(In-Medium) Similarity Renormalization Group and Effective Field Theory

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Similarity Renormaliztion Group

Similarity Renormalization Group

Basic Idea

continuous unitary transformation of the Hamiltonian to banddiagonal form w.r.t. a given "uncorrelated" many-body basis

• flow equation for Hamiltonian $H(s) = U(s)HU^{\dagger}(s)$:

$$\frac{d}{ds}H(s) = \left[\eta(s), H(s)\right], \qquad \eta(s) \equiv \frac{dU(s)}{ds}U^{\dagger}(s) = -\eta^{\dagger}(s)$$

• choose $\eta(s)$ to achieve desired behavior, e.g.,

$$\eta(s) \equiv \left[H_d(s), H_{od}(s) \right]$$

to suppress (suitably defined) off-diagonal Hamiltonian

• consistent evolution for all observables of interest

Similarity Renormalization Group



Induced Interactions



- SRG is a **unitary** transformation in **A-body space**
- up to **A-body interactions** are **induced** during the flow:

- state-of-the-art: evolve in three-body space, truncate induced four- and higher many-body forces
- flow parameter dependence of eigenvalues is a diagnostic for size of omitted induced interactions

What is "Magic" about EM1.8/2.0?





Roland Wirth

(now at DWD, Offenbach, Germany)

What is "Magic" about EM1.8/2.0?



- The **"magic" NN+3N interaction**, EM1.8/2.0 yields excellent ground-state energies across the nuclear chart, all the way to ²⁰⁸Pb [Simonis et. al., PRC **96**, 014303; Hu et al., Nat. Phys. **18**, 1196]
- **Construction:** [Hebeler et. al., PRC 83, 031301(R)]
 - **NN:** Entem-Machleidt N³LO @ 500 MeV cutoff, SRG evolved to $\lambda = 1.8 \text{ fm}^{-1}$
 - **3N:** N²LO, nonlocal regulator with $\Lambda_{3N} = 2.0 \text{ fm}^{-1}$, c_D and c_E fit to ³H g.s. energy and ⁴He charge radius
- Assumption: Induced 3N terms can be absorbed into c_D and c_E
- **Test:** Evolve EM2.0/2.0 to $\lambda = 1.8 \text{ fm}^{-1}$ and project 3N force onto N²LO topologies

Projecting 3N Forces



- Use chiral N²LO operators $O_{1,3,4,D,E}$ with $\Lambda_{3N} = 2.0 \, \text{fm}^{-1}$ as a basis for 3N force
 - represented as three-body Jacobi HO matrices
- Frobenius inner product:

$$\left\langle U, V \right\rangle \equiv \sum_{J^{\pi}T} \operatorname{tr} \left(U_{J^{\pi}T}^{\dagger} V_{J^{\pi}T} \right)$$

 $G\mathbf{c} = \mathbf{y}$

• basis is **not orthogonal** - introduce **metric** $G_{ij} \equiv \langle O_i, O_j \rangle$

• compute
$$\mathbf{y} = \left(\langle O_1, V \rangle, \dots, \langle O_E, V \rangle \right)^T$$
 and solve

• c contains the LECs of the projected interaction

Structure of N²LO Topologies





 $|\langle EJ^{\pi}T | V_{123}(\lambda = \infty) | E'J^{\pi}T \rangle|$

 $\hbar\omega = 36 \,\mathrm{MeV}$

- $n_{\rm reg} = 4$
- $\Lambda = 2.0 \, \mathrm{fm}^{-1}$
- low E, low J
- c_D similar to c_3
- c_E is S-wave only

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Evolving from $2.0 \,\mathrm{fm}^{-1}$ to $1.8 \,\mathrm{fm}^{-1}$



- SRG evolution in **three-body space**
- unitary transformation obeys

$$\frac{dU}{d\lambda} = -\frac{4}{\lambda^5}\eta(\lambda)U(\lambda)$$

• can use $U(\lambda)$ to separate induced 3N interaction, $V_{2\rightarrow 3}(\lambda)$, from evolved initial 3N interaction $V_{3\rightarrow 3}(\lambda)$

Structure of Induced Interactions



$$|\langle EJ^{\pi}T | V_{2\to 3} | E'J^{\pi}T \rangle|$$

$$\hbar\omega = 36 \,\text{MeV}$$

$$\lambda = 1.8 \, \mathrm{fm}^{-1}$$

- contributions from all energies (up to model space truncation)
- dominant diagonal
- different from N²LO topologies

Structure of Induced Interactions





Evolved 3N Interaction





- $|\langle EJ^{\pi}T | V_{3\to 3} | E'J^{\pi}T \rangle|$ $\hbar\omega = 36 \,\text{MeV}$
 - $\lambda = 1.8 \, \mathrm{fm}^{-1}$

- shape is similar to initial 3N force
- weak compared to $V_{2\rightarrow 3}$



LEC	2.0/2.0	2.0/2.0 →1.8		1.8/2.0
		Full	c_D, c_E	
<i>c</i> ₁	-0.81	-0.67	-0.81	-0.81
c_3	-3.20	-2.92	-3.20	-3.20
c_4	5.40	5.14	5.40	5.40
c_D	1.26	1.38	1.45	1.27
c_E	-0.12	-0.13	-0.11	-0.13

- Full: 10% correction to c_i s, 2PE suppressed, contacts enhanced
- c_D, c_E only: D term enhanced, E term (slightly) suppressed
- Final values quite different from EM1.8/2.0



- "Magic" of EM1.8/2.0 seems to be an accidental cancellation: induced 3N terms and excluded higher order NN, 3N, 4N, ... cancel, except for contact terms
- c_D, c_E have the **right size** to fit few-body observables and provide **correct shift** in E/A
- Use this protocol to analyze Δ -full interactions, impact of new 3N forces

SRG in Many-Body Systems

SRG Scales



- split the Hamiltonian: $H(s) = H_d(s) + H_{od}(s)$
- assume that

$$H_d(s) |n\rangle = E_n(s) |n\rangle, \qquad \langle n | H_{od}(s) | n\rangle = 0$$

• **generator** - e.g., Wegner:

$$\begin{split} \langle i | \eta | j \rangle &= \sum_{k} \left(\langle i | H_{d} | k \rangle \langle k | H_{od} | j \rangle - \langle i | H_{od} | k \rangle \langle k | H_{d} | j \rangle \right) \\ &= - \left(E_{i} - E_{j} \right) \langle i | H_{od} | j \rangle \end{split}$$

• flow equation:

$$\frac{d}{ds}\langle i | H | j \rangle = -(E_i - E_j)^2 \langle i | H_{od} | j \rangle$$
$$+ \sum_k (E - i + E_j - 2E_k) \langle i | H_{od} k \rangle \langle k | H_{od} | j \rangle$$

SRG Scales



 assume H_{od}(s) is small - should be a good assumption for some s > s₀ if the SRG flow is working as intended (or if there are perturbative arguments)

$$\frac{d}{ds}E_i = \frac{d}{ds}\langle i | H_d | i \rangle = 2\sum_k (E_i - E_k) |\langle i | H_{od} | k \rangle|^2 \approx 0$$

$$\frac{d}{ds}\langle i|H|j\rangle = \frac{d}{ds}\langle i|H_{od}|j\rangle \approx -(E_i - E_j)^2\langle i|H_{od}|j\rangle$$

• integrate:

$$\langle i | H_{od}(s) | j \rangle = \langle i | H_{od}(s_0) | j \rangle e^{-(E_i - E_j)^2 (s - s_0)}$$

- White generator: e^{-s}
- imaginary time / Brillouin: $e^{-|E_i E_j|s}$





- s characterizes decoupling of energy scales in the manybody system
 - $s \sim f(\Delta E^{-1})$
 - concrete interpretation depends on choice of generator
- carries forward from many-body states to operator
 formulation in IMSRG applies in the same way to 0B, 1B, 2B, ... operators
- Can this be used (more) ?

In-Medium SRG

Operator Bases for the IMSRG



 choose a basis of operators to represent the flow (might involve an educated guess about physics):

$$H(s) = \sum_{i} c_{i}(s)O_{i}, \qquad \eta(s) = \sum_{i} f_{i}(\{c(s)\})O_{i}$$

• close algebra by truncation, if necessary:

$$\left[O_i, O_j\right] = \sum_{ijk} g_{ijk} O_k$$

• flow equations for the coefficient (coupling constants):

$$\frac{d}{ds}c_k = \sum_{ij} g_{ijk} f_i(c)c_j$$

• "obvious" choice for many-body problems

$$\{O_{pq}, O_{pqrs}, \dots\} = \{a_p^{\dagger}a_q, a_p^{\dagger}a_q^{\dagger}a_s a_r, \dots\}$$

Transforming the Hamiltonian





Decoupling in A-Body Space



goal: decouple reference state $|\Phi\rangle$ from excitations

Flow Equation





IMSRG(2) Flow Equations





IMSRG(2) Flow Equations



2-body Flow



Examples of (IM)SRG Invariance



HH et al., Phys. Rept. **621**, 165 M. Frosini et al., EPJA **58**, 64



observables should be invariant under unitary evolution

 (tunable) combination of relevant operators in basis + degrees of freedom of many-body Hilbert space

Examples of (IM)SRG Invariance



S. K. Bogner et al., PPNP 65, 94







• explicit exponential ansatz for unitary transformation:

$$U(s) = S \exp \int_0^s ds' \ \eta(s') = e^{\Omega(s)}$$

• flow equation for Magnus operator:

$$\frac{d}{ds}\Omega = \sum_{k=0}^{\infty} \frac{B_k}{k!} \operatorname{ad}_{\Omega}^k(\eta), \quad \operatorname{ad}_{\Omega}(O) = [\Omega, O]$$

(B_k : Bernoulli numbers)

- construct $O(s) = U(s)OU^{\dagger}(s)$ using Baker-Campbell-Hausdorff expansion (Hamiltonian + effective operators)
- Magnus(2): two-body truncation (as in NO2B, IMSRG(2))

Magnus vs. Direct Integration





Perturbative Treatment of Subleading Forces



- Magnus formulation is convenient for treating specific interactions perturbatively or non-perturbatively
 - possible in MR-IMSRG formalism as well, but much more cumbersome can be used for cross-checks in the future
- generate $\Omega_0(s)$ nonperturbatively
- construct $H_1(s) = e^{\Omega_0(s)}H_1(0)e^{-\Omega_0(s)}$ and treat in finite-order MBPT
 - first tests with interactions from pionless EFT see
 Matthias Heinz' talk
- challenge: implementation of 4N force at NLO
 - will explore normal-ordered approximations

Epilogue





- Projection methods to help analyze 3N (or other) nuclear forces / operators
- Opportunities from **embracing the RG** in IMS**RG** ?
- IMSRG offers convenient pathways for treating specific components of the Hamiltonian perturbatively instead of non-perturbatively see Matthias Heinz' talk

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Supplements

Decoupling





Decoupling





Decoupling





absorb correlations into RG-improved Hamiltonian

$$U(s)HU^{\dagger}(s)U(s)|\Psi_{n}\rangle = E_{n}U(s)|\Psi_{n}\rangle$$

 reference state is ansatz for transformed, less correlated eigenstate:

$$U(\mathbf{s}) \left| \Psi_n \right\rangle \stackrel{!}{=} \left| \Phi \right\rangle$$

Emulators for the IMSRG





Davison, HH, J. Crawford, S. Bogner, arXiv:2504.xxxx

- non-invasive ROM
 emulator based on
 Dynamic Mode
 Decomposition
 - Δ NNLO_{GO}, NN+3N, $e_{max} = 12, E_{3max} = 14$
- O(10M) samples computational effort reduced by 5+ orders of magnitude

Parametric Matrix Model Emulators



B. Clark, P. Cook, ... also see: S. Yoshida, Particles 2025, 8



- DMD fails for Magnus operator if snapshots are taken during initial stages of flow...
- ... but PMMs seem to work

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