

Four-body systems at large cutoffs in short-range EFTs up to NLO

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Background: Goals and Key ingredients

Semi-analytical EFT studies on four-body (4B) systems

- QCD $\xrightarrow{\Lambda_\chi \sim \mathcal{O}(1)\text{GeV}}$ ChEFT $\xrightarrow{\Lambda_{NN} \sim 300\text{MeV}}$ KSW $\xrightarrow{\Lambda_{\not{E}} \sim 140\text{MeV}}$ ~~EFT~~
- Integral equations, e.g., Faddeev, diagrammatic
- cold ^4He atoms, ^4He nuclei, nuclear cluster/halo systems, etc

Background: Goals and Key ingredients

Renormalization group (RG) in few-body systems, e.g.,

- 2 body: Manifest power counting using dim. reg. with power divergence subtraction
- 3 body: three-body force needed at LO, Efimov physics, and discrete scaling symmetry
- 4 body: four-body force needed starting at NLO (to be shown at *large* cutoffs)

Universality in 4B systems

- Universal relations in the unitarity
- Expansion around unitarity
- Relevant for four-nucleon systems (Wigner-SU(4) symmetry)

- **Intro:** Contact EFT; 2B, 3B and 4B integral equations
- **Problem:** deep 3B poles make large-cutoff 4B calculations difficult
- **Method:** Addressing 3B poles
- **Results at LO and NLO:** 4B binding energies of cold ^4He atoms, their cutoff dependence, and 4B force
- **Universal relations and Comparing with FY results**

Two-Body Systems

2B System with a Pole

$$\begin{aligned} \frac{m}{2\pi} A_2 &= \frac{1}{-\frac{1}{a} + \frac{1}{2}rp^2 + \dots - ip} \quad (\text{effective range expansion}) \quad (1) \\ &= \frac{1}{-\frac{1}{a} - ip} \left[1 + \underbrace{\left(\frac{rp}{2} \right) \left(\frac{ap}{1 + iap} \right)}_Q + \dots \right], \text{ if } |ap| \sim 1, |ra^{-1}| \sim Q \ll 1 \end{aligned}$$

2B pole (non-perturbative) reproduced at the lowest order

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2B pole (non-perturbative) reproduced at the lowest order
Relevant 2B systems with $|a/r| \gg 1$?

(2N) 3S_1 deuteron : $a_t \approx 5.4$ fm, $r_t \approx 1.7$ fm

(2N) 1S_0 virtual state : $a_s \approx -24$ fm, $r_s \approx 2.7$ fm

(2 Boson) Cold atomic ^4He : $a_{^4\text{He}} \approx 100$ Å, $r_v \approx 7$ Å

2B System with a Pole


$$iA = \frac{4\pi}{m_N} \frac{i}{k \cot(\delta) - ik} \stackrel{ERE}{=} \frac{4\pi}{m_N} \frac{i}{-\frac{1}{a} + \frac{rk^2}{2} + \dots - ik}$$
$$\approx \frac{4\pi}{m_N} \frac{i}{-\frac{1}{a} - ik} = \text{Diagram 1} + \text{Diagram 2} + \text{Diagram 3} + \dots$$


Figure 1: Reproduce 2B pole using Born series with 2B contact interaction (geometric sum)

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Why sum to all orders? How to treat different interactions?

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- Interactions ordered by power counting; Observables expanded in, e.g., $M_{lo}/M_{hi} \sim 10\%$ for cold ^4He atoms.
- **Renormalization (cutoff dependence) can inform power counting**

Three-Body Systems

3B System in Contact EFT

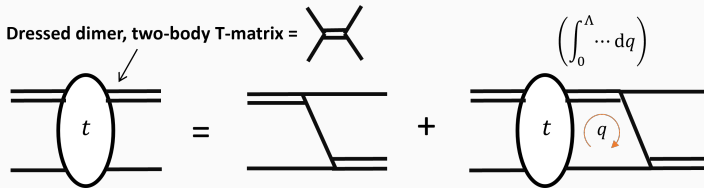


Figure 2: 3B integral equation, equivalent to Faddeev equation

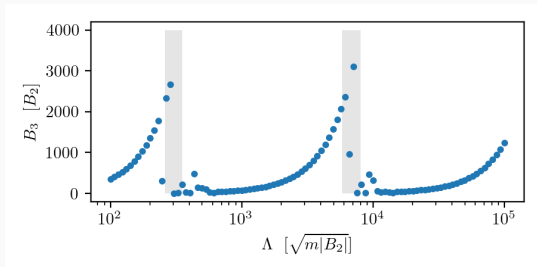


Figure 3: 3B binding energies with 2B contact interaction only

3B System in Contact EFT

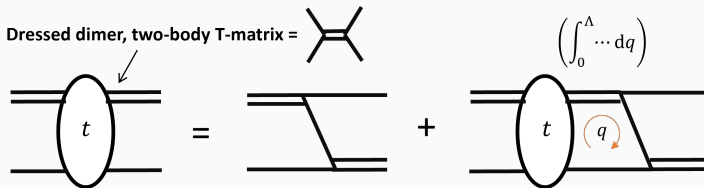


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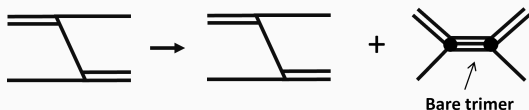


Figure 5: Promoting 3B force to LO to absorb cutoff dependence*

- LO contact EFT: 2B and 3B force (non-perturbative)**

*Bedaque, Hammer, and Kolck 1999

3B System in Contact EFT

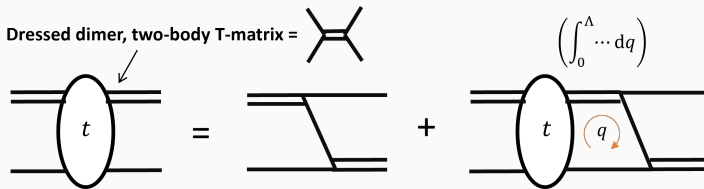
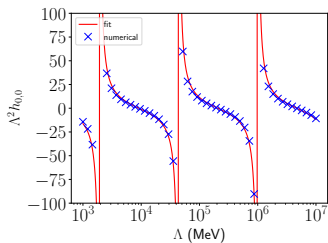
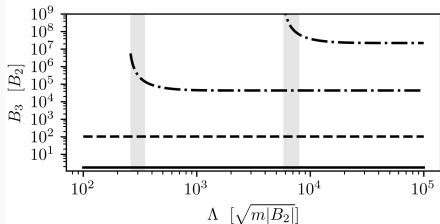


Figure 6: 3B integral equation, equivalent to Faddeev equation



(a) Cutoff dependence of 3B force (illustrative only, ignore units)



(b) Binding energies of Efimov states as a function of cutoff

2B and 3B Systems in Short

Recap of Key concepts in 2B and 3B systems:

- LO interactions: 2B and 3B contact (zero-derivative)
- 2B and 3B bound states occur at LO
- 3B interaction required at LO for RG invariance

Four-Body Systems at LO

4B integral equation

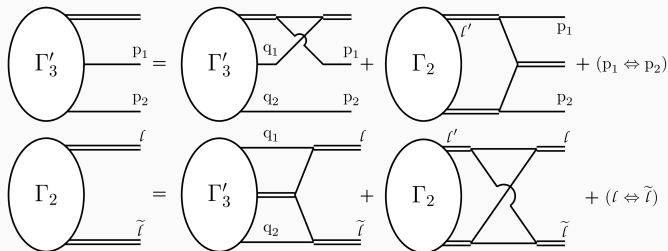


Figure 8: 4B integral equation in terms of 2B force *

- Two types fragmentation: (3B+1B) and (2B+2B)
- 3B poles occurs from iteration of (3B+1B) diagrams

4B integral equation

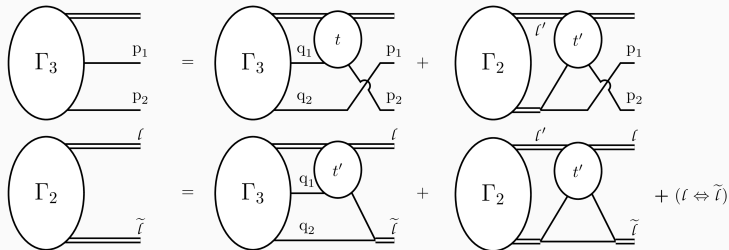
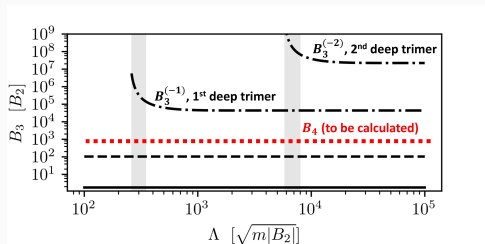


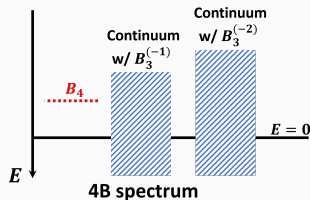
Figure 9: 4B integral equation using 3B amplitude *

- Can **directly address trimer poles** of 3B amplitude. **But why do we care?**

Why worry about 3B poles?



(a) (Schematic) Location of 4B binding energy



(b) 4B "bound" state and continuum states associated with deep trimers

- 3B+1B decay channel available at large cutoffs
- 4B bound state become unstable (become resonance) at large cutoffs
- **Problem: Deep trimers lead to difficult on assessing cutoff dependence**

Dealing with 3B poles in 4B equation

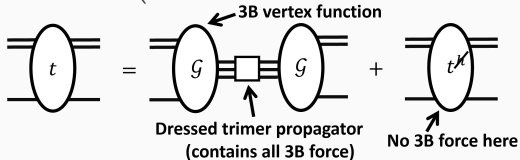
&

LO Results (for atomic ^4He)

Dealing with trimer poles (Method A)

Method A: * **isolate the singularity and then include/exclude it** by Cauchy principal value prescription

$$\hat{G}_3 = \frac{|\psi\rangle R_\psi \langle \psi|}{E + B_3} + \dots$$
$$= |\psi\rangle R_\psi \langle \psi| \left(\pm i\pi\delta(E + B_3) + \text{p.v.} \left(\frac{1}{E + B_3} \right) \right) + \dots$$



- Large-cutoff calculations now possible with 4B equation in terms of 3B amplitudes
- Can compute 4B decay width perturbatively

Results: $B_4^{(n)}$ at LO (cold ^4He atoms)

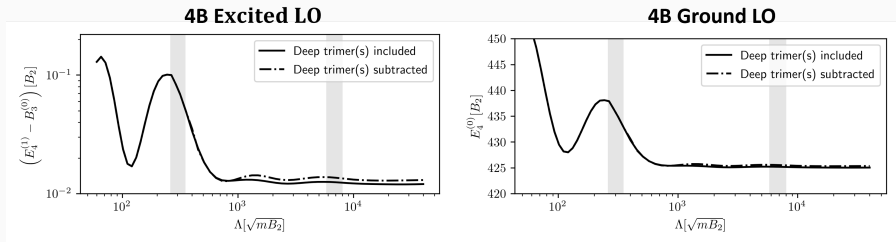


Figure 11: Convergence of $B_4^{(n)}$ at LO

- Gray bands indicate occurrence of deep trimer poles
- Convergence \Rightarrow No 4B force needed at LO.*
- Large cutoffs (regulator-dependent) needed for convergence

*See Platter, Hammer, and Meissner 2004 at relatively low cutoffs

Results: $B_4^{(n)}$ at LO (cold ^4He atoms)

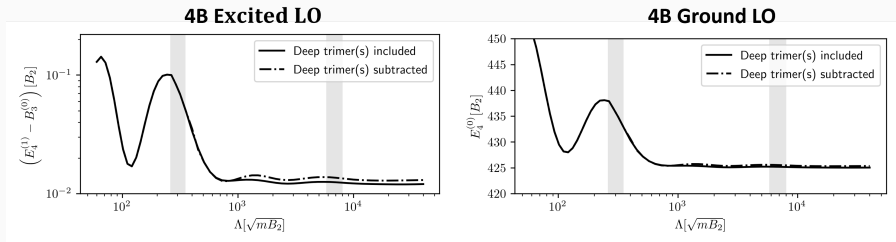


Figure 12: Convergence of $B_4^{(n)}$ at LO

	$B_3^{(0)}$ [mK]	$E_4^{(1)}$ [mK]	$E_4^{(0)}$ [mK]
This work	*128.500	128.517(1)	526.1(5)
Platter et al. 2004	127	128[3]	492[25]
H&K 2012	126.4	127.33	558.98

Table 1: Results for cold ^4He atoms. First two rows are LO EFT calculations. Hiyama and Kamimura 2012 uses realistic potential.

Four-Body System at NLO

&

Preliminary Results (for atomic ^4He)

Results: $B_4^{(n)}$ at NLO (cold ^4He atoms)

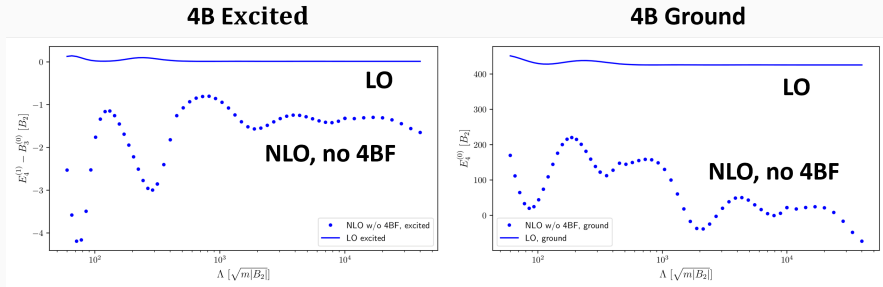
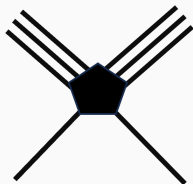


Figure 13: Cutoff dependence of $B_4^{(n)}$ at NLO without 4B force

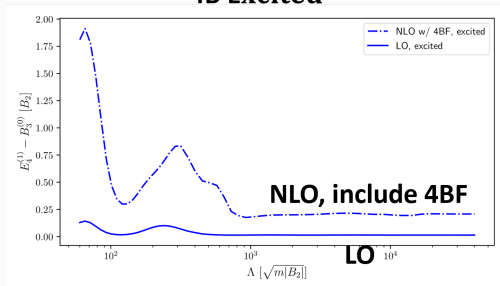
- **Strong cutoff dependence if no 4B force**

Results: 4B force and $B_4^{(1)}$ at NLO (cold ^4He atoms)



(a) 4B force

4B Excited

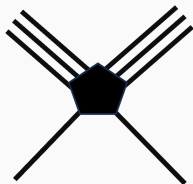


(b) $B_4^{(1)}$ w/ 4B force fitted to $B_4^{(0)}$

- **4B force needed at NLO for renormalization group invariance***

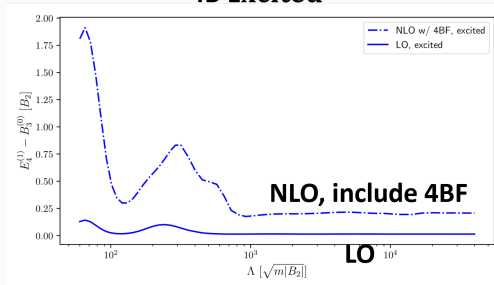
*See Bazak et al. 2019 at low cutoffs

Results: 4B force and $B_4^{(1)}$ at NLO (cold ^4He atoms)



(a) 4B force

4B Excited



(b) $B_4^{(1)}$ w/ 4B force fitted to $B_4^{(0)}$

	$B_3^{(0)}$ [mK]	$E_4^{(1)}$ [mK]	$E_4^{(0)}$ [mK]
This work NLO	*128.500	128.83(1?)	*557
H&K 2012	126.4	127.33	558.98

Table 2: Results for cold ^4He atoms. First row is NLO EFT calculation. Hiyama and Kamimura 2012 uses realistic potential.

4B Universal Relations

Universal scalings

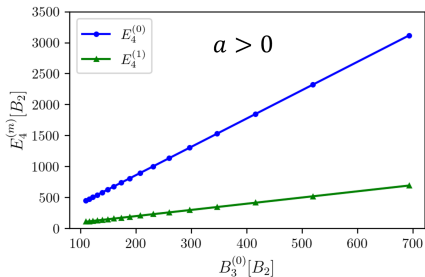
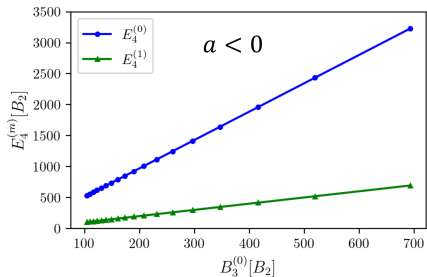


Figure 16: Correlations between trimer and tetramer binding energies; Tjon line* in nuclear physics. a : 2-Boson scattering length.

	$E_4^{(0)}/B_3^{(0)}$	$E_4^{(1)}/B_3^{(0)}$
This work	4.60(1)	1.0022(3)
Deltuva 2010	4.6108	1.00228

Table 3: Universal scalings in the unitary limit.

*Tjon 1975

Four-body results near unitary limit

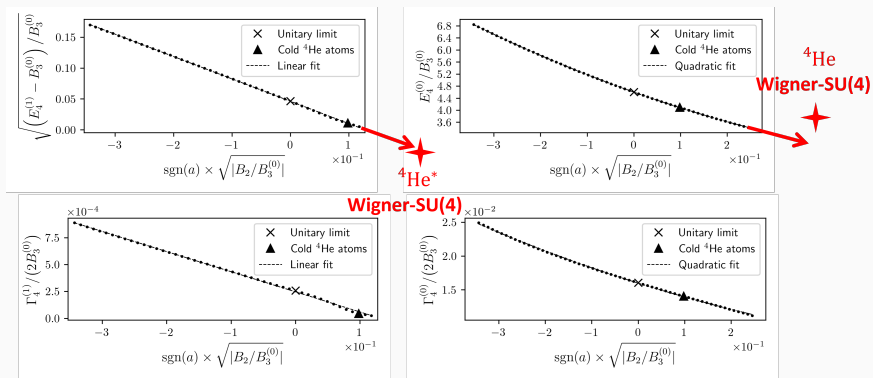


Figure 17: Universal relations between trimer and tetramer binding energies

Four-body results near unitary limit

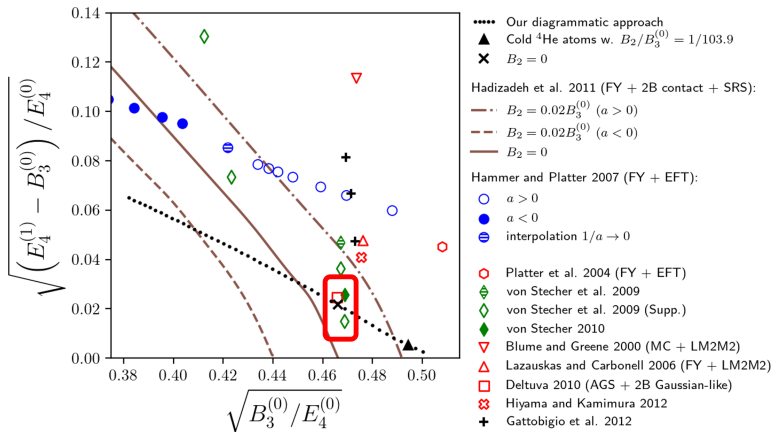
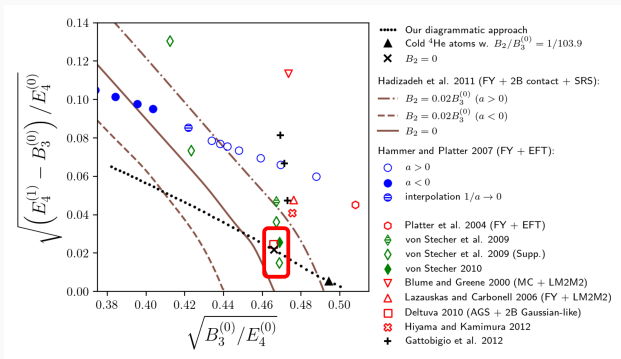


Figure 18: Four-body scaling function*

*Hadizadeh et al. 2011

Four-body results near unitary limit



	$E_4^{(0)}/B_3^{(0)}$	$\Gamma_4^{(0)}/(2B_3^{(0)})$	$E_4^{(1)}/B_3^{(0)}$	$\Gamma_4^{(1)}/(2B_3^{(0)})$
This work	4.60(1)	0.0160(1)	1.0022(3)	$2.57(2) \times 10^{-4}$
Deltuva 2010	4.6108	0.01484	1.00228	2.38×10^{-4}
von Stecher 2010	4.55	-	1.003	-
von Stecher et al. 2009 (Supp., with V_{3b})	4.55	-	1.001	-

Four-body results near unitary limit

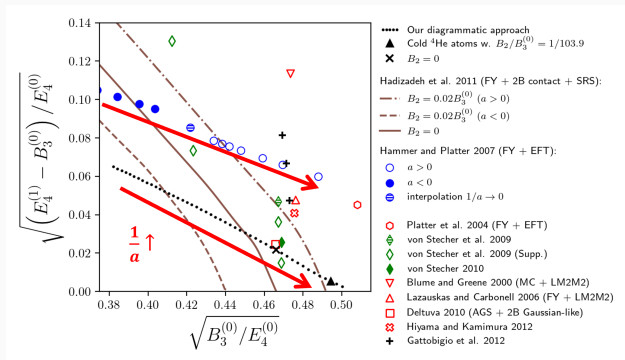


Figure 18: Four-body scaling function*

- Similar flow from unitary limit to cold ^4He atoms;
- Origin of difference?

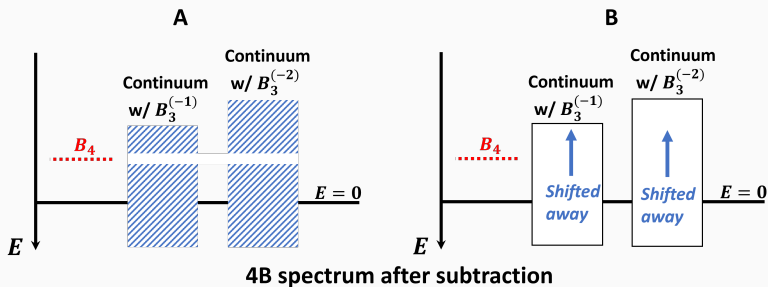
*Hadizadeh et al. 2011

Comparing with FY results

Dealing with trimer poles (Method B)

Method B:* **exclude unwanted deep-trimer spectrum** by shifting them away

$$\begin{aligned}\hat{V}_3 &= |p\rangle V_3 \langle q| \\ \rightarrow \hat{\tilde{V}}_3 &= |p\rangle V_3 \langle q| + |\psi\rangle \eta \langle \psi| \\ \Rightarrow \lim_{\eta \rightarrow \infty} \hat{\tilde{G}}_3 &= \hat{G}_3 - \frac{|\psi\rangle R_\psi \langle \psi|}{E + B_3}\end{aligned}$$



*See, e.g., Lehman 1982

Comparison with FY results in unitary limit

	$E_4^{(0)}/B_3^{(0)}$	$E_4^{(1)}/B_3^{(0)}$
Lin 2024	4.60(1)	1.0022(3)
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Table 4: Universal scalings in the unitary limit.

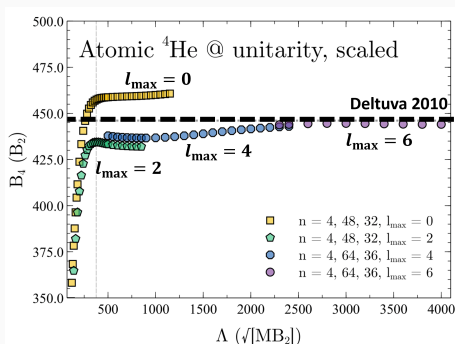
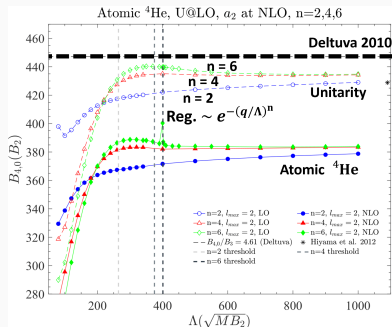


Figure 19: $B_4^{(0)}$ by Faddeev-Yakubovsky equation. Deep trimers shifted by $\eta|\psi\rangle\langle\psi|$. Calculations and figures by Wu, Koenig, and van Kolck.

Conclusion and Outlook

Summary:




- LO and NLO tetramer binding energies of cold ^4He atoms (at large cutoffs);
- 4B force needed starting at NLO in contact EFT;
- 4B universal relations;




Outlook:

- Scattering
- Unitary limit \rightarrow finite $a \rightarrow$ 4-nucleon system (adding Wigner-SU(4)-breaking NN interaction)
- Halo/Cluster systems (suggestion?)




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References

-  Lin, X. **“Four-body systems at large cutoffs in effective field theory”**. In: *Phys. Rev. C* 109.2 (2024), p. 024002. DOI: 10.1103/PhysRevC.109.024002. arXiv: 2304.06172 [nucl-th].
-  Bazak, B. et al. **“Four-Body Scale in Universal Few-Boson Systems”**. In: *Phys. Rev. Lett.* 122 (2019), p. 143001. DOI: 10.1103/PhysRevLett.122.143001. arXiv: 1812.00387 [cond-mat.quant-gas].
-  Hadizadeh, M. R. et al. **“Scaling Properties of Universal Tetramers”**. In: *Phys. Rev. Lett.* 107 (13 Sept. 2011), p. 135304. DOI: 10.1103/PhysRevLett.107.135304. URL: <https://link.aps.org/doi/10.1103/PhysRevLett.107.135304>.

-  Deltuva, A. **“Efimov physics in bosonic atom-trimer scattering”**. In: *Phys. Rev. A* 82 (2010), p. 040701. DOI: 10.1103/PhysRevA.82.040701. arXiv: 1009.1295 [physics.atm-clus].
-  Brodsky, I. V. et al. **“Exact diagrammatic approach for dimer-dimer scattering and bound states of three and four resonantly interacting particles”**. In: *Phys. Rev. A* 73 (3 Mar. 2006), p. 032724. DOI: 10.1103/PhysRevA.73.032724. URL: <https://link.aps.org/doi/10.1103/PhysRevA.73.032724>.
-  Platter, L., H. W. Hammer, and Ulf-G. Meissner. **“The Four boson system with short range interactions”**. In: *Phys. Rev. A* 70 (2004), p. 052101. DOI: 10.1103/PhysRevA.70.052101. arXiv: cond-mat/0404313.

References iii

-  Bedaque, Paulo F., H. W. Hammer, and U. van Kolck. **“The Three boson system with short range interactions”**. In: *Nucl. Phys. A* 646 (1999), pp. 444–466. DOI: 10.1016/S0375-9474(98)00650-2. arXiv: nucl-th/9811046.
-  Lehman, D. R. **“Excluded bound state in the $S_{\frac{1}{2}}$ $\alpha - N$ interaction and the three-body binding energies of ${}^6\text{He}$ and ${}^6\text{Li}$ ”**. In: *Phys. Rev. C* 25 (6 June 1982), pp. 3146–3154. DOI: 10.1103/PhysRevC.25.3146. URL: <https://link.aps.org/doi/10.1103/PhysRevC.25.3146>.
-  Tjon, J.A. **“Bound states of 4He with local interactions”**. In: *Physics Letters B* 56.3 (1975), pp. 217–220. ISSN: 0370-2693. DOI: [https://doi.org/10.1016/0370-2693\(75\)90378-0](https://doi.org/10.1016/0370-2693(75)90378-0). URL: <https://www.sciencedirect.com/science/article/pii/0370269375903780>.