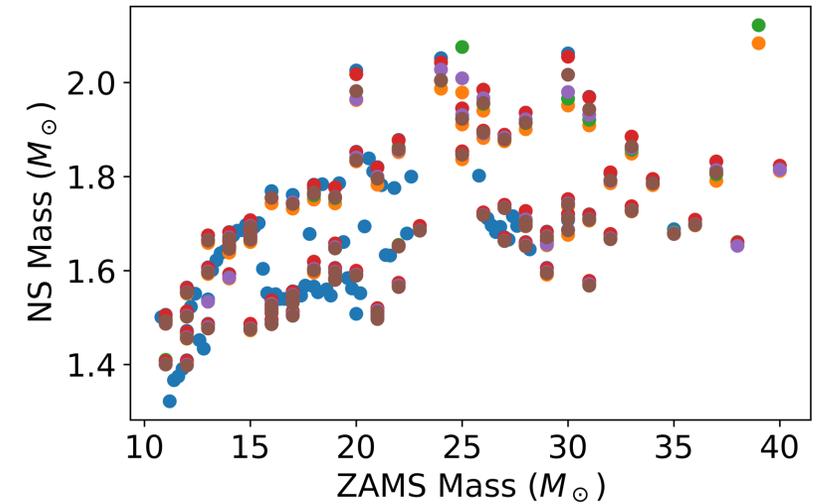
# Simulated data

- Binary outcome: explosion / no explosion
  - Explosion  $\rightarrow$  NSs of a given mass
  - No explosion  $\rightarrow$  BHs of a given mass



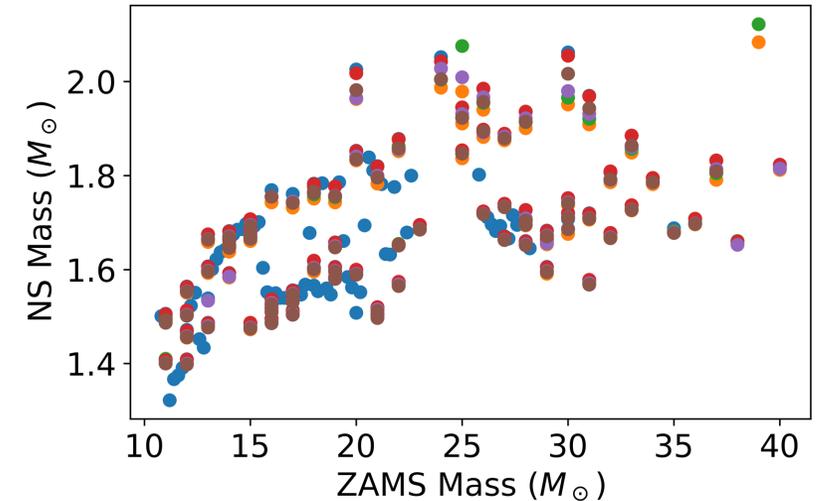
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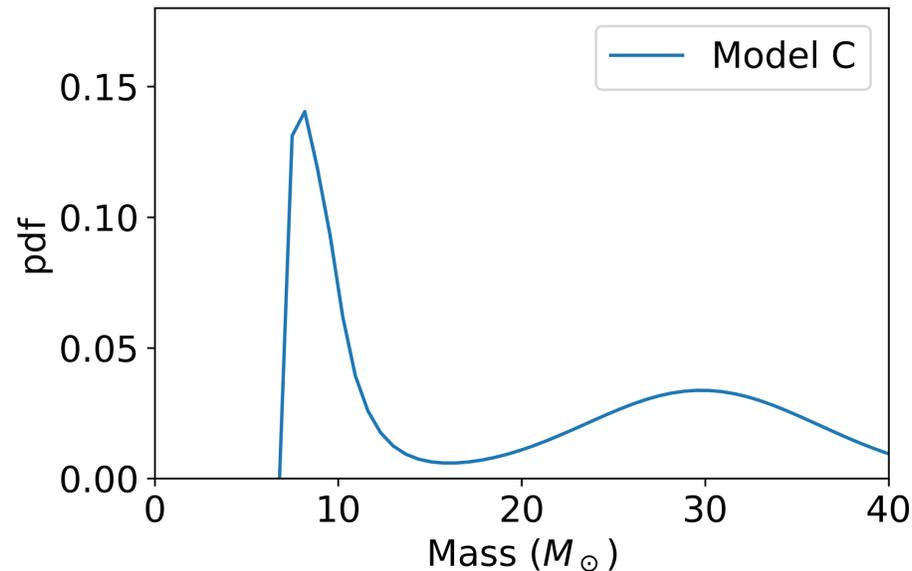


- Distribution of mass:
  - Kroupa initial mass function
- Distribution of metallicity:
  - uniform

$$\xi(M_{\text{ZAMS}}, z) = \begin{cases} 0.035 M_{\text{ZAMS}}^{-1.3} & \text{for } M_{\text{ZAMS}} < 0.5 \\ 0.019 M_{\text{ZAMS}}^{-2.2} & \text{for } 0.5 \leq M_{\text{ZAMS}} < 1.0 \\ 0.019 M_{\text{ZAMS}}^{-2.7} & \text{for } M_{\text{ZAMS}} \geq 1.0 \end{cases}$$

# Data from Nature: Observations of Black Holes (BHs)

- Black holes:  $p(M_{\text{BH}}) = C M_{\text{BH}}^{-\alpha} \Theta(M_{\text{BH}}^+ - M_{\text{BH}}) S(M_{\text{BH}}^-, \delta_m)$



$$M_{\text{BH}}^- = 6.8$$

$$M_{\text{BH}}^+ = 75,$$

$$\delta_m = 3.$$

$$\alpha = 7.1$$

- Model C from Abbott et al (2019)
- No BHs below  $M_{\text{BH}}^-$
- Truncated power law from  $M_{\text{BH}}^-$  to  $M_{\text{BH}}^+$
- Gaussian distribution of high-mass BHs from pair-instability SNe

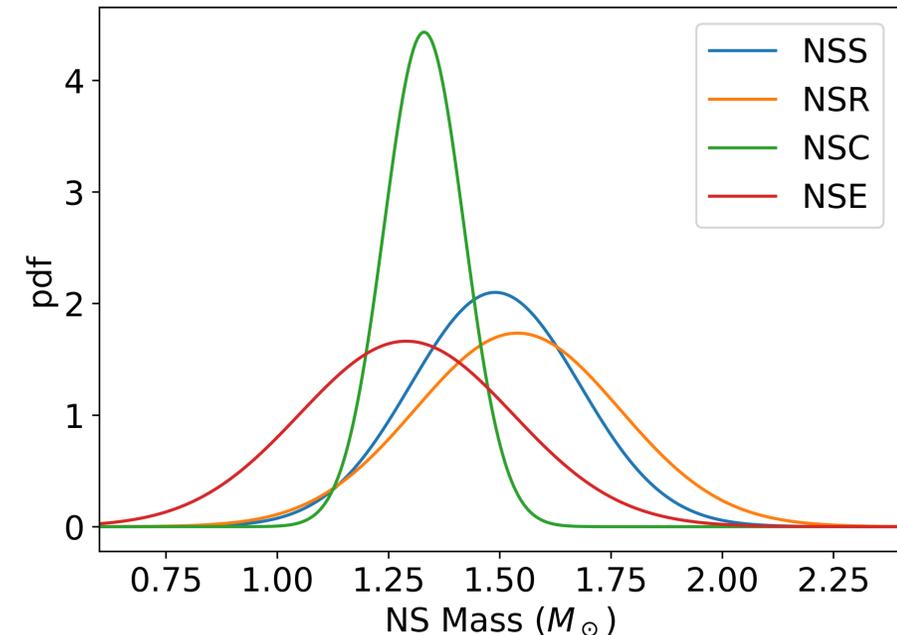
We drop this term because PUSH does not capture PISNe

# Data from Nature: Observations of Neutron Stars (NSs)

- Neutron stars: 
$$p(M_{\text{NS}}) = \frac{1}{\sqrt{2\pi\sigma^2}} e^{-(M_{\text{NS}} - M_0)^2 / (2\sigma^2)}$$

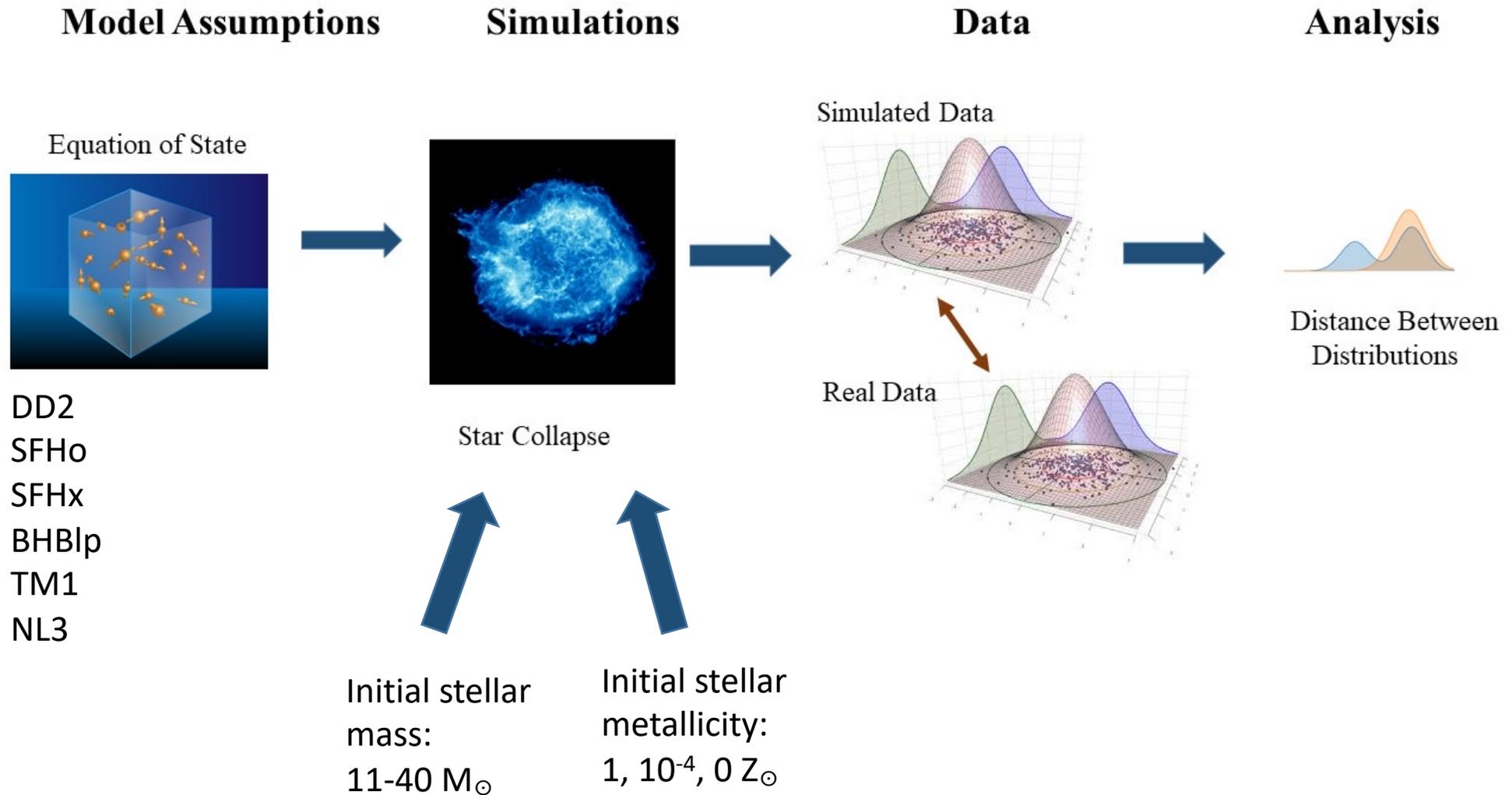
- Split into different astrophysical systems

- For slow pulsars:  $M_0 = 1.49$  and  $\sigma = 0.19$
- For recycled pulsars:  $M_0 = 1.54$  and  $\sigma = 0.23$
- For NSs in binaries with another degenerate object (concentric orbits):  $M_0 = 1.33$  and  $\sigma = 0.09$ .
- For NSs in binaries with another degenerate object (eccentric orbits):  $M_0 = 1.29$  and  $\sigma = 0.24$



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# Overview of methodology



# Measurement Error

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- Observed mass distributions are subject to observational error
- We convert the ensemble of measurements to a distribution via the central limit theorem

- Error  $e$  can be drawn from a Gaussian distribution with mean zero and standard deviation  $\sigma$  for each population

$$p(e; \sigma) = \frac{1}{\sqrt{2\pi\sigma^2}} e^{-e^2/2\sigma^2}$$

- Standard deviation:

- For BHs: from 90% confidence interval in Abbott et al (2019)

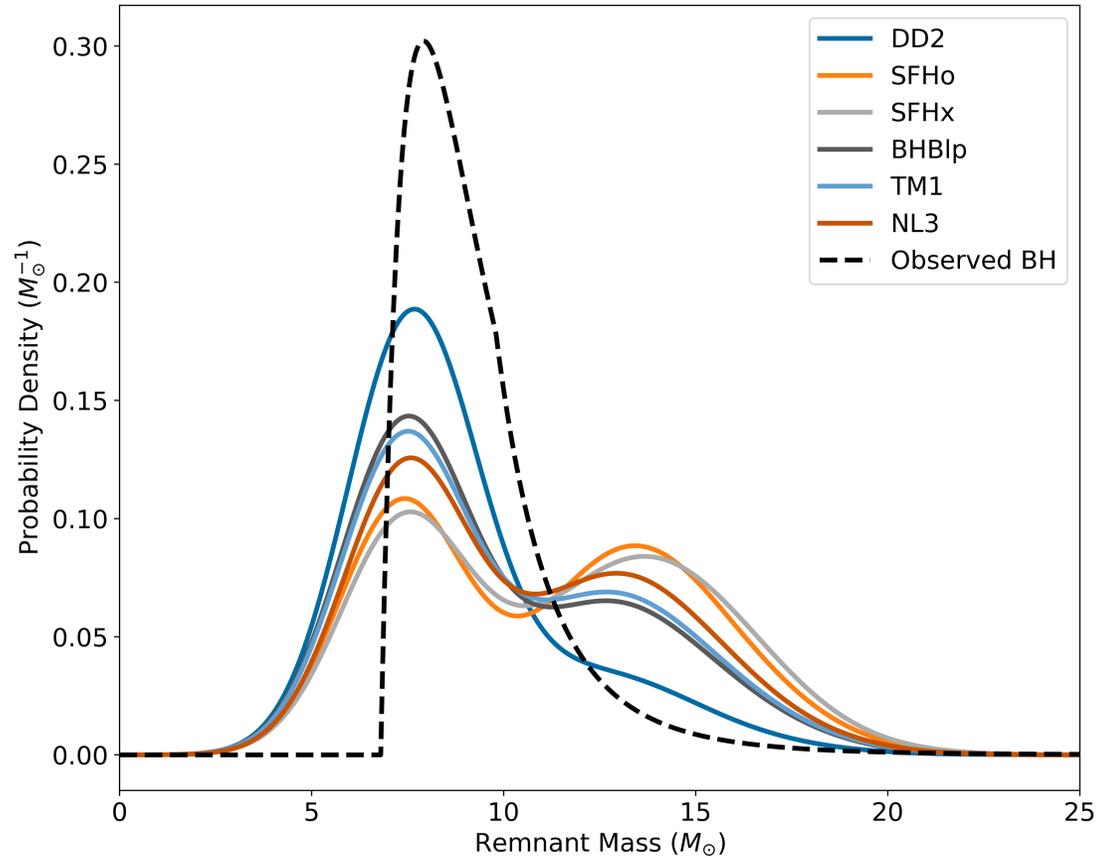
$$\sigma_{\text{BH}}(M_{\text{BH}}) = 0.120213M_{\text{BH}} + 0.355936$$

- For NSs: average width of 90% confidence intervals of NS observations

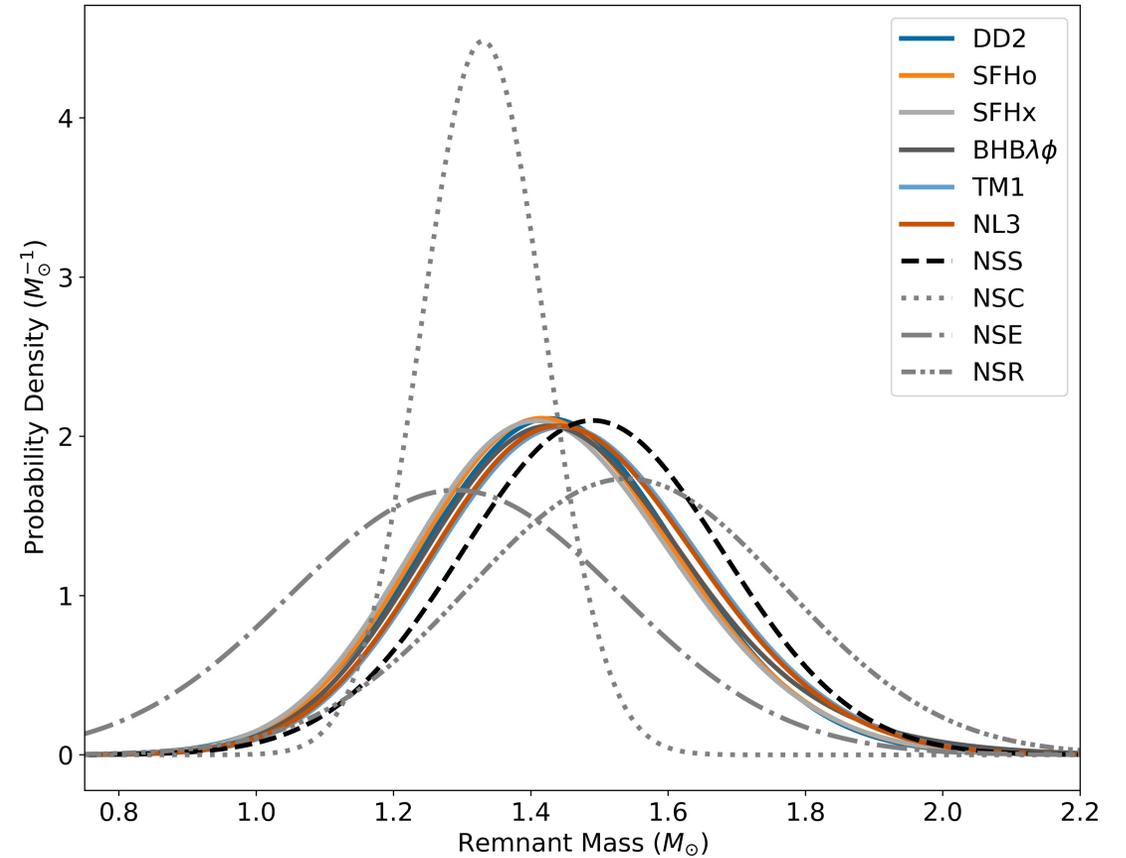
$$\sigma_{\text{NS}} = 0.1.$$

# Results: Probability Density Functions

## Black holes



## Neutron stars



# Statistical distances between $p$ and $p_s$

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- Use two different distance measures:

- **Kullback-Leibler (KL) divergence**  $D_{\text{KL}}(p || p_s) = \int p(\tilde{M}_{\text{BH}}) \log \left( \frac{p(\tilde{M}_{\text{BH}})}{p_s(\tilde{M}_{\text{BH}})} \right) d\tilde{M}_{\text{BH}}$

- Measures the total entropy between  $p$  and  $p_s$

- **Total variation (TV) distance**  $D_{\text{TV}}(p, p_s) = \frac{1}{2} \int |p(\tilde{M}_{\text{BH}}) - p_s(\tilde{M}_{\text{BH}})| d\tilde{M}_{\text{BH}}$

- Measures the maximum distance between the probabilities assigned to an event by two probability distributions

# Results

EOS	BH		NSS	
	$D_{\text{KL}}$	$D_{\text{TV}}$	$D_{\text{KL}}$	$D_{\text{TV}}$
DD2	$0.528^{+0.101}_{-0.068}$	$0.324^{+0.050}_{-0.036}$	$0.066^{+0.058}_{-0.043}$	$0.125^{+0.045}_{-0.052}$
SFH <sub>o</sub>	$1.011^{+0.836}_{-0.270}$	$0.523^{+0.155}_{-0.092}$	$0.071^{+0.081}_{-0.054}$	$0.133^{+0.060}_{-0.069}$
SFH <sub>x</sub>	$0.999^{+1.510}_{-0.312}$	$0.519^{+0.187}_{-0.109}$	$0.082^{+0.098}_{-0.063}$	$0.143^{+0.066}_{-0.075}$
BHB $\lambda\phi$	$0.729^{+0.270}_{-0.135}$	$0.427^{+0.089}_{-0.060}$	$0.044^{+0.066}_{-0.033}$	$0.105^{+0.061}_{-0.064}$
TM1	$0.752^{+0.295}_{-0.152}$	$0.437^{+0.093}_{-0.065}$	$0.018^{+0.056}_{-0.016}$	$0.068^{+0.068}_{-0.050}$
NL3	$0.808^{+0.459}_{-0.178}$	$0.457^{+0.121}_{-0.073}$	$0.025^{+0.060}_{-0.023}$	$0.080^{+0.066}_{-0.062}$

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$D_{KL}$  can suffer from stability issues due to small density values, eg at the tails of the distribution

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normalized to 1

# Results: Distances between distributions

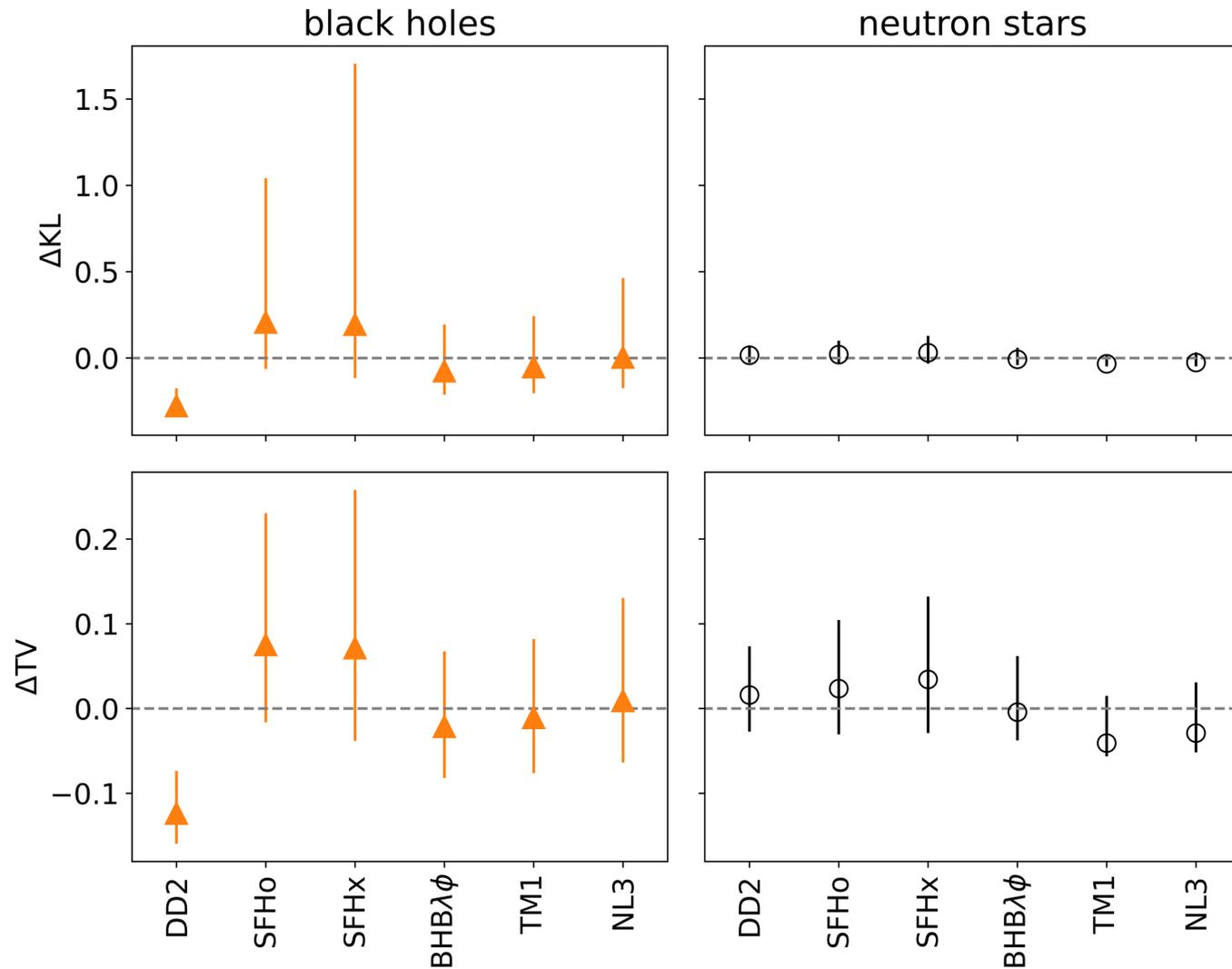
$$\Delta = D_{EOS} - \bar{D}$$

$\Delta < 0 \rightarrow$  more favored

$\Delta > 0 \rightarrow$  less favored

Error bars:

95% confidence interval



# Conclusions...

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- Rankings are slightly different between BHs and NSs
  - Results for NSs are not statistically significant
  - We show results for NSS (most similar to our simulation setup)
- DD2 is most favored
  - Caveat: DD2 was used to calibrate the PUSH parameters
- SFHo and SFHx are mildly disfavored
- BLB $\lambda\phi$ , TM1, NL3 are mildly favored

## ... and improvements

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- Include progenitors with ZAMS masses of 8 -11  $M_{\odot}$ 
  - We miss NSs at the low mass end ( $\sim 1.4 M_{\text{sun}}$ )
- Include binary stars (currently assume isolated stars as progenitors)
- Influence of the PUSH calibration on the results
  - Re-calibrated PUSH using TM1 instead of DD2  $\rightarrow$  no relevant difference
- Understand the influence of the progenitors used
  - All progenitors are from the same stellar evolution code
- Use more observables: NS radius, Ni mass, explosion energy, ...

# More (future) observables

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# Gravitational Waves and the nuclear EOS

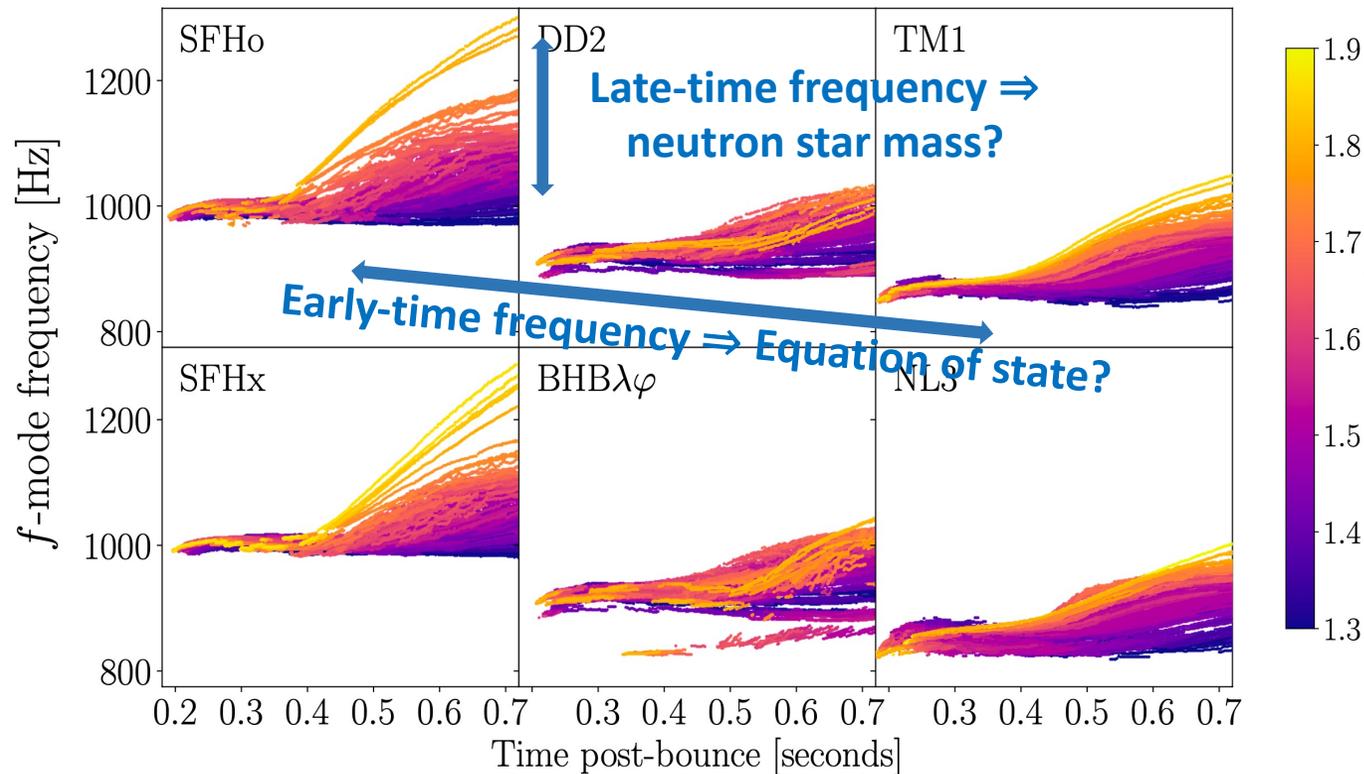
Wolfe+23



Noah Wolfe:  
undergraduate at NCSU  
→ PhD at MIT

## Gravitational Wave Eigenfrequencies from Neutrino-driven Core-collapse Supernovae

Noah E. Wolfe<sup>1,2,3</sup> , Carla Fröhlich<sup>1</sup> , Jonah M. Miller<sup>4,5</sup> , Alejandro Torres-Forné<sup>6,7</sup> , and Pablo Cerdá-Durán<sup>6,7</sup> 



- Linear perturbation analysis of general-relativistic hydrodynamics background  
*Torres-Forné+18, Morozova+18, Torres-Forné+19, Sotani+20*
- Calculate time-frequency evolution (no amplitudes) from spherically-symmetric proto-neutron star background
- Identify frequencies that characterize astrophysical properties of proto-neutron star
  - Universal relations for PNS surface gravity  
*Torres-Forné+19*
  - Multi-messenger observations of core-collapse  
*Warren+20, Nakamura+22*
  - Parameter estimation  
*Bizouard+21, Powell+22*

# Gravitational Waves and the nuclear EOS

Wolfe+23



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