

# STAR Results in the Upcoming EIC Era

Grigory Nigmatkulov (for the STAR experiment) University of Illinois Chicago

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# Outline

- J/ $\psi$  production in ultraperipheral collisions
- Vector meson spin interference
- Baryon number carrier and its transport
- EM-jet  $A_N$  studies at 200 GeV
- Nonlinear gluon effects in QCD



# $J/\psi$ production in ultraperipheral collisions

## The strongest EM-fields in UPCs



● In UPCs,

 $\mathrm{E}_{max} = 10^{18}$  V/m ,  $\mathrm{B}_{max} \sim 10^{14} - 10^{18}$  T

=> Strongest EM-field in the universe, but transient

EM-field treated in terms of quasi-real photons

$$\mathrm{W}_{\gamma,max}\sim\gamma\hbar c/R$$
 ;

$$W_{\gamma,max} \sim 3 \text{ GeV (RHIC)}$$
  
 $W_{\gamma,max} \sim 80 \text{ GeV (LHC)}$ 

=> EM-fields are quantized as photons in UPCs

## **Photon-gluon scattering**





Photoproduction of Vector Mesons (VM) in UPC

## UPC VM: Powerful probe of parton densities inside nuclei





 Probes parton density & fluctuations inside nuclei – constraints for A+A initial state

Modification of parton densities in heavy nuclei
 => VMs help to probe parton density inside nuclei before EIC era

y [fm]

x [fm]

## **UPC events with STAR detector**





Neutron(s) detected in ZDCs

=> Method to trigger UPC events

- ZDC signals show peak structure for neutrons
- No activity in both BBCs => Diffractive events ( $\eta$ -gap)

## $J/\psi$ measurements in 200 GeV Au+Au UPCs

STAR, PRC 110 (2024) 014911





=> Coherent and incoherent contributions can be disentangled via the combined fit of mass and  $p_T$ 





## Rapidity dependence $J/\psi$ production cross-section



 Measured for coherent and incoherent contributions for different neutron emission in ZDCs

- Systematic unc. in incoherent to coherent cross-section ratio are largely cancelled
- Sensitive to the nuclear structure and deformation

=> Important to constrain theoretical models related to nuclear geometry



STAR, arXiv:2311.13637

## Incoherent J/ $\psi$ production cross-section vs $p_T^2$

STAR, arXiv:2311.13632

 Incoherent production compared with H1 data with free proton

 Strong nuclear suppression (~49%) seen (Mäntysaari et. al, Phys. Rev. Lett. 117 (2016) 5, 052301)
 Models found H1 data supports subnucleonic fluctuations

(Mäntysaari et. al, Phys. Rev. D **106** (2022) 7, 074019) STAR data shows the bound nucleon has similar shape as the free proton — similar sub-nucleonic fluctuations in heavy nuclei

=> Strong nuclear suppression and subnucleonic fluctuations in Au nucleus





# Vector meson spin interference

## **Polarized Photons from colliding nuclei**

#### STAR, Phys. Rev. Lett. 127 (2021) 52302



Transverse view of Lorentz contracted nuclei



=> Photons in UPC are linearly polarized

Experimental access to photon polarization demonstrated by STAR, measuring the Breit-Wheeler process,  $\gamma\gamma \rightarrow e^+e^-$ 



#### => The cos(2φ) modulation in VM momentum distribution w.r.t photon polarization direction



Photon polarization correlated with Impact parameter -> random from one event to the next

= Event average washes out the cos(2 $\phi$ ) modulation w.r.t photon polarization direction

## Interference makes the modulation observable in experiment some



Light Source Metal Sheet Screen

# **Double Slit Experiment**

Best analogy: Double slit experiment in Optics

=> Two indistinguishable paths may interfere and make the  $cos(2\phi)$  modulation observable

Photon source ambiguity: Interference among amplitudes of two possible paths

# Observation of interference for $\rho^0 \rightarrow \pi^+ \pi^-$ at STAR





STAR, Sci. Adv. 9, eabq 3903 (2023)

Observed the interference for coherent  $\rho^0$  photoproduction in UPCs

#### SCIENCE ADVANCES | RESEARCH ARTICLE

**PHYSICS** STAR, Sci. Adv. 9, eabq 3903 (2023) Tomography of ultrarelativistic nuclei with polarized photon-gluon collisions

STAR Collaboration

A linearly polarized photon can be quantized from the Lorentz-boosted electromagnetic field of a nucleus traveling at ultrarelativistic speed. When two relativistic heavy nuclei pass one another at a distance of a few nuclear radii, the photon from one nucleus may interact through a virtual quark-antiquark pair with gluons from the other nucleus, forming a short-lived vector meson (e.g.,  $\rho^0$ ). In this experiment, the polarization was used in diffractive photoproduction to observe a unique spin interference pattern in the angular distribution of  $\rho^0 \rightarrow \pi^+\pi^-$  decays. The observed interference is a result of an overlap of two wave functions at a distance an order of magnitude larger than the  $\rho^0$  travel distance within its lifetime. The strong-interaction nuclear radii were extracted from these diffractive interactions and found to be 6.53 ± 0.06 fm (<sup>197</sup>Au) and 7.29 ± 0.08 fm (<sup>238</sup>U), larger than the nuclear charge radii. The observable is demonstrated to be sensitive to the nuclear geometry and quantum interference.

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#### Measured in 3 different collision systems: Au+Au, U+U, p+Au -> Sensitive to nuclear shape/size

## The p<sub>T</sub> dependence of interference for $\rho^0 \rightarrow \pi^+\pi^-$ at STAR



STAR, Sci. Adv. 9, eabq 3903 (2023)

Clear  $p_T$  dependence of interference observed

#### Interference gets weak at higher p<sub>T</sub> — Incoherent processes take over

SCIENCE ADVANCES | RESEARCH ARTICLE

**PHYSICS** STAR, Sci. Adv. 9, eabq 3903 (2023) Tomography of ultrarelativistic nuclei with polarized photon-gluon collisions

 $\mathbf{D}^{\mathsf{U}}$ 

**STAR Collaboration** 

A linearly polarized photon can be quantized from the Lorentz-boosted electromagnetic field of a nucleus traveling at ultrarelativistic speed. When two relativistic heavy nuclei pass one another at a distance of a few nuclear radii, the photon from one nucleus may interact through a virtual quark-antiquark pair with gluons from the other nucleus, forming a short-lived vector meson (e.g.,  $\rho^0$ ). In this experiment, the polarization was used in diffractive photoproduction to observe a unique spin interference pattern in the angular distribution of  $\rho^0 \rightarrow \pi^+\pi^-$  decays. The observed interference is a result of an overlap of two wave functions at a distance an order of magnitude larger than the  $\rho^0$  travel distance within its lifetime. The strong-interaction nuclear radii were extracted from these diffractive interactions and found to be 6.53 ± 0.06 fm (<sup>197</sup>Au) and 7.29 ± 0.08 fm (<sup>238</sup>U), larger than the nuclear charge radii. The observable is demonstrated to be sensitive to the nuclear geometry and quantum interference of nonidentical particles.

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# Radius measurement with interference for $\rho^0 \rightarrow \pi^+ \pi^-$ at STAR

STAR, Sci. Adv. 9, eabq 3903 (2023)



Impact of spin interference on |t| distribution studied in different φ bins

Improved measurement of mass radii using spin interference effect

R (Au) =  $6.53 \pm 06$  fm; R (U) =  $7.29 \pm 08$  fm

## Spin interference with $J/\psi \rightarrow e^+e^-$







Mass: 0.7 GeV/c2Mass: 3.1 GeV/c2Lifetime: 1.3 fm/cLifetime: 2160 fm/c

J/U

Measured sign of the interference tells us the level of interference

 $-> J/\psi$  heavier than  $\rho^0$  and  $J/\psi$  has much longer lifetime

Interference of quantum particles -> Spin interference –> J/ψ decay length much
 longer than typical distance b/w
 two colliding nuclei in UPCs

 –> Probes finer structure and captures high quality images of the gluon distributions

## Measured spin interference with $J/\psi \rightarrow e^+e^-$





Observable for  $J/\psi$  spin interference

Interference signal fitted with:  $1 + a_2$  $\cos(2\phi) \Rightarrow a_2$  is the measure of the modulation



Observed spin interference for  $J/\psi \rightarrow e^+e^-$ 

## Corrections of interference signal due to 2y background



• The  $\gamma + \gamma \rightarrow e^+ + e^-$  has also the  $J/\Psi$  interference like pattern due to detector effect

• We correct for the 2
$$\gamma$$
 process with :  $a_2 = f \times a_2^{bkg} + (1 - f) \times a_2^{sig}$ , with  $f = \frac{N_{bkg}}{N_{sig} + N_{bkg}}$ 

=> Background correction is done to extract true modulation signal

## Corrections of interference signal due to bremsstrahlung process





• We considered the Bremsstrahlung process and  $J/\Psi \to e^+ + e^- + \gamma$ , using the STARLight+Geant simulations

=> Bremsstrahlung correction performed for true modulation signal

## Signal for J/ψ Spin interference



 $\odot$  Measured and corrected signal for  $J/\Psi$  spin interference:

 $a_2 = 0.102 \pm 0.027 \pm 0.029$ 

Measurement has ~3σ significance above zero

 Compared with STARLight and theory calculations

STARLight has no spin interference physics
 – consistent with zero

 Theory (Diffractive+Interference) predicts negative modulation

Diff+Int predictions : Mäntysaari et al. Phys.Rev.C 109 (2024) 2, 024908

=> Observed spin interference signal ~10% in the measured kinematic range

## The $p_{\rm T}\text{-dependent}$ interference of J/ $\psi$



 $_{\odot}$  Interference signal shows strong  $p_{T}$  dependence and rises toward positive

- STARLight predicts zero
- ${\ensuremath{\, \circ}}$  Diffractive+interference calculations are negative at low and high  $p_T$
- Diffractive+interference with additional soft  $\gamma$  radiation predicts negative at low  $p_T$  and rises towards positive value at higher  $p_T$



Diff+Int predictions : Mäntysaari et al. Phys.Rev.C 109 (2024) 2, 024908 Diff+Int+Rad predictions : Brandenburg et. al, Phys. Rev. D 106, 074008 (2022)

=> Modulation strength in data positively increases with  $p_T$  in the measured kinematics



# Baryon number carrier and its transport

#### What carries the baryon number?

#### https://en.wikipedia.org/wiki/Proton https://en.wikipedia.org/wiki/Baryon

In particle physics, the baryon number is a strictly conserved additive quantum number of a system.

Baryons, along with mesons, are hadrons, particles composed of quarks. Quarks have baryon numbers of  $B = \frac{1}{3}$  and antiquarks have baryon numbers of  $B = -\frac{1}{3}$ . The term "baryon" usually refers to *triquarks*—baryons made of three quarks ( $B = \frac{1}{3} + \frac{1}{3} + \frac{1}{3} = 1$ ).



1963-70



# In conventional picture, baryon number is assumed to be carried by the valence quarks each carrying 1/3

Goldberg and Y. Ne'eman, Nuovo Cimento 27 (1963) 1 Gell-Mann, Zweig, 1964, SLAC 1970 Review: hep-ph/9301246

# Baryon number may flow with the flow of the Y-shaped string junction (QCD topology)

X. <u>Artru</u>, Nucl. Phys. B 85, 442–460 (1975), G.C. Rossi and G. Veneziano, Nucl. Phys.B123(1977) 507; Phys. Rep.63(1980) 149 Kharzeev, Phys. Lett. B, 378 (1996) 238-246

#### No experiment has conclusively established the true carrier of baryon number, two different carriers for Q & B inside a baryon possible

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## Gluonic junction as a carrier of baryon number

Kharzeev, Phys. Lett. B, 378 (1996) 238-246, Lewis et. al, arXiv:2205.05685





#### **Junction-Junction**

D. Kharzeev <sup>a, b</sup>





Baryon junction:  $e^{-\alpha_B(y-Y_{\text{beam}})}$   $0.42 \le \alpha_B \le 1$ 

PYTHIA 6 (Quarks):  $e^{-2.5(y-Y_{\rm beam})}$ 

Regge theory can predict rapidity dependence of baryon stopping for junctions Larger transport to mid-rapidity for gluonic junction than valence quarks as baryon carrier

#### Strategies for tracing the baryon carrier

Check if charge and baryon are carried by the same object



Compare electric-charge with baryon transport

Q <-> Z/A x B

Test expectations for valence quark transport with rapidity & centrality

A

Α

A



b



Test if the baryon carrier is a gluonic object by colliding with a photon of very small stopping power



# Rapidity dependence of dN/dy(B) in $\gamma$ +A collisions

#### Centrality dependence of dN/dy(B) vs. y-Y<sub>beam</sub>

<~~~

#### Electric charge vs. baryon transport





Charge stopping  $\simeq \frac{Z}{A} \times Baryon$  stopping



Valence quarks carry electric charge & junction carry baryon Charge stopping  $< \frac{Z}{A} \times Baryon$  stopping Baryon transport at mid-rapidity:  $B = (N_p - N_{\bar{p}}) + (N_n - N_{\bar{n}})$ 

Not difficult except for "n" measurement

Charge transport at mid-rapidity:

$$Q = (N_{\pi^+} + N_{K^+} + N_p) - (N_{\pi^-} + N_{K^-} + N_{\bar{p}})$$

Precision measurement is difficult : isospin conservation, efficiency effects

 $\Delta Q$  and B transport should correlated for valence quark picture not for junctions

#### Precision measurements in isobar collisions





Zirconium: A=96 (Total baryon) Z=40 (Total charge)

Overcome precision problem: 1) compare two isobars, 2) express difference as ratios:

$$R2_{\pi} = \frac{(N_{\pi^+}/N_{\pi^-})^{\mathrm{Ru}}}{(N_{\pi^+}/N_{\pi^-})^{\mathrm{Zr}}}$$

Q transport difference between isobars:

$$\Delta Q = N_{\pi} \left[ (R2_{\pi} - 1) + \frac{N_K}{N_{\pi}} (R2_K - 1) + \frac{N_p}{N_{\pi}} (R2_p - 1) \right]$$

Neutron using deuteron, proton

B transport, same in two isobars:  $B = (N_p - N_{\bar{p}}) + (N_n - N_{\bar{n}})$ 

Using isobar goal is to test:  $\Delta Q \leftrightarrow \frac{\Delta Z}{\Delta X} \times B$ 

$$\frac{N_{\bar{n}}}{N_n} = \frac{N_p}{N_{\bar{p}}} \frac{N_{\bar{d}}}{N_d}$$



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#### Precision measurements in isobar collisions



First measurements of electric charge stopping using isobar collisions

Data: More baryon transported to central rapidity than electric charge

Non-junction models (UrQMD, HERWIG): equal or less baryon compared to electric charge

#### Not compatible with same carrier of electric charge and baryon

#### Rapidity distribution of baryon production: Global data

STAR data: N. Lewis, et. al., arXiv:2205.05685, BRAHMS+NA49: F. Videbaek, 1st workshop on baryon dynamics, SBU, 2024

#### Baryon transport with rapidity loss (y-Y<sub>beam</sub>)



#### BRAHMS + NA49 data (wider y-Y<sub>beam</sub>)



Exponential with slope 0.63±0.2, no change with centrality for 2<Y<sub>beam</sub> <5.5 At higher energy rapidity slope closer to~0.5 lower energy (ly-Y<sub>beam</sub>l<2) rapidity slope ~1

Rapidity slope of baryon density: centrality independent, depends on ly-Y<sub>beam</sub>l range Grigory Nigmatkulov, INT, Aug. 19-23, 2024, Seattle (WA)



Search for non-zero net-baryon in photon-ion collisions near central-rapidity

#### Triggering inclusive photon-induced processes by the STAR detector



#### Time Projection Chamber (TPC)

- Track reconstruction
- Identify particles using dE/dx

#### Time-Of-Flight detector (TOF)

- Extend particle identification to high pT
- Pile-up rejection

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Lewis et. al, arXiv: 2205.05685, BeAGLE: W. Chang, et al PRD 106, 012007 (2022)

**BeAGLE**  $v^*Au \rightarrow X$  $p_{T} > 0.2 \text{ GeV/c}$ 

Use characteristic asymmetric particle production to trigger inclusive  $\gamma$ +Au events with help of:

- Beam-Beam counter (BBC),
- Zero-Degree Calorimeter (ZDC),
- Vertex Position Detector (VPD)

#### Results: Rapidity distribution of net-proton in $\gamma$ +Au events



p and net-proton dN/dy with y described by an exponential with slope:  $1.13 \pm 0.32$ 

Anti-proton distribution is near constant with y

Compared Au+Au slope:  $0.63 \pm 0.02$  (2<Y<sub>beam</sub> <5.5)

Compared to PYTHIA, which does not include a baryon junction mechanism, predicts a slope of 2.5

Exponential slope of rapidity dependence of net-proton lower than PYTHIA predictions

## Rapidity slope of net-proton: Global data



X. Artru, M. Mekhfi, Nucl. Phys. A 532 (1991) 351 BRAHMS+NA49: Videbaek, 1st workshop on baryon dynamics, SBU 2024

Au+Au slope same for all centrality Slope  $\gamma$ +Au >~ Slope Au+Au:

Closer to the fit to BRAHMS + NA49 data slope to ~1 for  $Y_{beam} < 2$ (NA49 energy ~17 GeV closer to  $\gamma$ +Au cm energy ~ 10 GeV)

Slope has  $Y_{beam}$  (energy) dependence  $\alpha_B = \alpha_B (|y-Y_{beam}|)$ 

Consistent with Regge theory baryon-junction prediction but smaller than PYTHIA/HERWIG

Rapidity dependence of net-proton in  $\gamma$ +Au collisions compatible with junction picture



# EM-jet A<sub>N</sub> studies at 200 GeV

# STAR forward detectors used in the current STAR analyses:

- Forward Meson Spectrometer (FMS):
   2.6 < η < 4.2, φ ∈ (0, 2π); detect γ, π<sup>0</sup>, η
- Roman Pot detector (RP): Located about 15 m away from interaction point on both sides; detects slightly scattered protons



## Multi-dimensional studies for inclusive EM-jet $A_N$ at 200 GeV



The ElectroMagnetic jets (EM-jets) are



- The EM-jet  $A_N$  decreases with increasing photon multiplicity for  $x_F > 0$ 
  - A<sub>N</sub> is larger for the EM-jets consisting of 1 or 2 photons
- A<sub>N</sub> increases with x<sub>F</sub> for all the cases of photon multiplicity

#### Single diffractive EM-jet $A_N$ at 200 GeV



- The EM-jet A<sub>N</sub> for x<sub>F</sub> > 0 (>2σ significance of non-zero) is observed for 1 or 2 photon multiplicity EM-jets in the single diffractive process
- A<sub>N</sub> for the three processes consistent with each other within uncertainty
- The single diffractive processes fail to provide evidence for its significant contribution to large A<sub>N</sub> in the inclusive processes



## Semi-exclusive process EM-jet $A_{\rm N}$ at 200 GeV





- Semi-exclusive process: polarized proton intact; constrain the energy of EM-jet at FMS and west side proton to less than beam energy
- A non-zero  $A_N$  for  $x_F > 0$  is observed with 3.3 $\sigma$  significance for the semi-exclusive process
- The sign of  $A_N$  is negative. Theoretical inputs are needed to understand the different sign





# Nonlinear gluon effects in QCD



- DGLAP (Dokshitzer Gribov Lipatov Altarelli Parisi): Evolution in resolution Q<sup>2</sup>, resums terms in  $\alpha_{s} \log Q^{2} \rightarrow$  resolving "smaller" partons at high Q
- BFKL (Balitski Fadin Kuraev Lipatov (BFKL): Evolution in energy x, resums terms in α<sub>S</sub> log 1/x → Large parton densities at small x
- Saturation region at very small x
- Important to understand QCD evolution, parton densities

## Nonlinear gluon effects in QCD

M.S. Abdallah et al., Phys. Rev. Lett. 129, 092501



First measurement of the A dependence of nonlinear gluon effects





- At low p<sub>T</sub> regime, a clear suppression is observed in p+A compared to the p+p data
- Such suppression scaling with A<sup>1/3</sup> matches gluon saturation models

# Summary



- STAR measured the coherent and incoherent J/ $\psi$  production in Au+Au UPCs
- STAR observed the spin interference of the photoproduced  $\rho^0$  and  $J/\psi$ 
  - The measured interference signal increases with  $\ensuremath{p_{T}}$
  - Measurements are sensitive to nuclear geometry and useful to constrain the theoretical models
- Baryon number carrier and its transport
  - Three approaches to test the carrier of baryon number & transport:
    - Isobar data: less electric-charge transport than baryon transport
    - Au+Au BES/global data: exponential rapidity dependence with slope showing no centrality dependence, flavor blind
    - Significant net-proton in γ+Au at midrapidity: exponential rapidity slope is compatible with the prediction
      of Regge theory on baryon junction
- First diffractive  $A_{\rm N}$  is studied, but diffractive  $A_{\rm N}$  can not have a significant contribution to large  $A_{\rm N}$
- STAR di- $\pi^0$  correlation study shows strong suppression at low  $p_T$  in p+A, following expected  $A^{1/3}$  dependence



# Backup slides

#### Results: characteristic features of y+Au events

Model calculations: Lewis et. al, arXiv: 2205.05685

— Peripheral Events: 60-80%

---- East-Going Photon Events

West-Going Photon Events

 $\gamma$ +Au-rich events



γ+Au events produce rapidity asymmetry that is expected from model predictions

Most photonuclear events have low multiplicity, consistent with very peripheral Au+Au collisions

(|η|<0.5)

30 N<sup>offline</sup>

STAR Preliminary Au+Au 54.4 GeV

Au+Au

20

Bulk features of  $\gamma$ +Au events are consistent with expectations from models

 $10^{-7}$ 

10'

10<sup>t</sup>

10

10<sup>2</sup>

10 🛓

10

Pevent N

## STAR forward upgrade

**Coverage:** 2.5 < η < 4.0 **Status:** 

- Installation and commission completed in 2021
- Start taking data since 2022

#### **Requirements:**

Detector	pp and pA	AA
ECal	$\sim$ 10 % / $\sqrt{E}$	$\sim$ 20 % / $\sqrt{E}$
HCal	$\sim$ 50 % / $\sqrt{E}$ + 10%	-
Tracking	Charge separation photon suppression	$\delta p_T/p_T \sim$ 20 - 30% for 0.2 $< p_T <$ 2 GeV/c

Measures:

- h<sup>+/-</sup>, e<sup>+/-</sup> (with good e/h separation)
- Photon, π<sup>0</sup>, jets



#### **Combines:**

- Forward Tracking System (FTS)
  - Forward Silicon Tracker (FST)
  - small-strip Thin Gap Chambers (sTGC)
- Forward Colorimeter System (FCS)
  - Electromagnetic Calorimeter (ECal)
  - Hadronic Calorimeter (HCal)

#### STAR: Forward tracking system (FTS)

Forward Tracking System (FTS):

- Forward Silicon Tracker (FST)
- small-strip Thin Gap Chambers (sTGC)



#### Forward Silicon Tracker (FST):

- 3 disks, each with 12 modules
- Each module includes 3 single-sided doublemetal mini-strip sensors (Si from Hamamatsu)
  - Fine granularity in  $\varphi$  and coarse in R
  - Material **budget ~ 1.5% X<sub>0</sub> per disk**

#### small-strip Thin Gap Chambers (sTGC):

- 4 planes, each with 5 pentagonal modules
  - Double-sided sTGC with diagonal strips give x, y, u in each layer
  - Position resolution < 200 µm
- Material budget  $\sim 0.5\% X_0$  per layer
- Readout based on VMM chips

## STAR: Forward calorimeter system (FCS)



#### ECal:

- Reuse PHENIX Pb-Scintillator calorimeter
  - <sup>•</sup> 1496 channels: 5.52 × 5.52 × 33 cm<sup>3</sup>
  - 18  $X_0$ ; 0.85 nuclear interaction length
- SiPM readout



#### HCal:

- Fe/Sc (20 mm/3 mm) sandwich
  - 520 channels: 10 × 10 × 84 cm<sup>3</sup>
  - $\sim$  4.5 nuclear interaction length
- Same SiPM readout as ECal

#### Physics opportunities with the STAR forward upgrade

#### Cold QCD:

- Sivers asymmetries for hadrons, (tagged) jets, and di-jets
- Collins measurements at high x
- GPD E<sub>g</sub> : gluon spin-orbit correlations
- Gluon PDFs for nuclei: R<sub>pA</sub> for direct photons & DY
- Test of Saturation predictions through dihadrons, γ-jets

Hot QCD:

- Temperature dependence of viscosity through flow harmonics up to  $\eta \sim 4$
- Longitudinal decorrelation up to  $\eta \sim 4$
- Global Lambda Polarization: test predictions
   of strong rapidity dependence

#### Physics opportunities with the STAR forward upgrade







No increase in the width of the azimuthal angular correlation is seen within experimental uncertainties

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