

Providing crustal parameters for the NS cooling simulations

Nikolai Shchechilin

With N. Chamel, A.I. Chugunov, M.E. Gusakov, A.Y. Potekhin, W. Ryssens, J.M. Pearson

Université Libre de Bruxelles



IReNA-INT Joint Workshop, Seattle 2024

On the way of providing crustal parameters for the NS cooling simulations

Nikolai Shchechilin

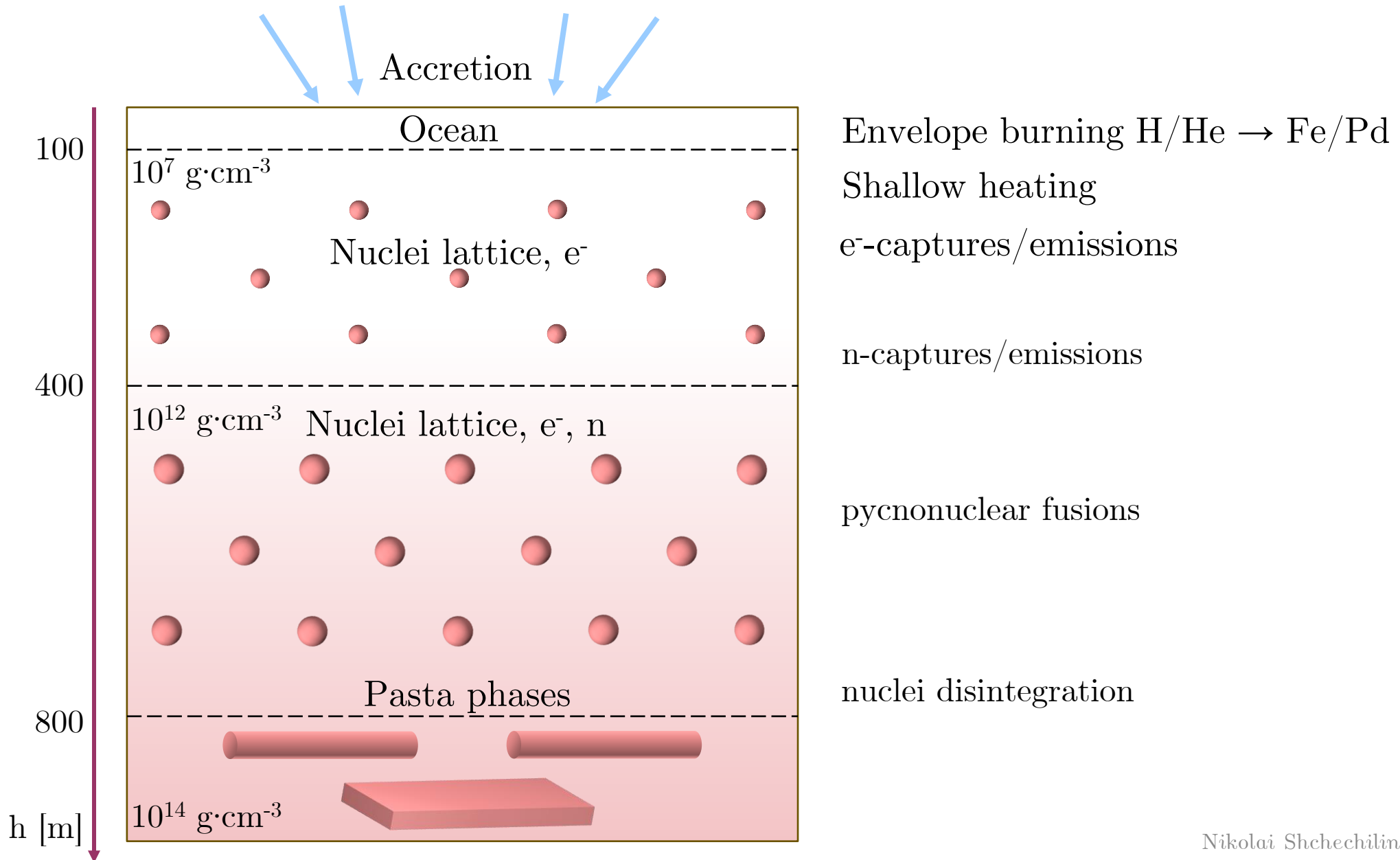
With N. Chamel, A.I. Chugunov, M.E. Gusakov, A.Y. Potekhin, W. Ryssens, J.M. Pearson

Université Libre de Bruxelles



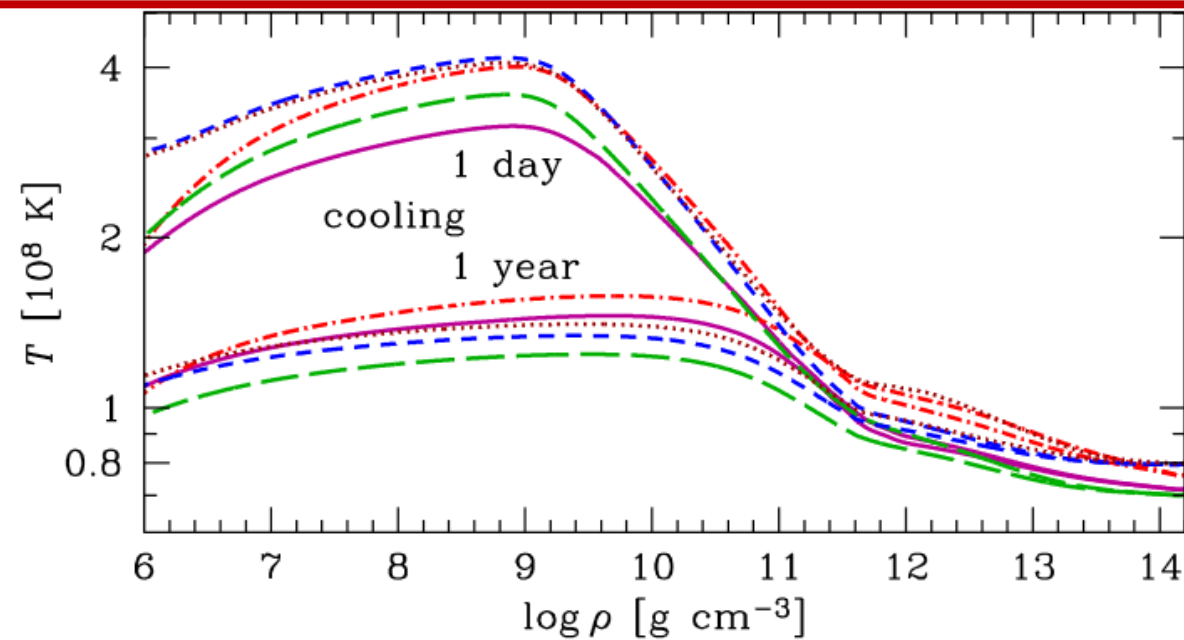
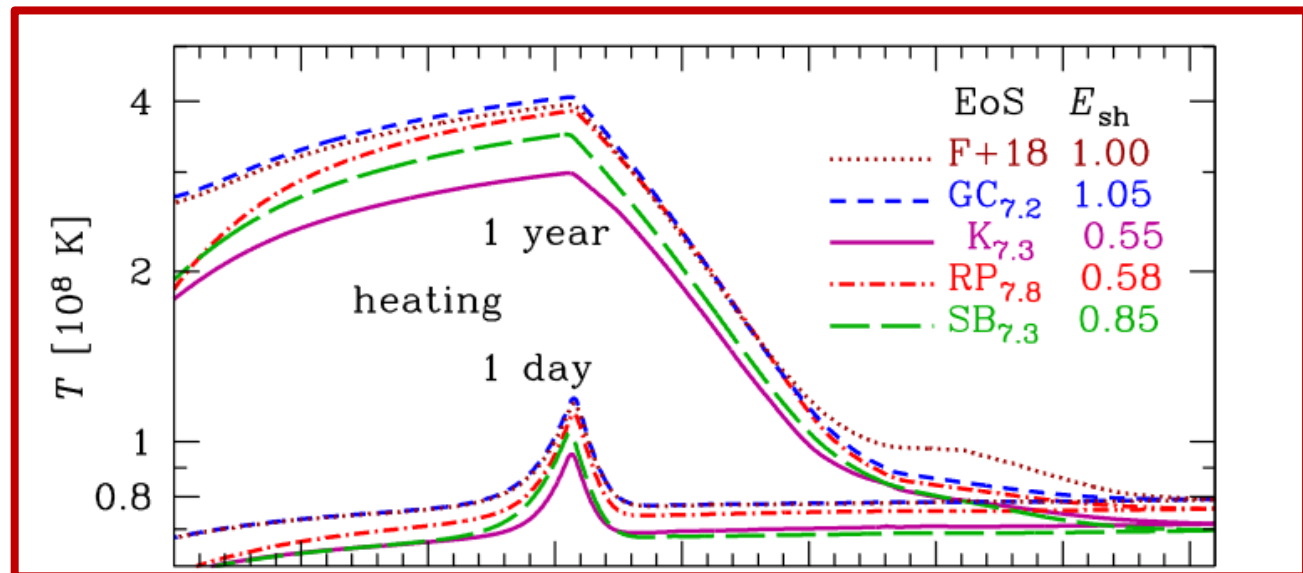
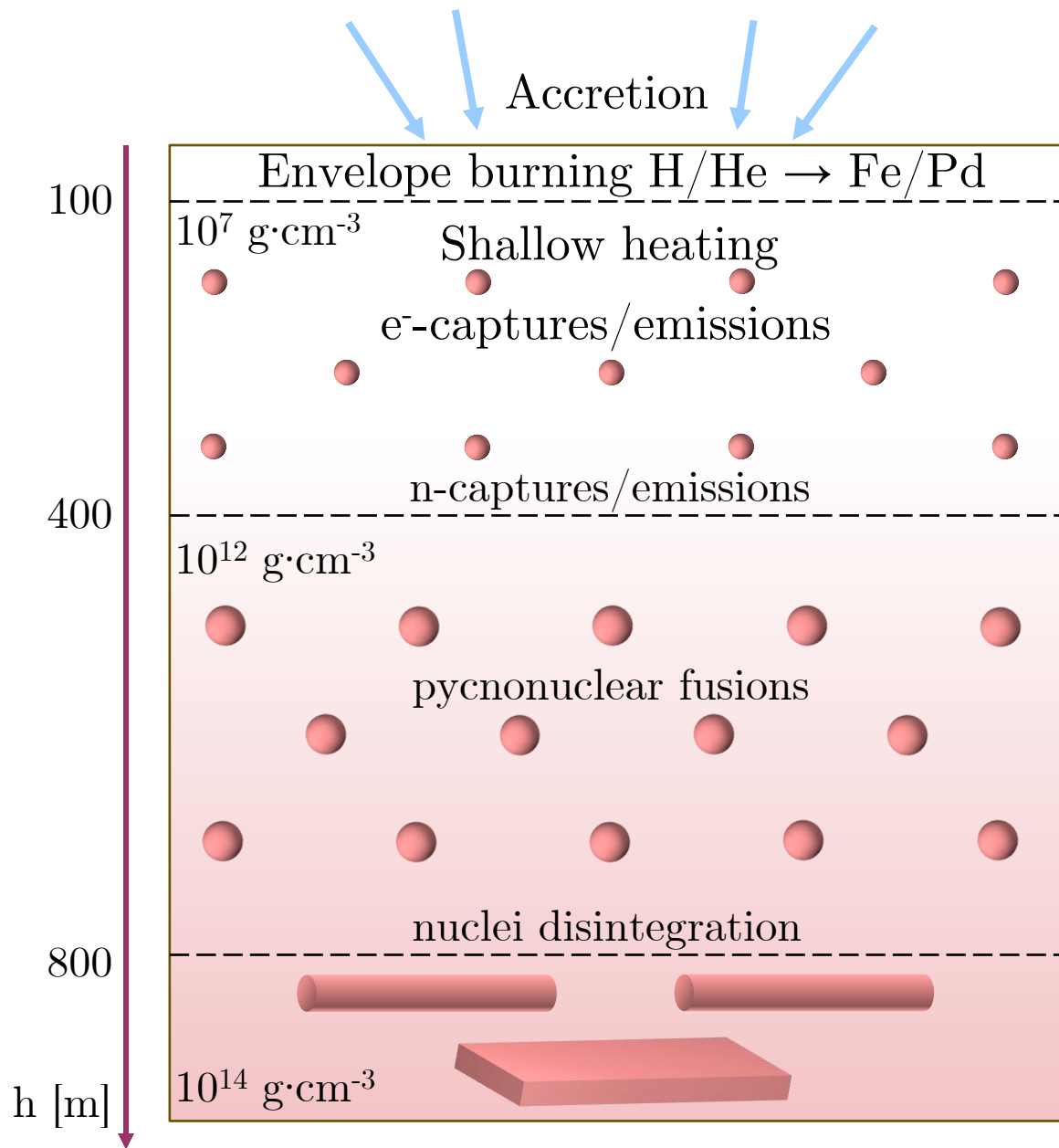
IReNA-INT Joint Workshop, Seattle 2024

General paradigm of accreted crust



Crustal thermal evolution

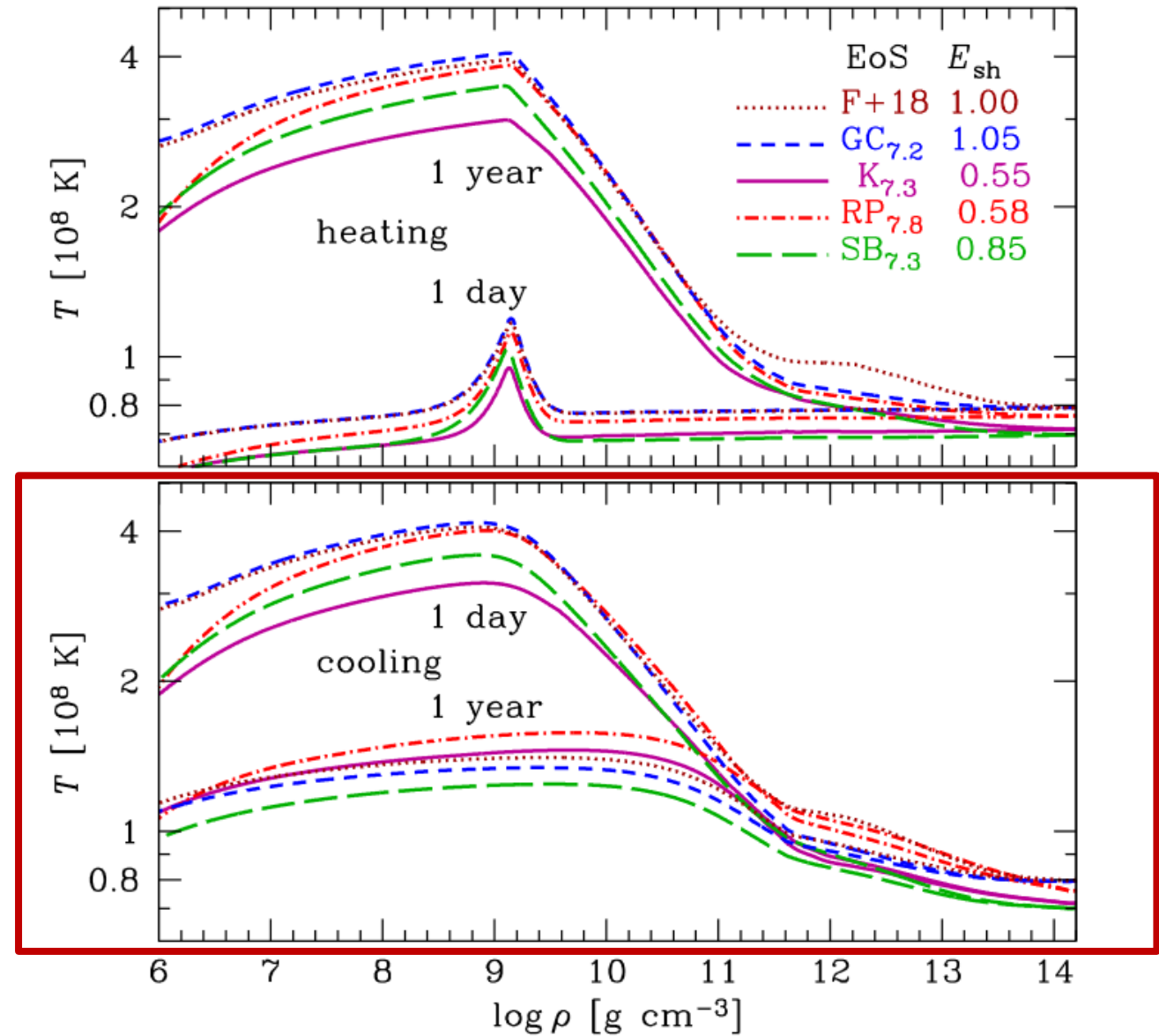
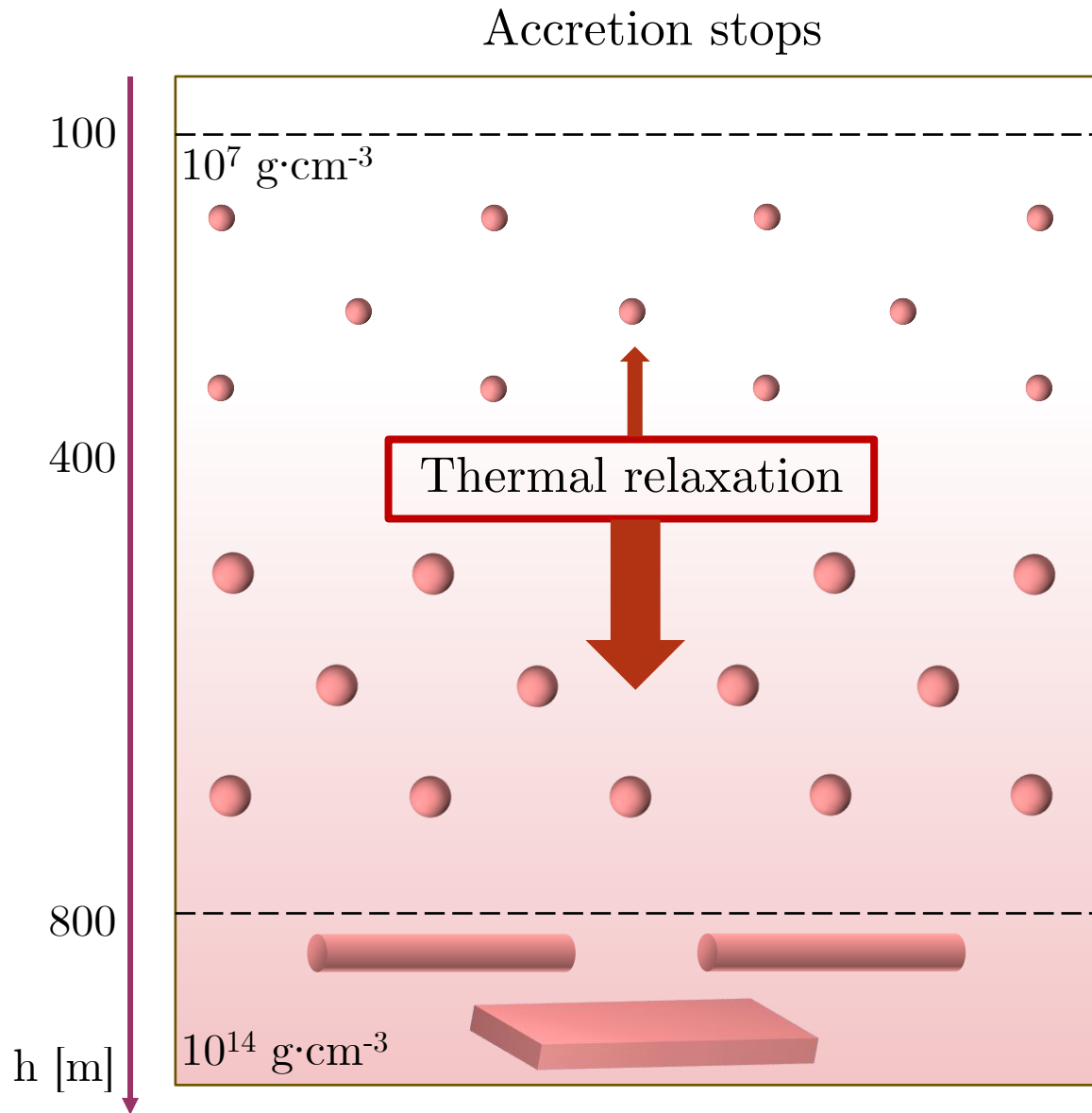
Temperature profiles in the crust



[Potekhin, Chugunov, NS, Gusakov '24]

Crustal thermal evolution

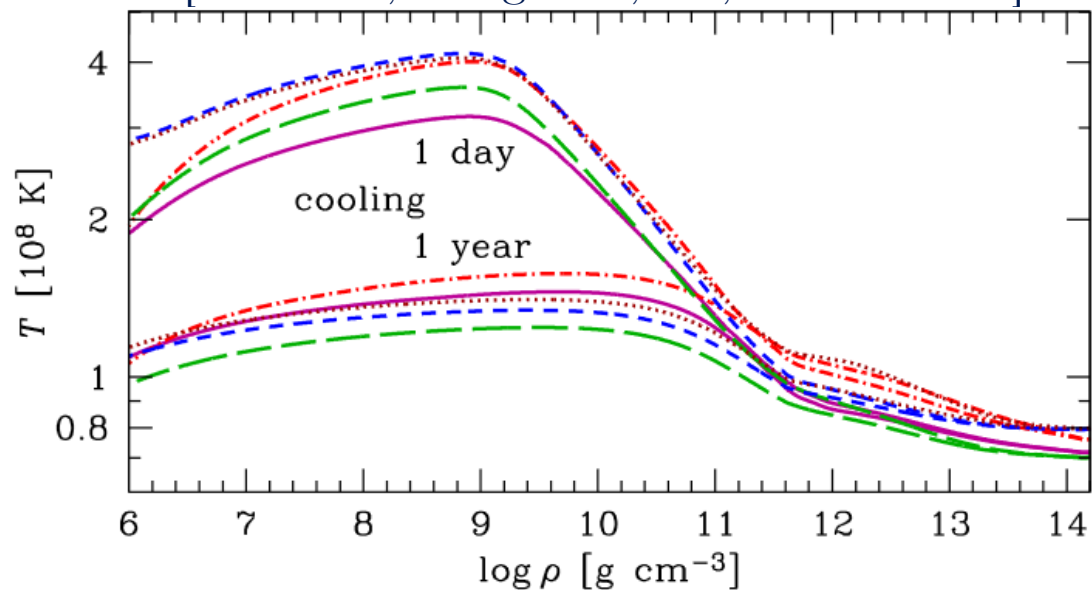
Temperature profiles in the crust



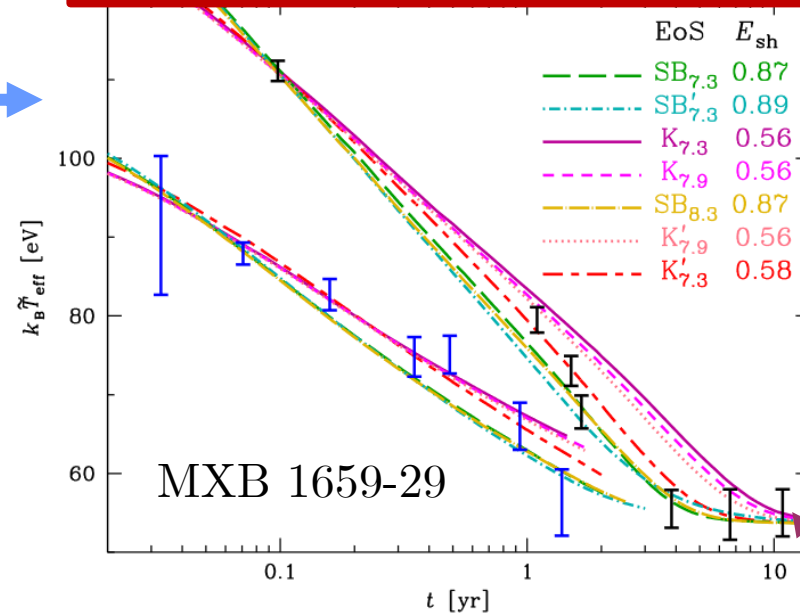
[Potekhin, Chugunov, NS, Gusakov '24]

Crustal thermal evolution

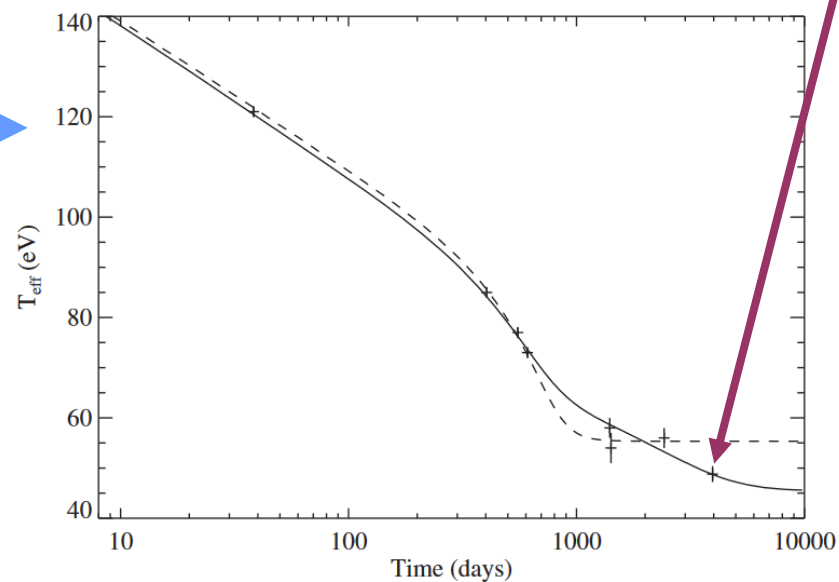
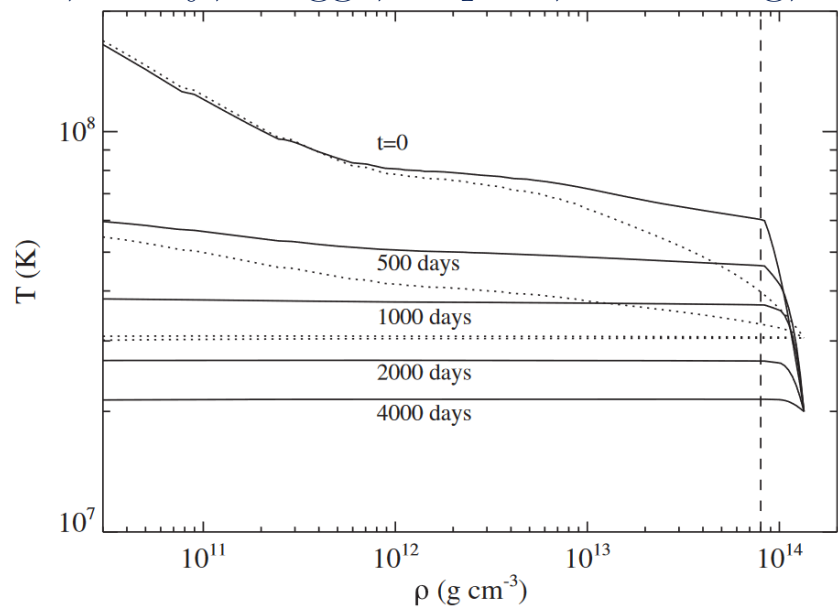
[Potekhin, Chugunov, NS, Gusakov '24]



Cooling curves vs observations



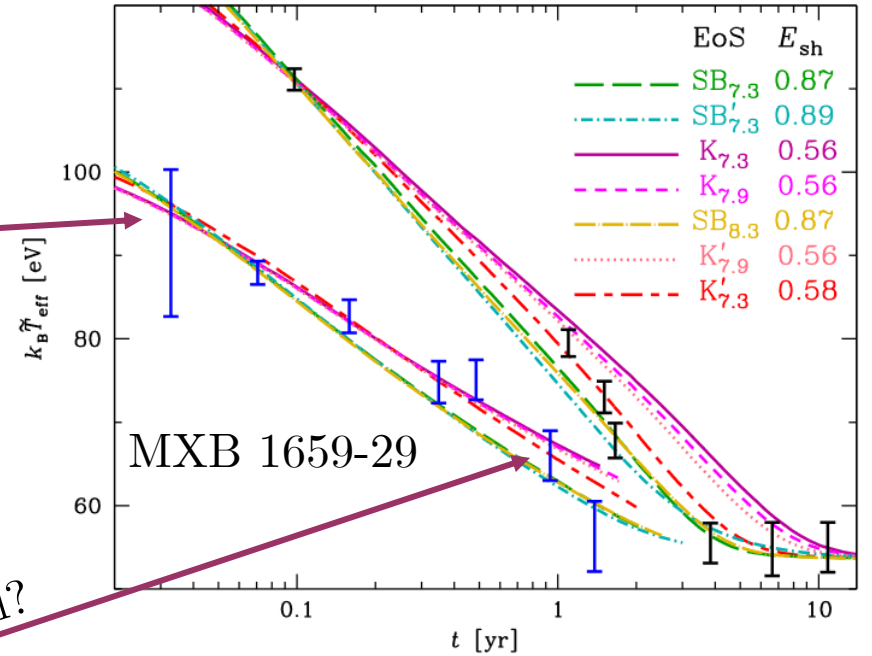
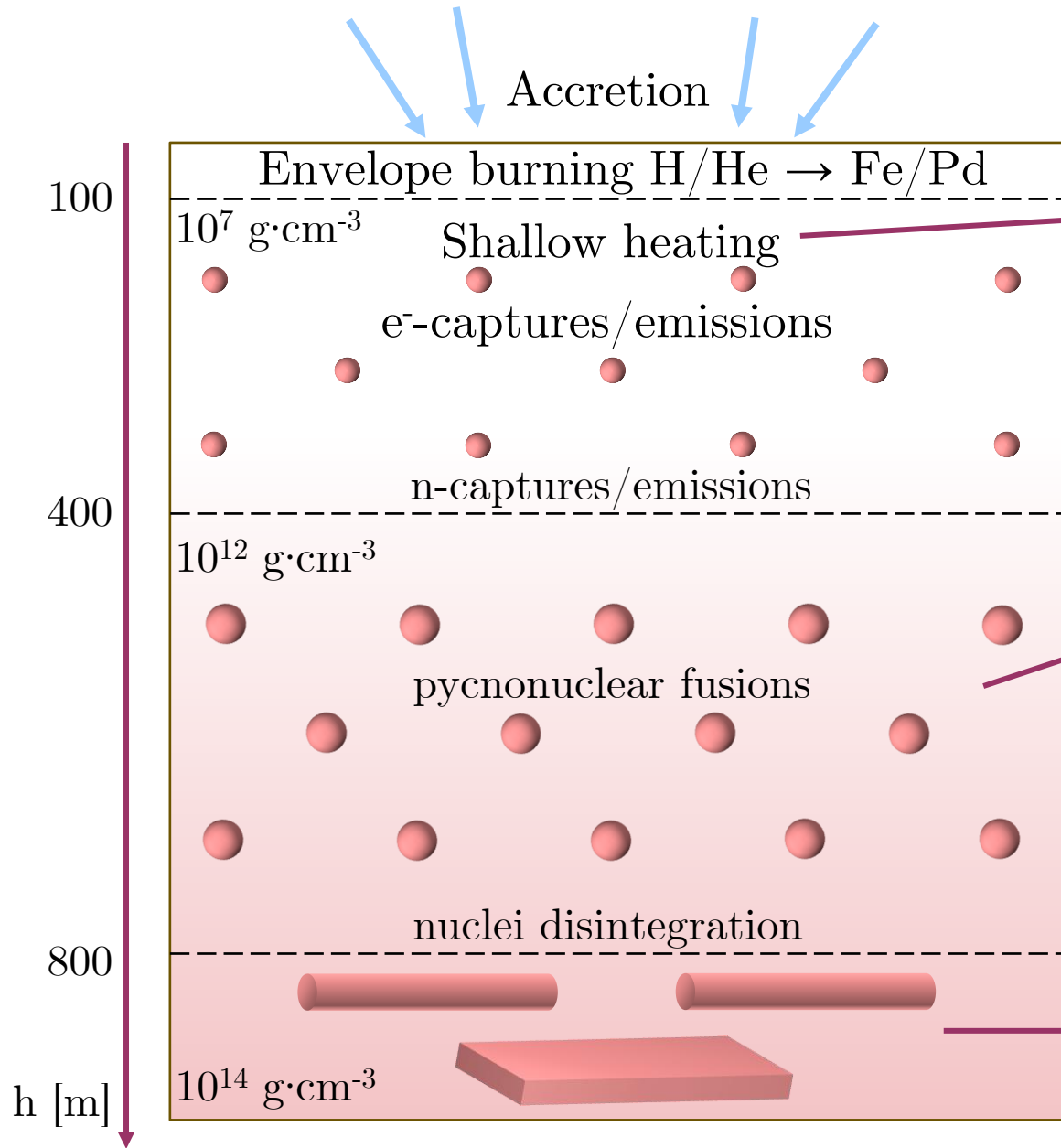
[Horowitz, Berry, Briggs, Caplan, Cumming, Schneider '15]



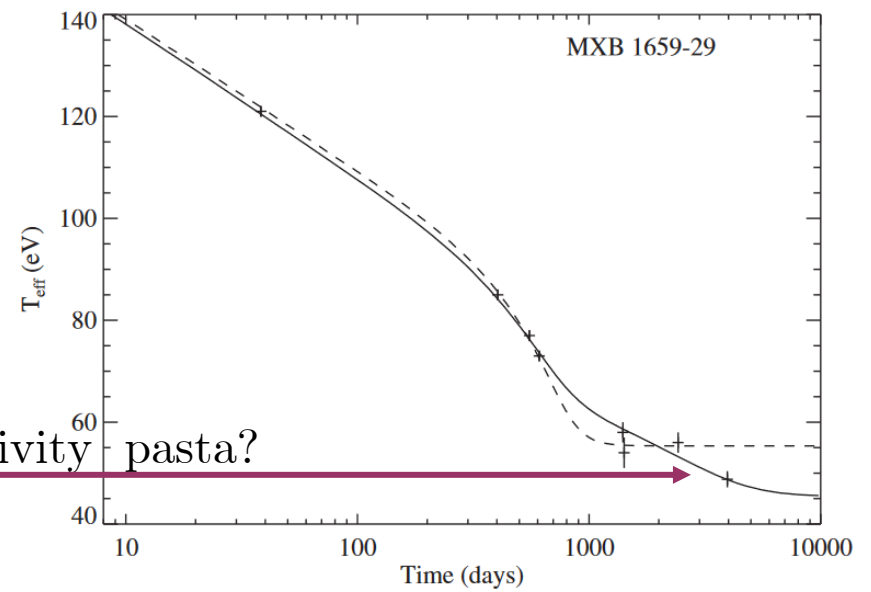
Depends on the spectral model [Cackett+ '13]

Crustal imprints

[Potekhin, Chugunov, NS, Gusakov '24]



Pure crystal?

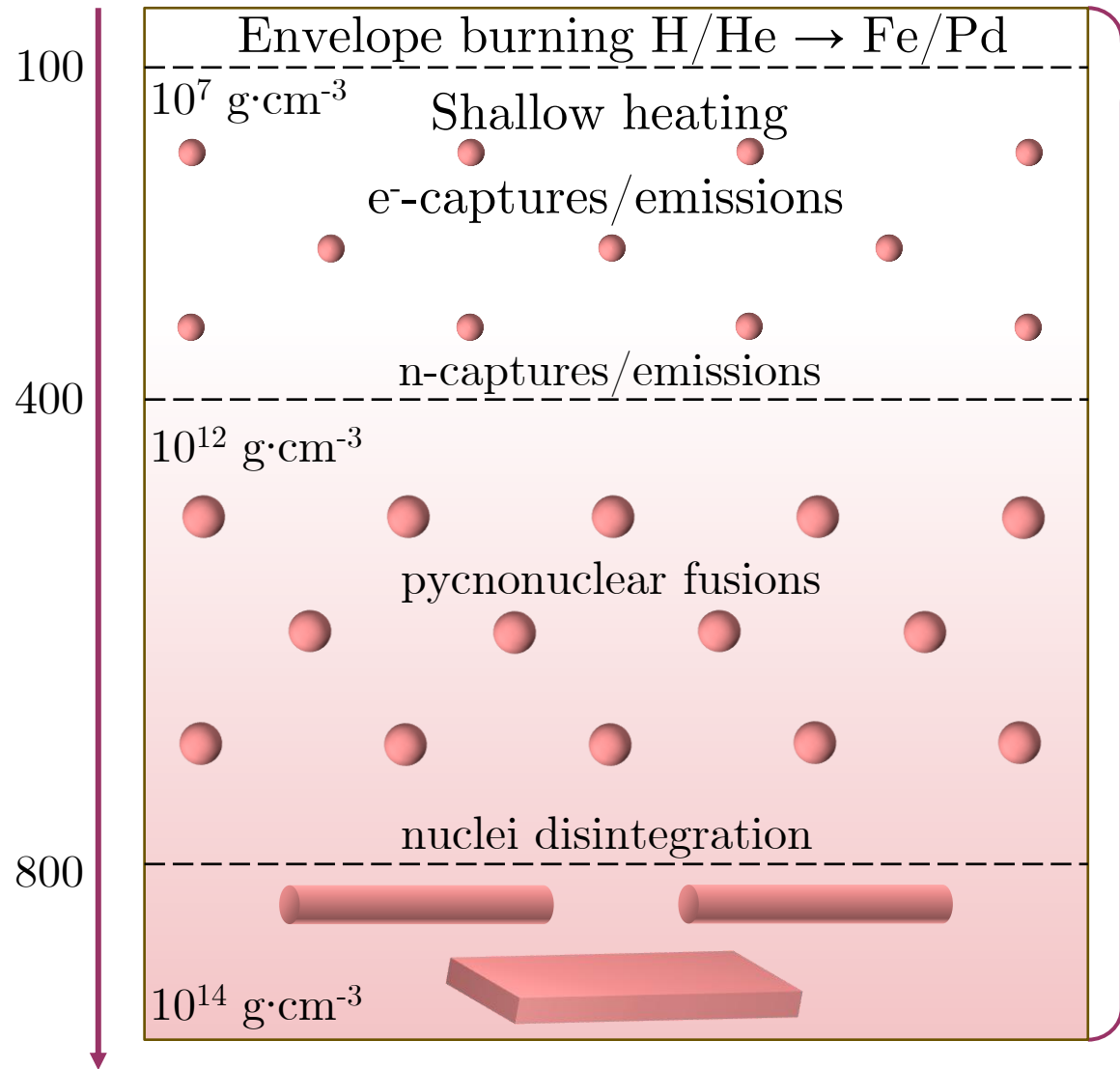


Low-conductivity pasta?

[Horowitz, Berry, Briggs, Caplan, Cumming, Schneider'15]

Crustal inputs for cooling simulations

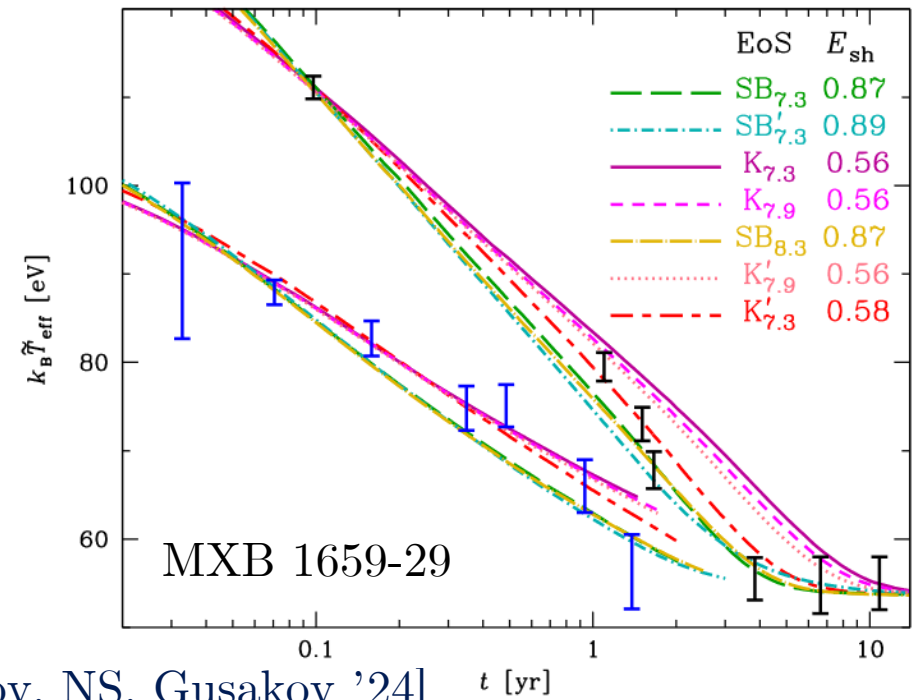
Accretion



Heating sources, composition, superfluidity model,

$\langle Z \rangle$, Q_{imp} , Δ_n

heat capacity, conductivity



[Potekhin, Chugunov, NS, Gusakov '24]

t [yr]

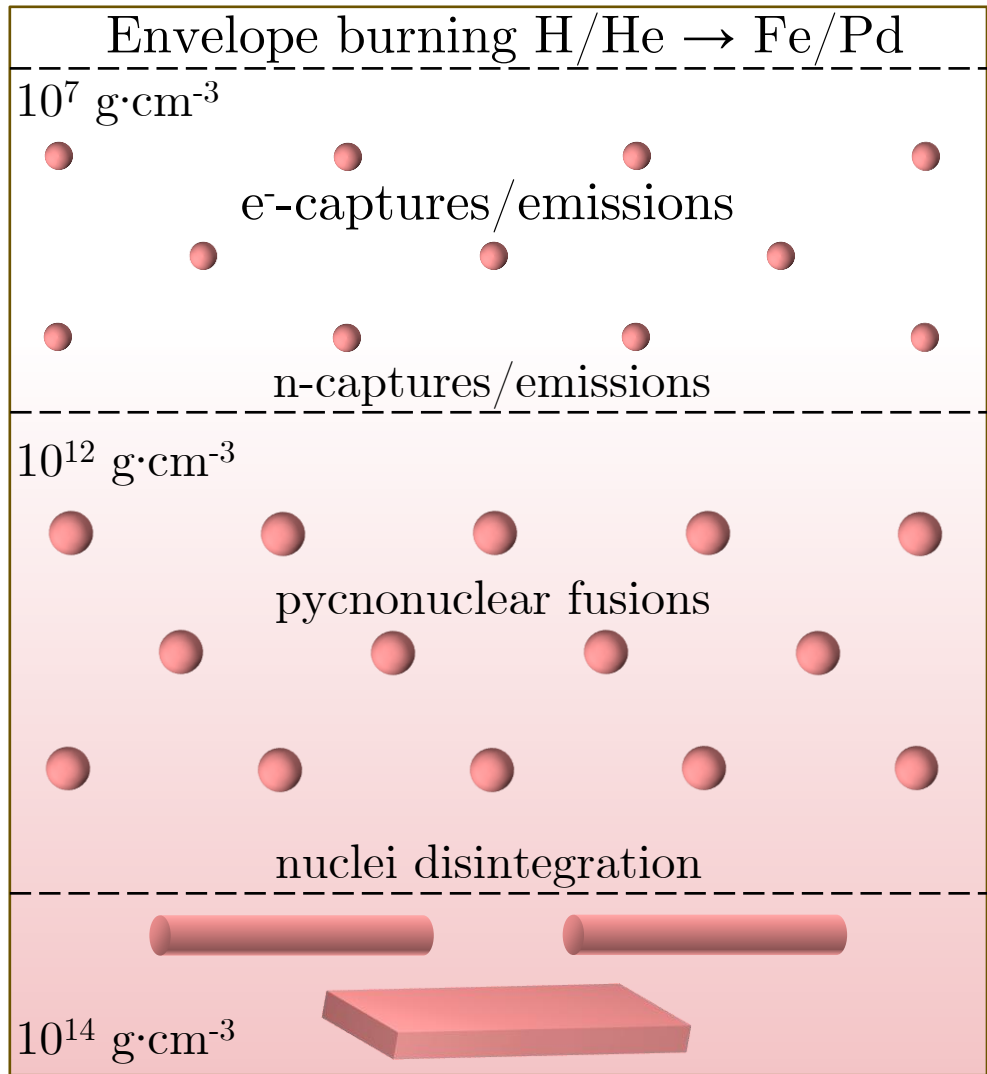
Total heating: general considerations

$$Q_{\text{tot}}^{\infty} = m_{\text{b,H}} e^{\nu_s/2} - \mu_{\text{b,core}}^{\infty} =$$

redshift factor

$$m_{\text{b,H}} e^{\nu_s/2}$$

Accreted material conversion



$$(m_{\text{b,H}} - m_{\text{b,ash}}) e^{\nu_s/2}$$

+

$$(m_{\text{b,ash}} e^{\nu_s/2} - \mu_{\text{b,core}}^{\infty})$$

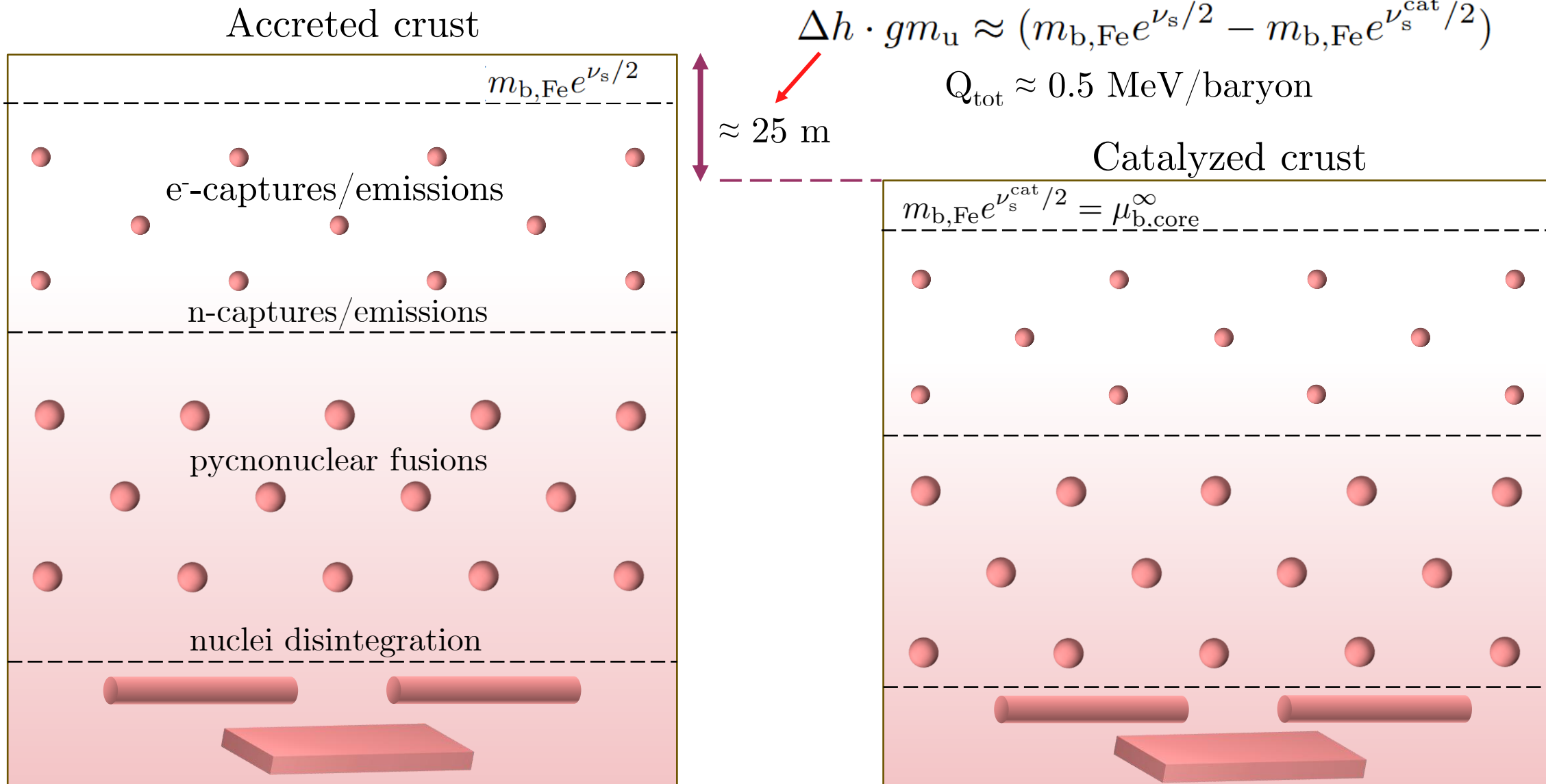
Heating in a “net” reaction

[Gusakov & Chugunov '21]

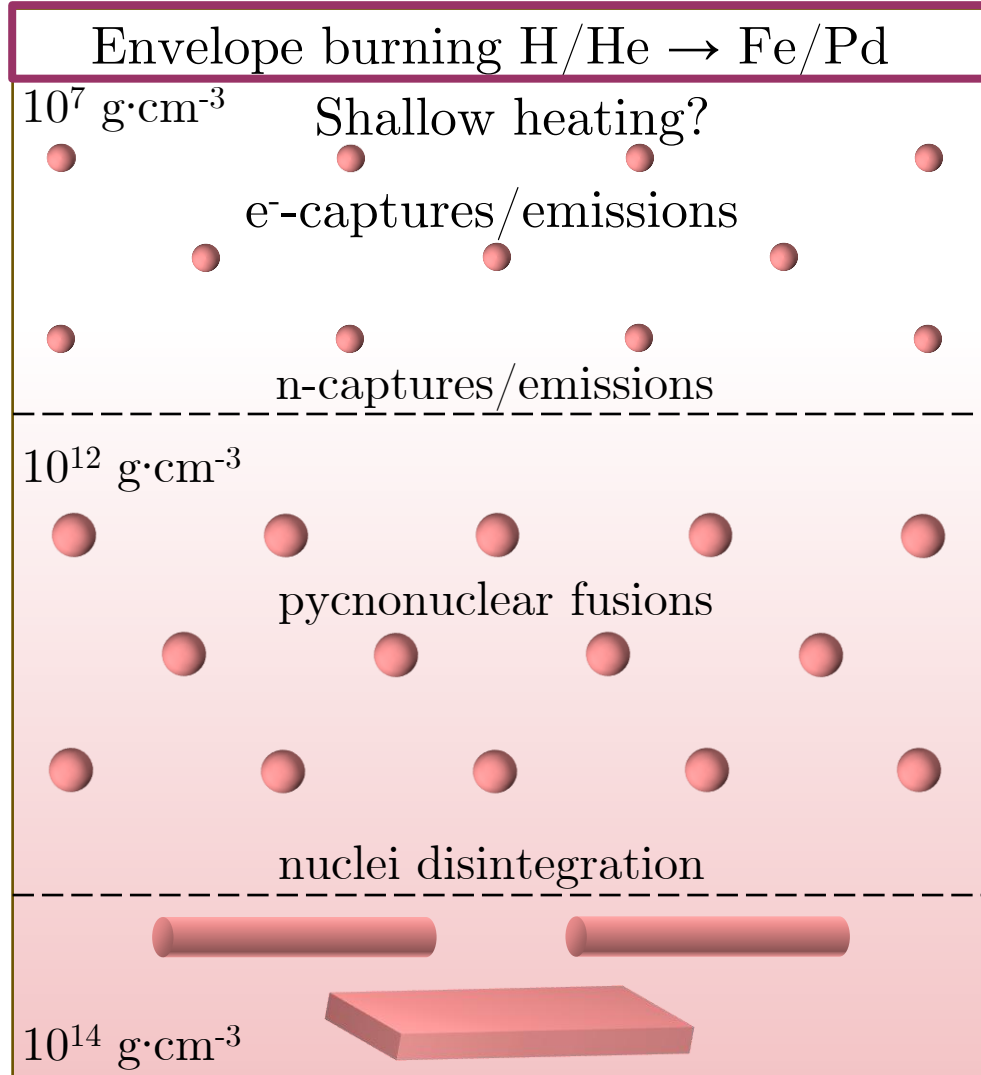
$$\mu_{\text{b,core}}^{\infty}$$

Total heating: general considerations

[Gusakov & Chugunov '21]

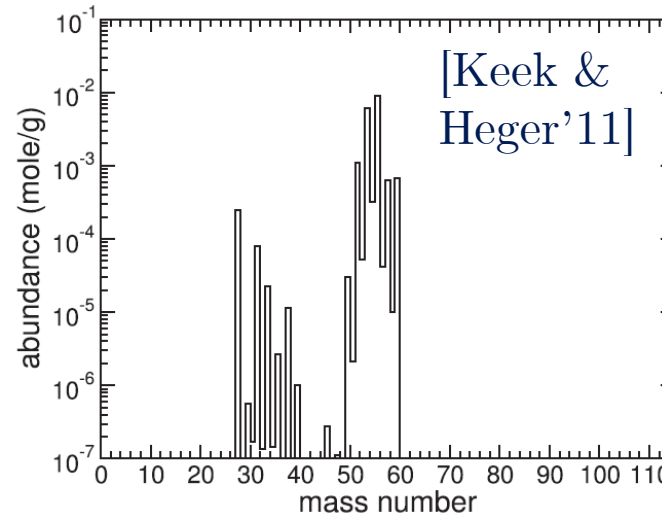


Initial composition for the crustal simulations

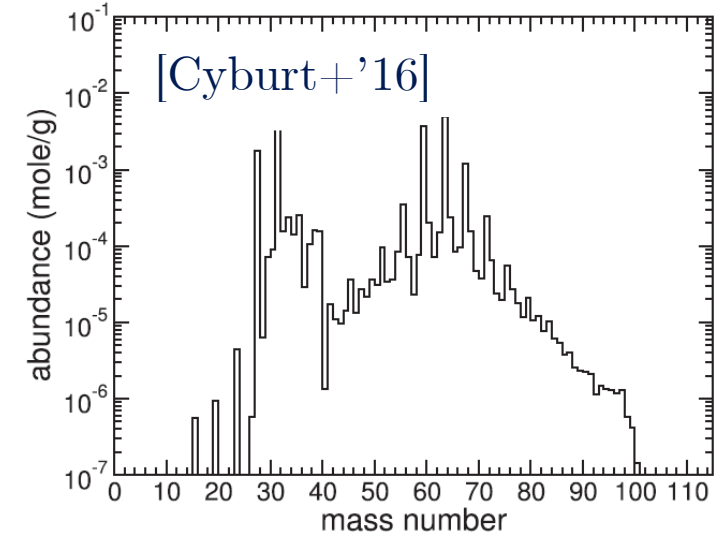


Three models of thermonuclear ashes

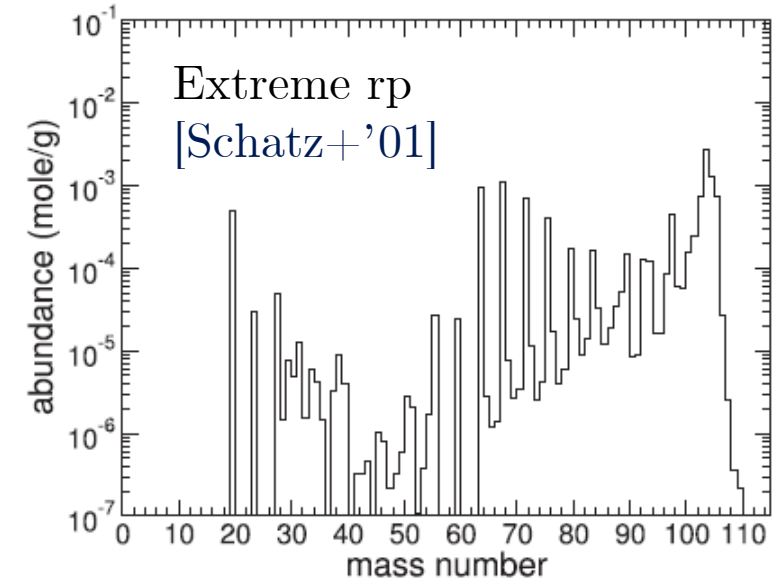
Superburst



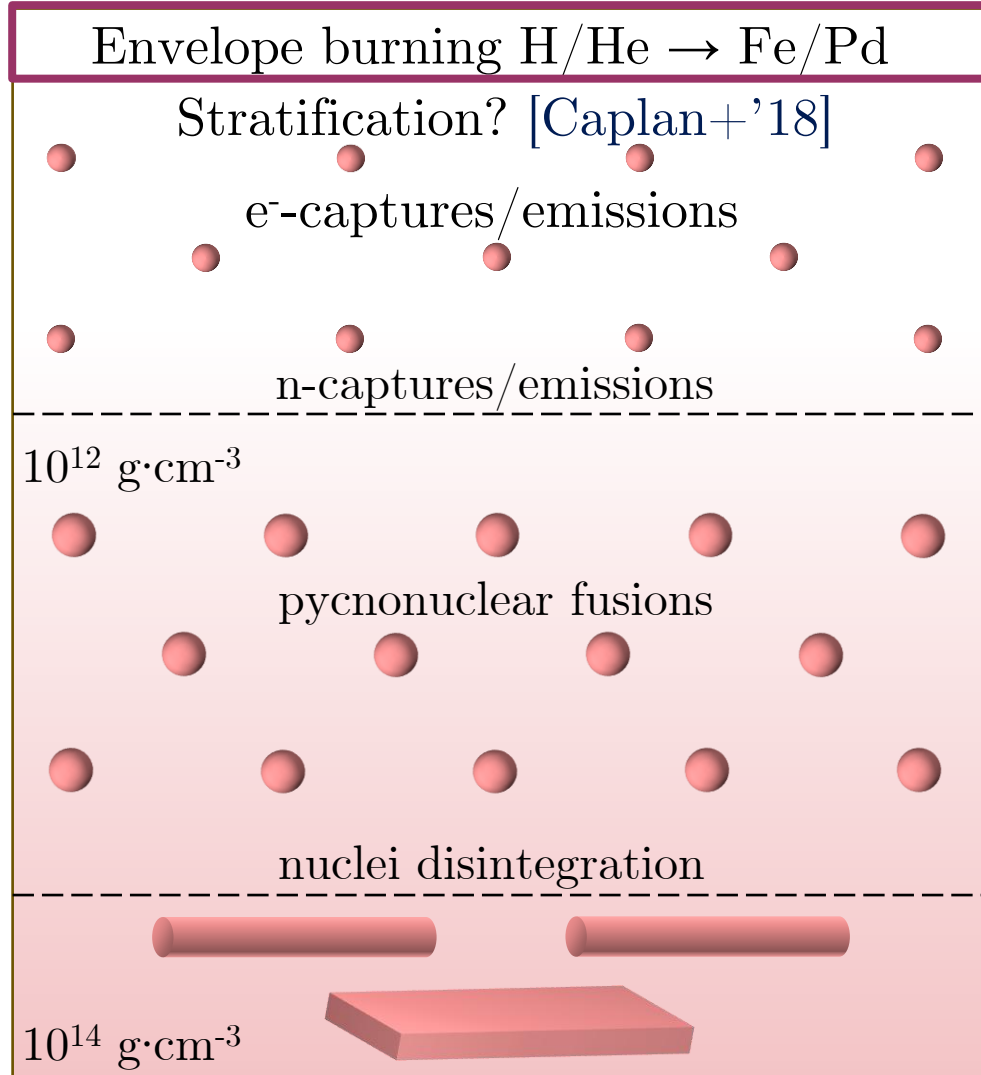
Kepler X-ray burst



Talks of
 D. Galloway
 H. Schatz
 D. Page

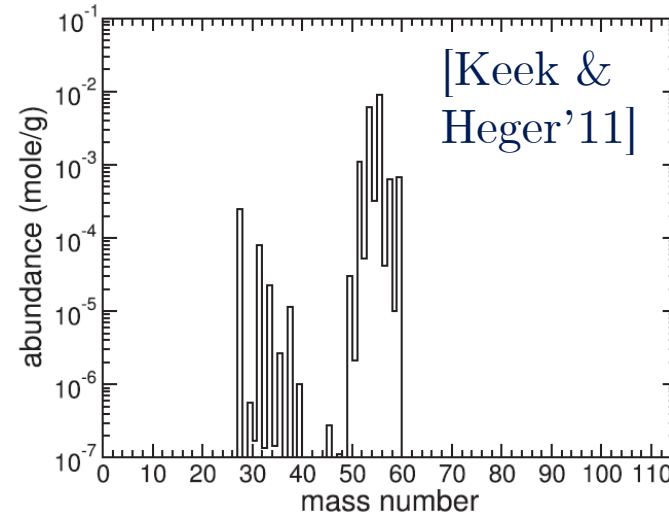


Initial composition for the crustal simulations

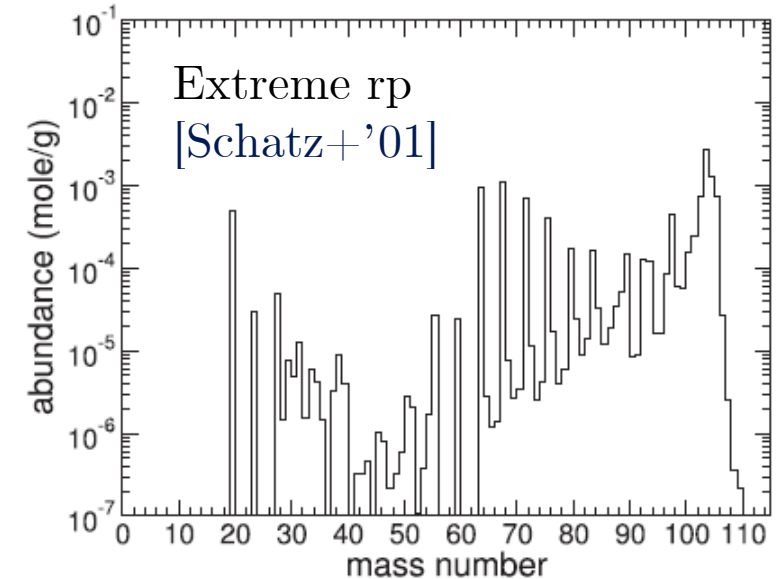
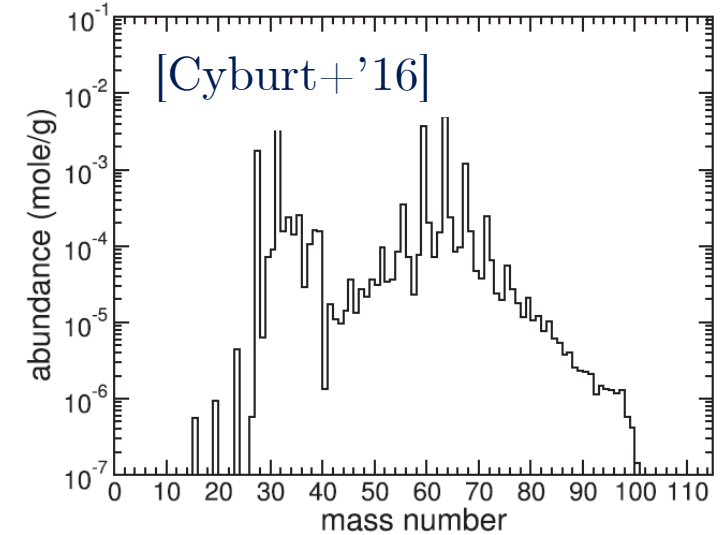


Three models of thermonuclear ashes

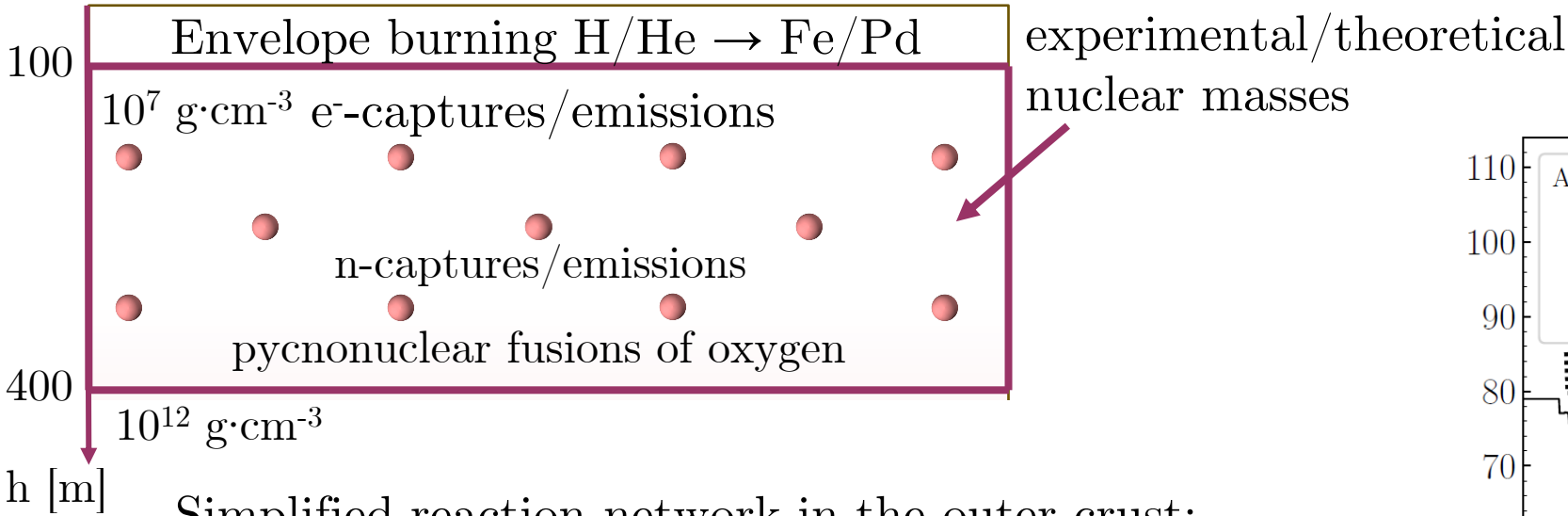
Superburst



Kepler X-ray burst

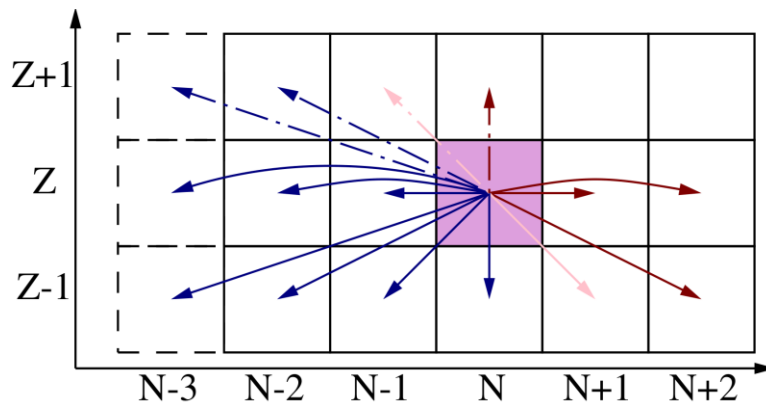


Evolution in the outer crust

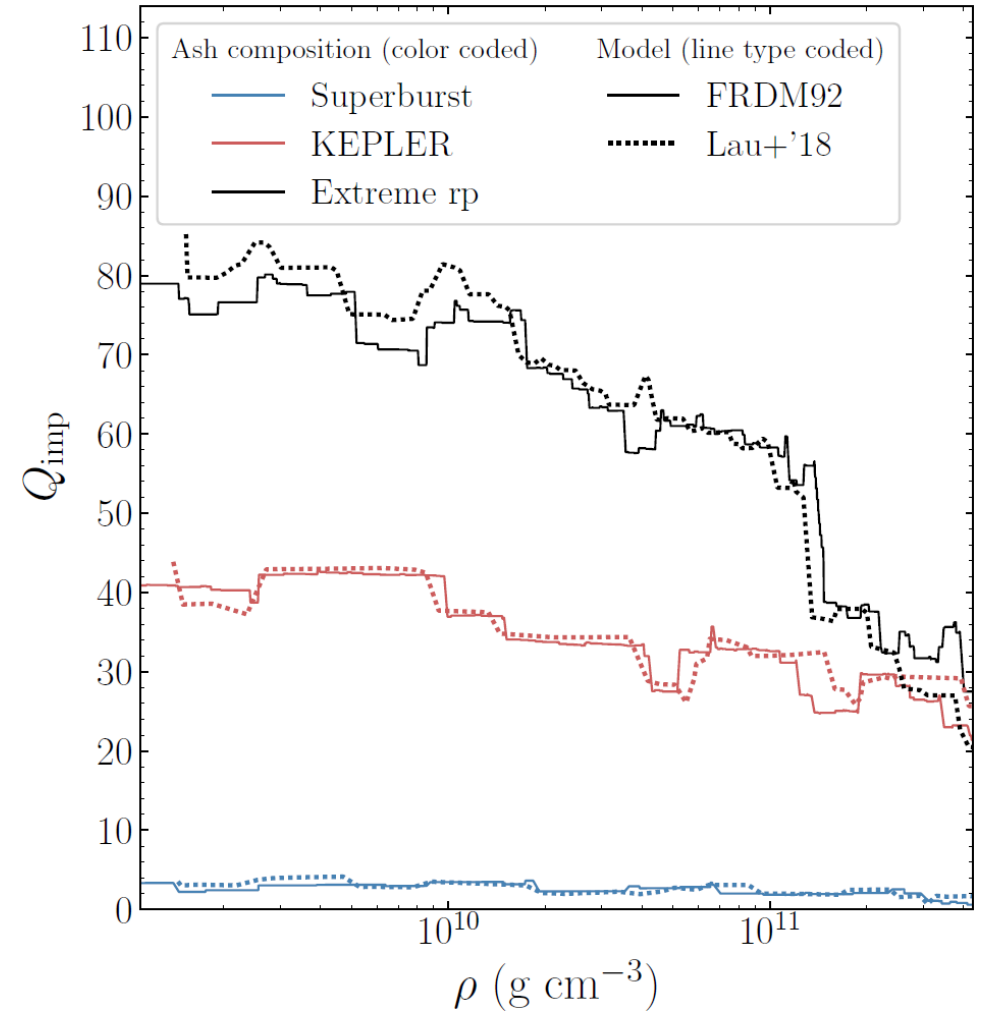


Simplified reaction network in the outer crust:

- Check for allowed reactions (min Gibbs energy)
- Small part of nuclei reacting according to the priority rules (based on τ_{reac})
- Increase the pressure



Comparison with [Lau+'18]



[NS, Gusakov, Chugunov'21]

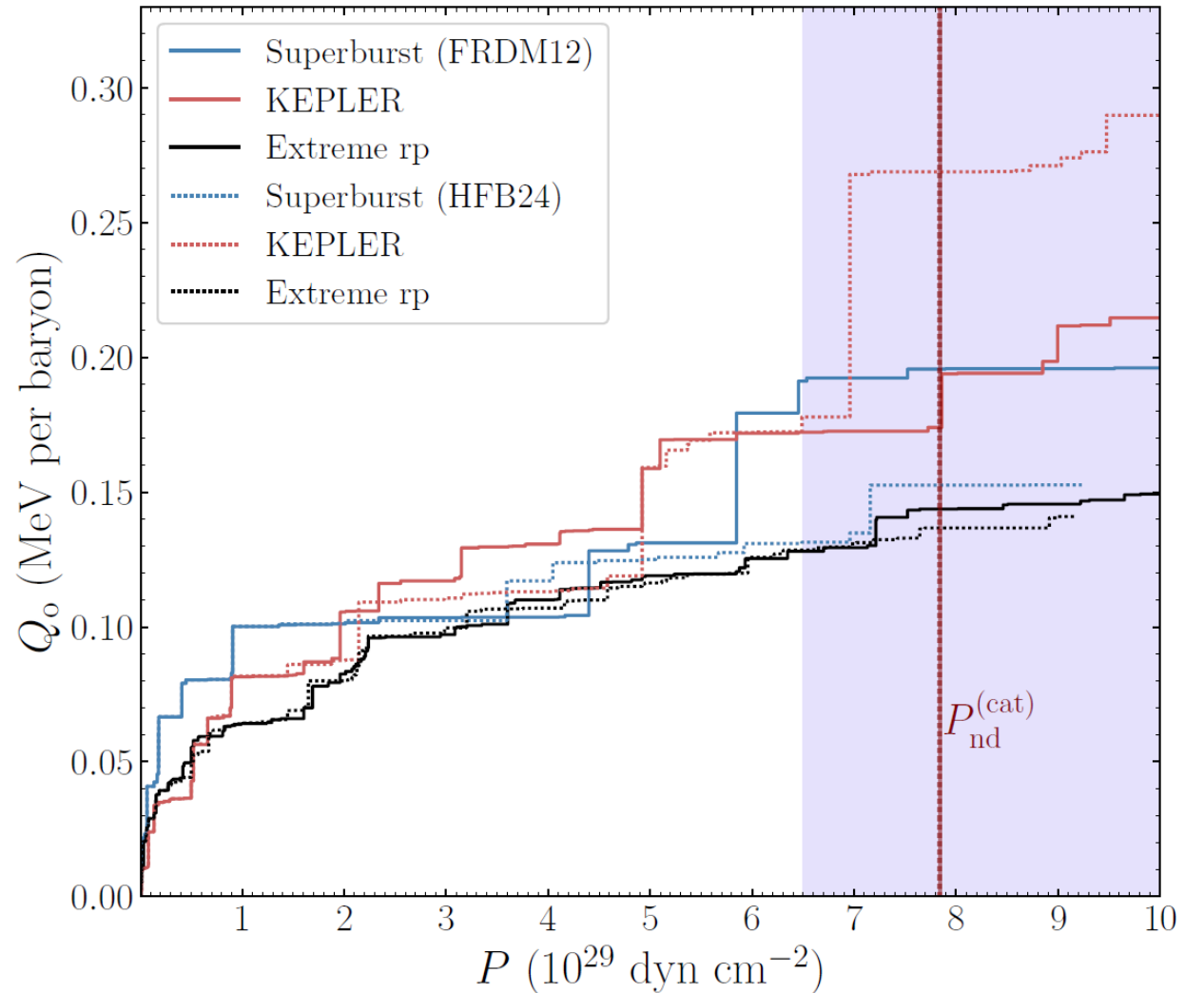
Heating in the outer crust

Total accumulated heat in the **outer crust** with experimental masses AME2020 plus HFB24/FRDM12 theoretical mass evaluations

- Superburst $Q_o \approx 0.15-0.20$ MeV/baryon
- Kepler $Q_o \approx 0.17-0.26$ MeV/baryon
- Extreme rp $Q_o \approx 0.12-0.15$ MeV/baryon

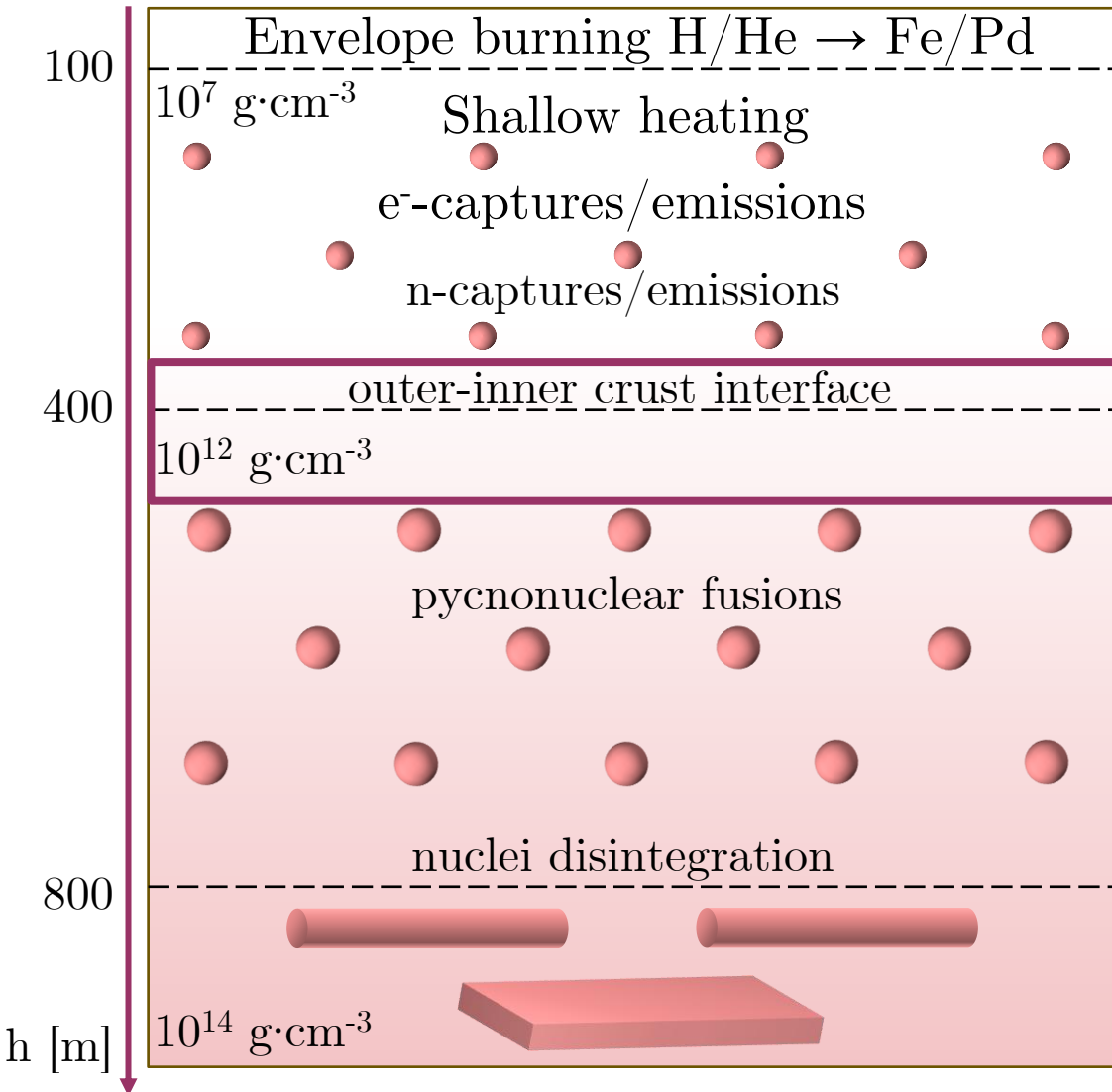
Without neutrino cooling and shallow heating!

Heat distribution in the outer crust



[NS, Gusakov, Chugunov'21]

Transition to the inner crust



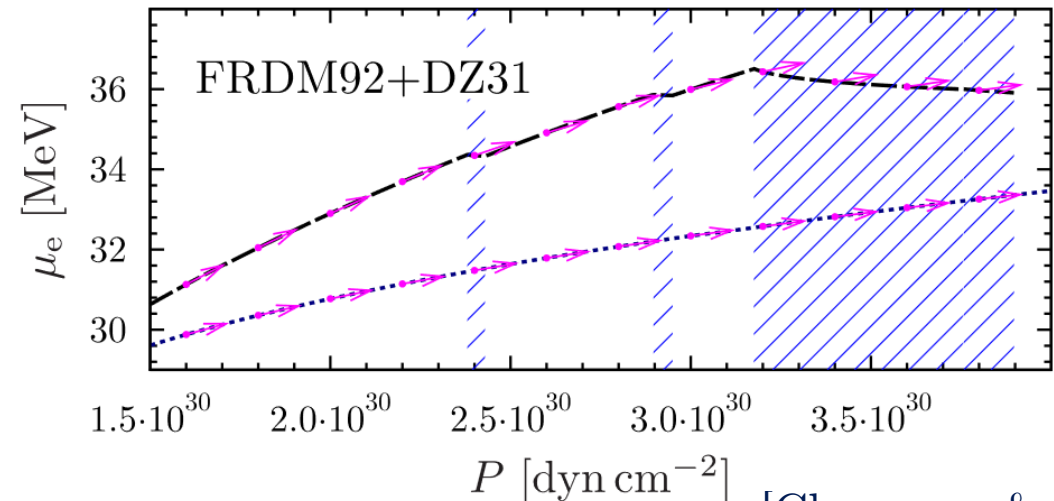
Traditional approach

- Compression till neutrons become abundant (not captured by the neighboring nuclei)
- Keeping neutrons in the compressed volume

Such crust can not exist!

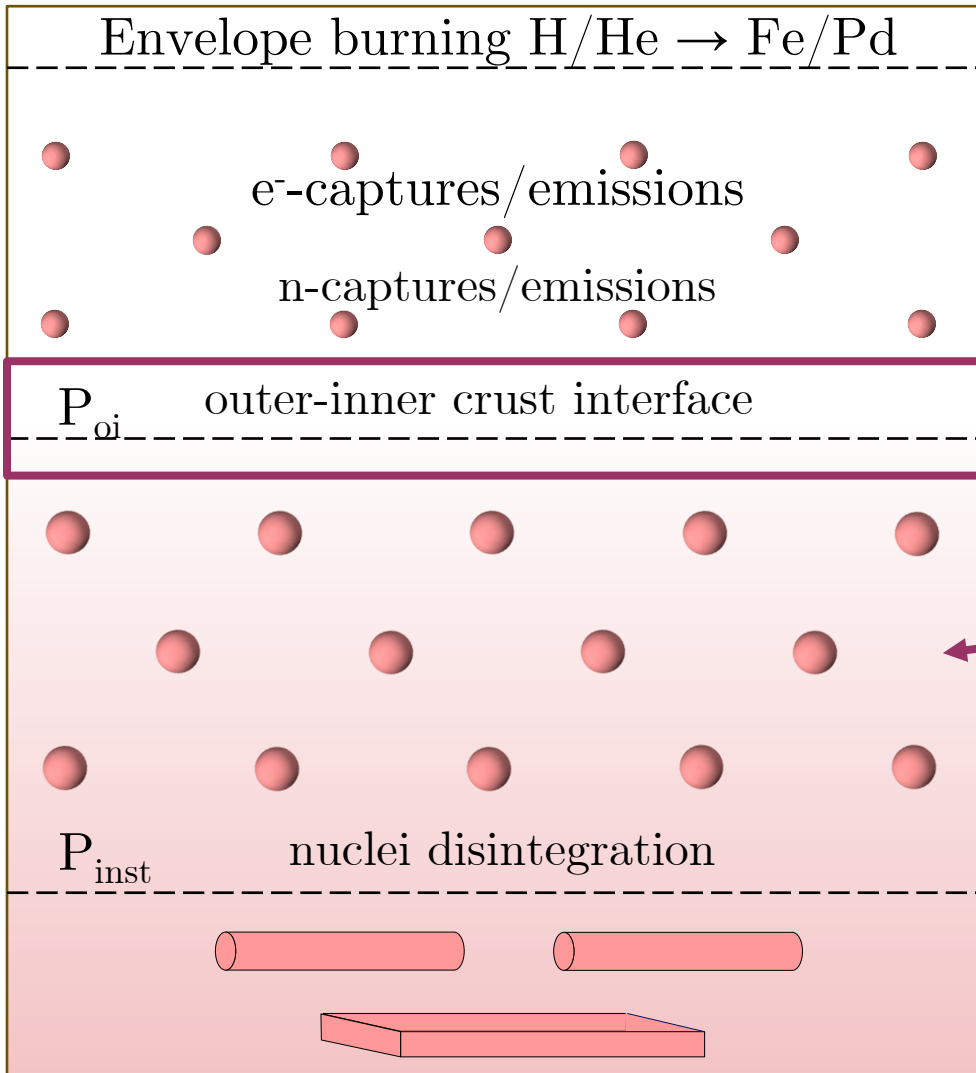
In the regions of active neutron emission electron chemical potential decreases violating the force balance

$$Z \nabla \mu_e \downarrow - m_i g \neq 0$$



[Chugunov & NS'20]

Transition to the inner crust



Neutron Hydrostatic and Diffusion (nHD) equilibrium superfluidity/rapid diffusion allows considering static picture in the limit of fast redistribution of neutrons

$$\mu_n^\infty = \text{const}$$

$$\mu_n^\infty = m_n e^{\nu_{oi}/2}$$

Unbound neutrons form a 'sea'

On course of accretion nuclei sink in the sea of neutrons

Thermodynamic analyses [Gusakov & Chugunov'24] shows

$$P_{oi} \lesssim P_{nd}^{(cat)}$$

$$\mu_n^\infty = \mu_{b,core}^\infty$$

Evolution in the upper layers of the inner crust

Nuclei sink in the sea of neutrons

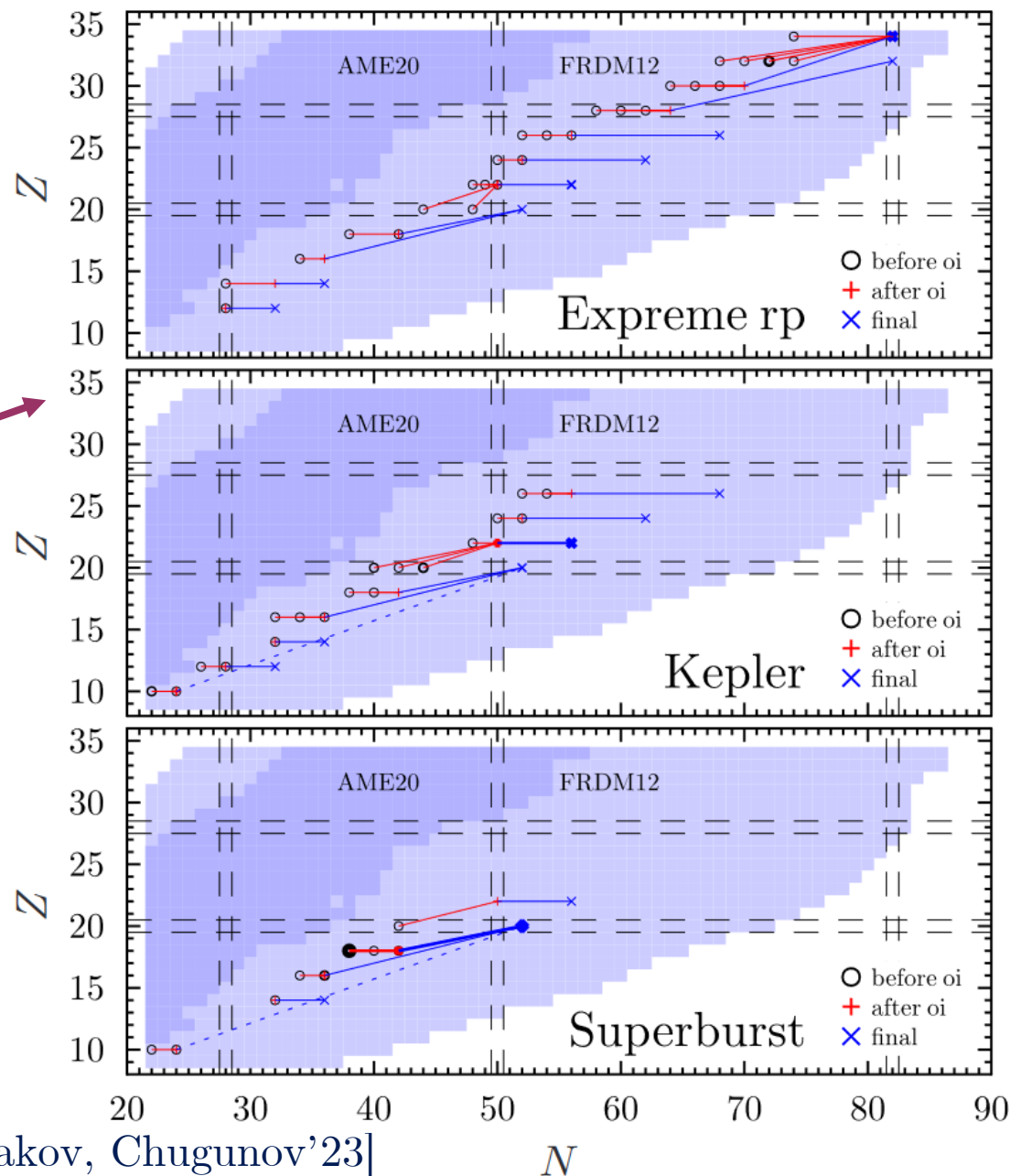
Main reactions at the oi interface:

- neutron captures
- electron emissions

In this shallow layers nuclei mass tables are still applicable [Lau+'18]

Catalyzed crust nuclei have $Z \approx 40$

In contrast to the traditional approach,
for which reverse reactions dominate



Evolution in the upper layers of the inner crust

Simplified reaction network in the inner crust:

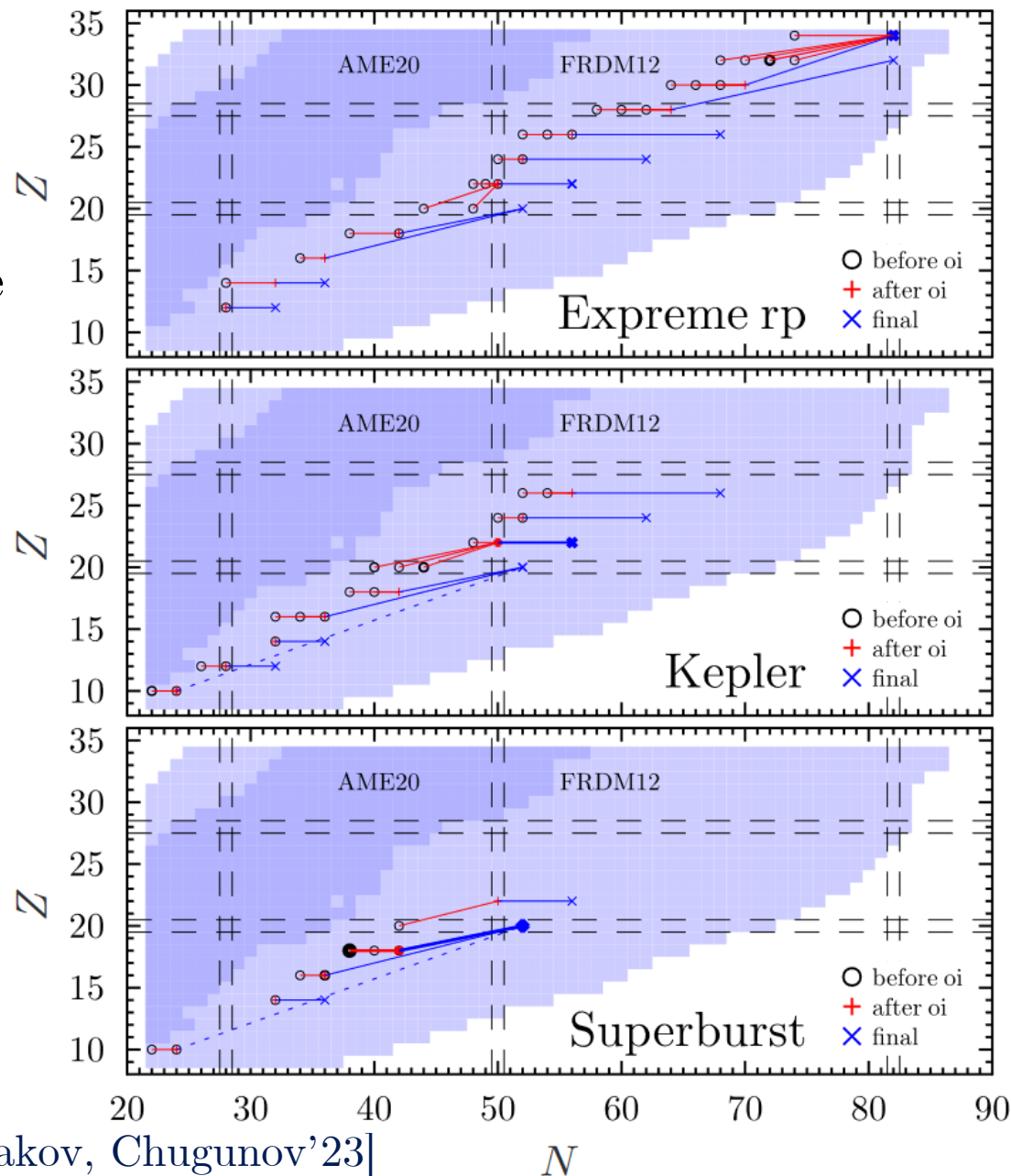
- Check for allowed reactions, which minimize $\Psi = G - \mu_n N_b$ [Gusakov, Kantor, Chugunov'21]
- Small part of nuclei reacting according to the priority rules (based on τ_{reac})
- Increase the pressure and μ_n

Main reactions at the oi interface:

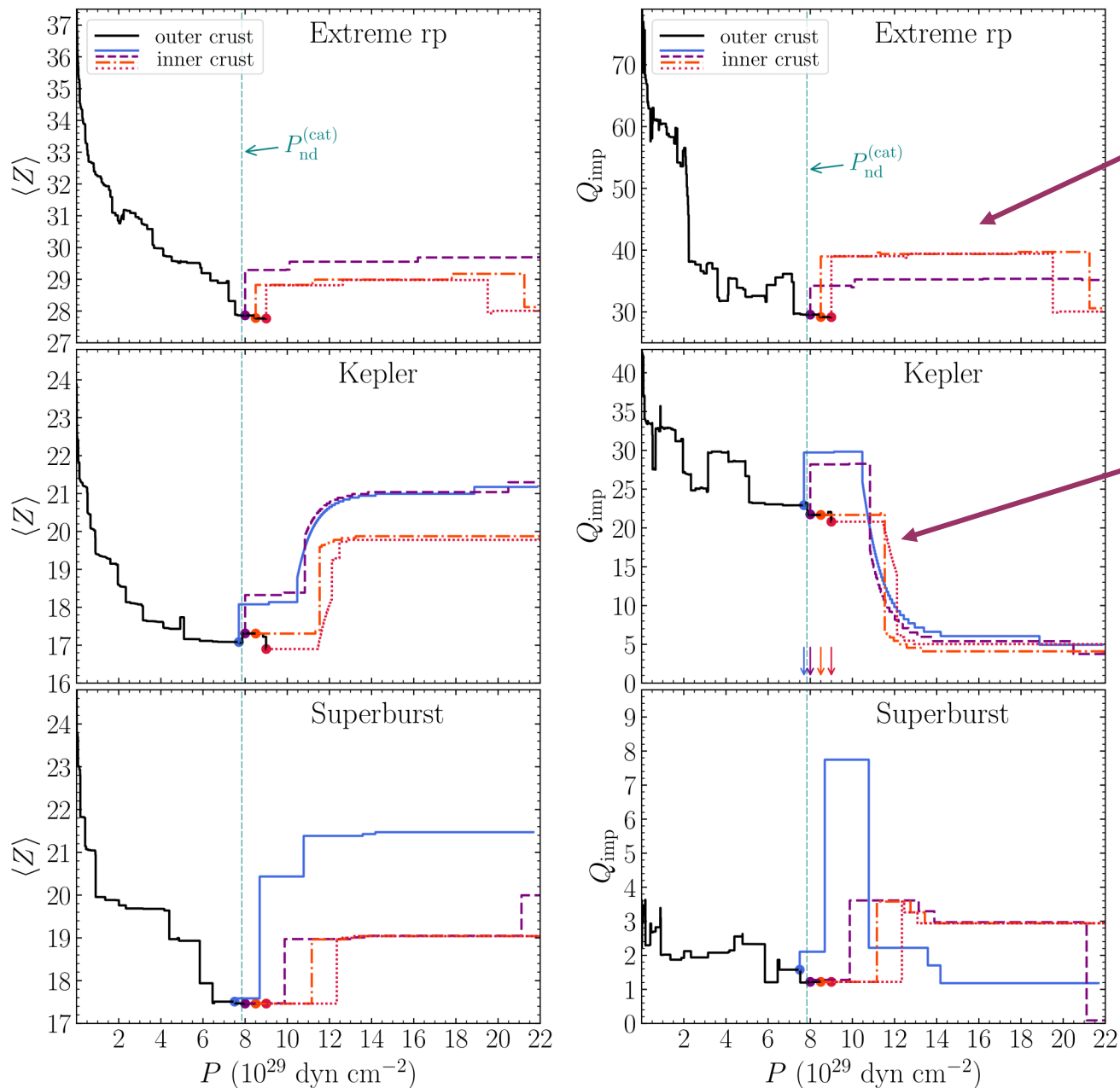
- neutron captures
- electron emissions

Pycnonuclear fusions play some role only for Kepler ashes

In contrast to the traditional approach, for which reverse reactions dominate



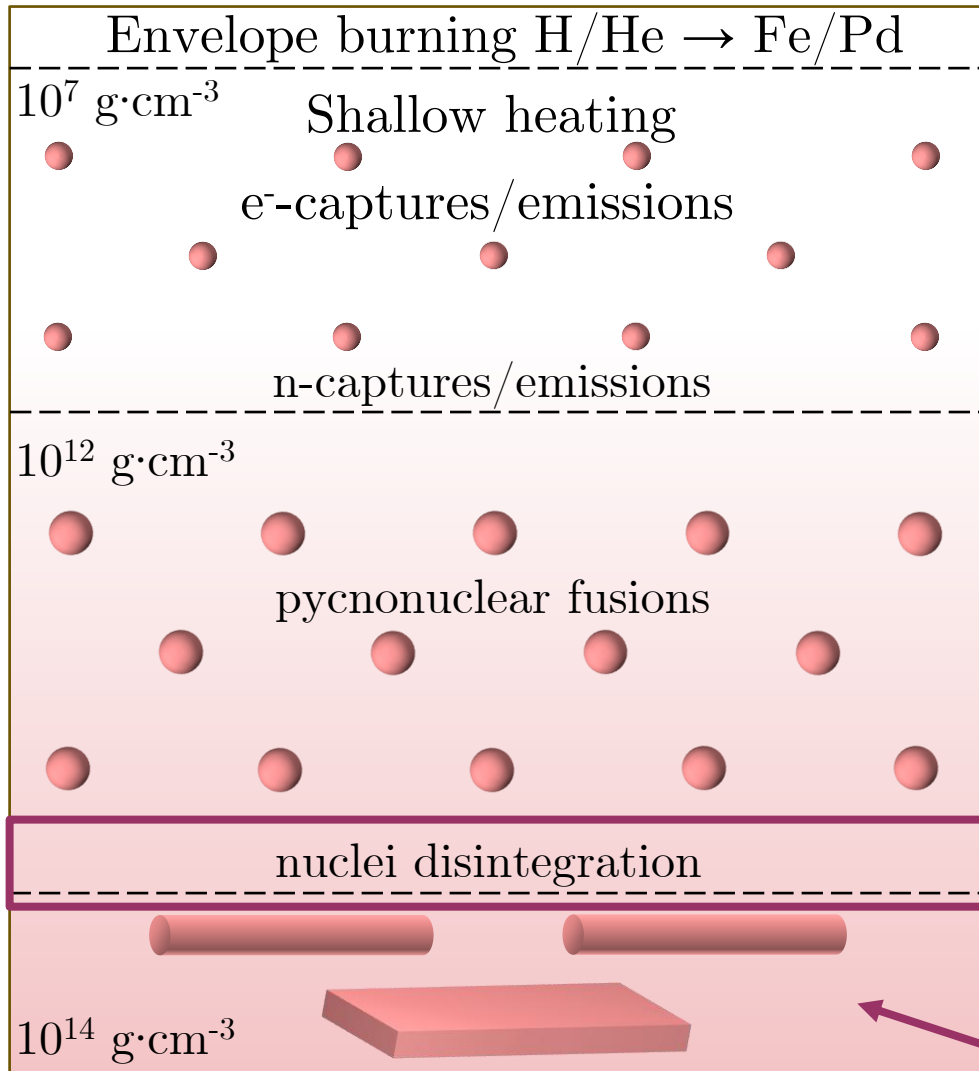
Evolution in the upper layers of the inner crust



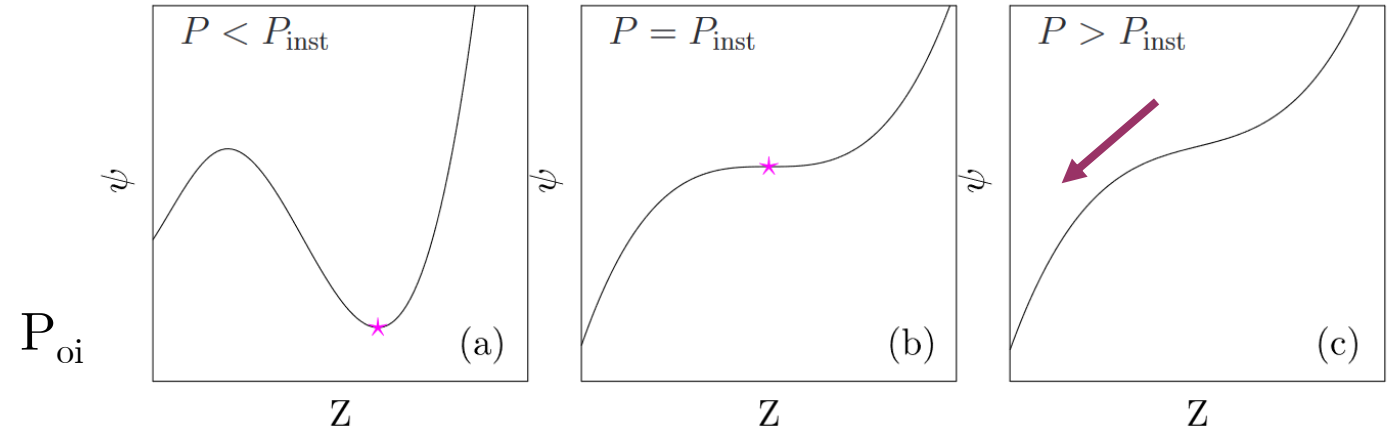
For Extreme rp ashes Q_{imp} stays high

Pycnonuclear fusions purify the crust only for Kepler ashes

Transition to the core



Absence of minimum in potential ψ signals instability



[Gusakov & Chugunov'24]

Nuclei are converted into the npe matter by unstoppable electron capture cascades (and possible pycnonuclear fusions)

[Gusakov & Chugunov'24]: most of the heat in the inner crust is released at the instability point

$$Q_{\text{inst}} \gtrsim 0.05 \text{ MeV per baryon}$$

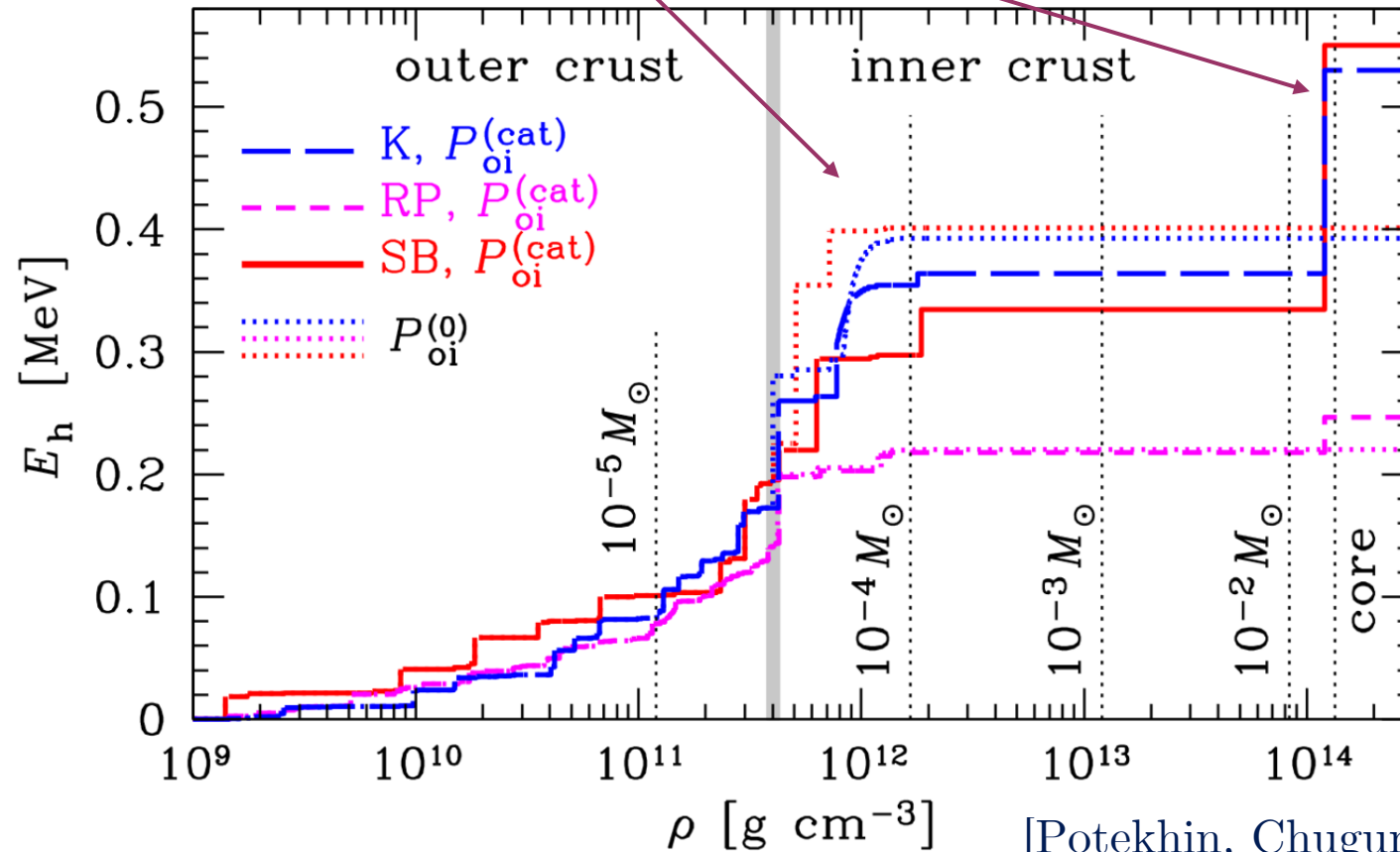
Some relic layer consisting of nuclear pasta can survive?

Heat profile in the crust

Main reactions are near the outer-inner crust transition and at the crust-core interface

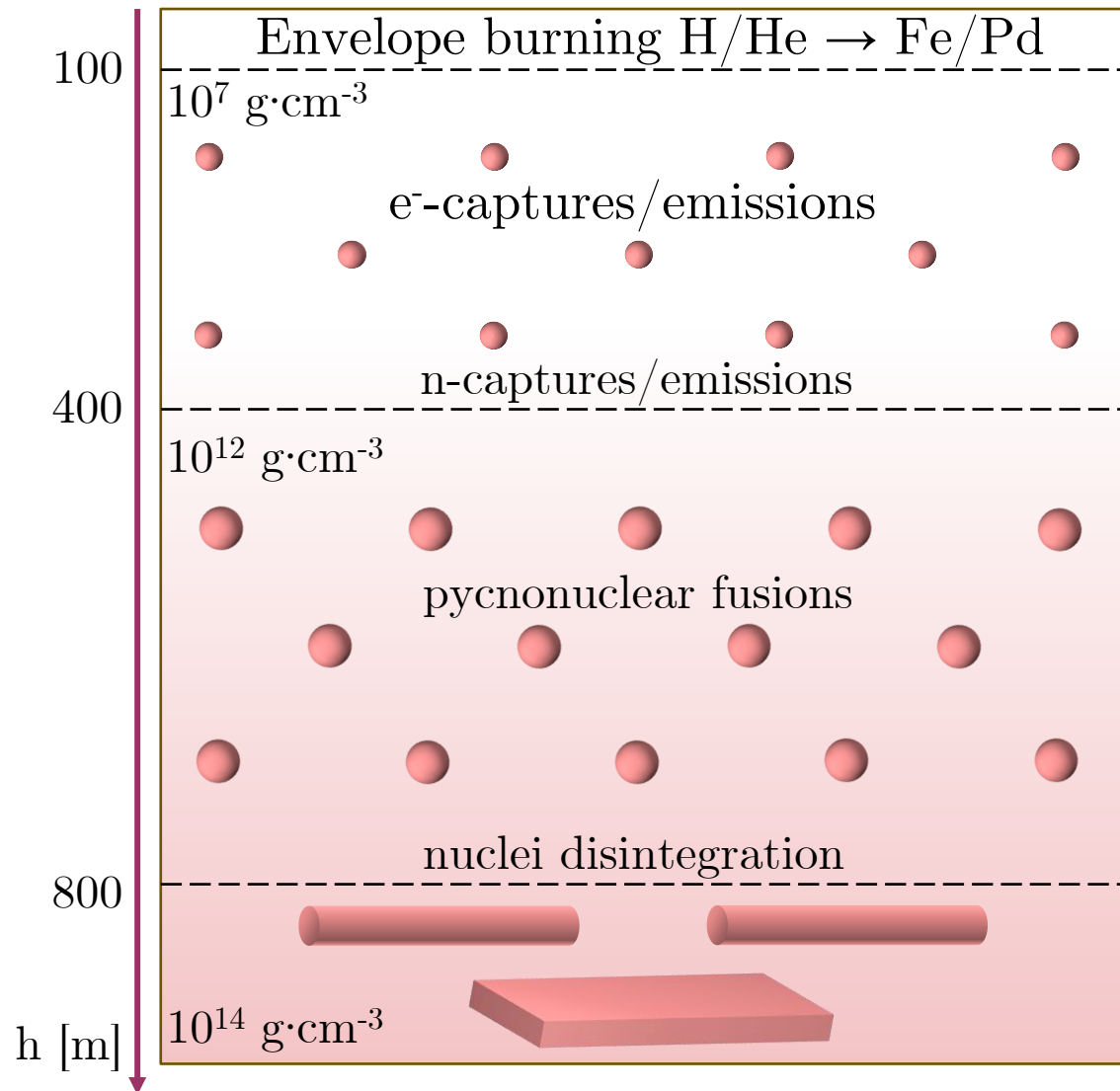
Total heat release

- Kepler and Superburst ashes $Q_{\text{tot}} \approx \mathbf{0.5}$ MeV/baryon
- Extreme rp $Q_{\text{tot}} \approx \mathbf{0.2}$ MeV/baryon



[Potekhin, Chugunov, NS, Gusakov'24]

Accreted crust models



Evolution of multicomponent composition with FRDM12 and heating sources are available at <https://www.ioffe.ru/astro/NSG/accrust/>

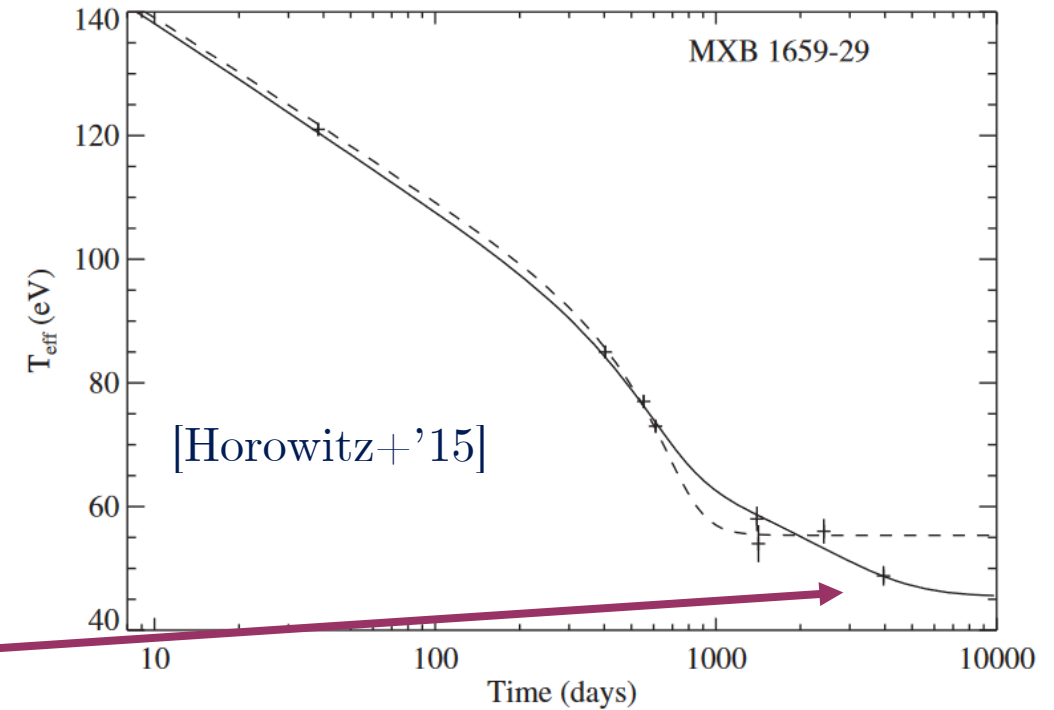
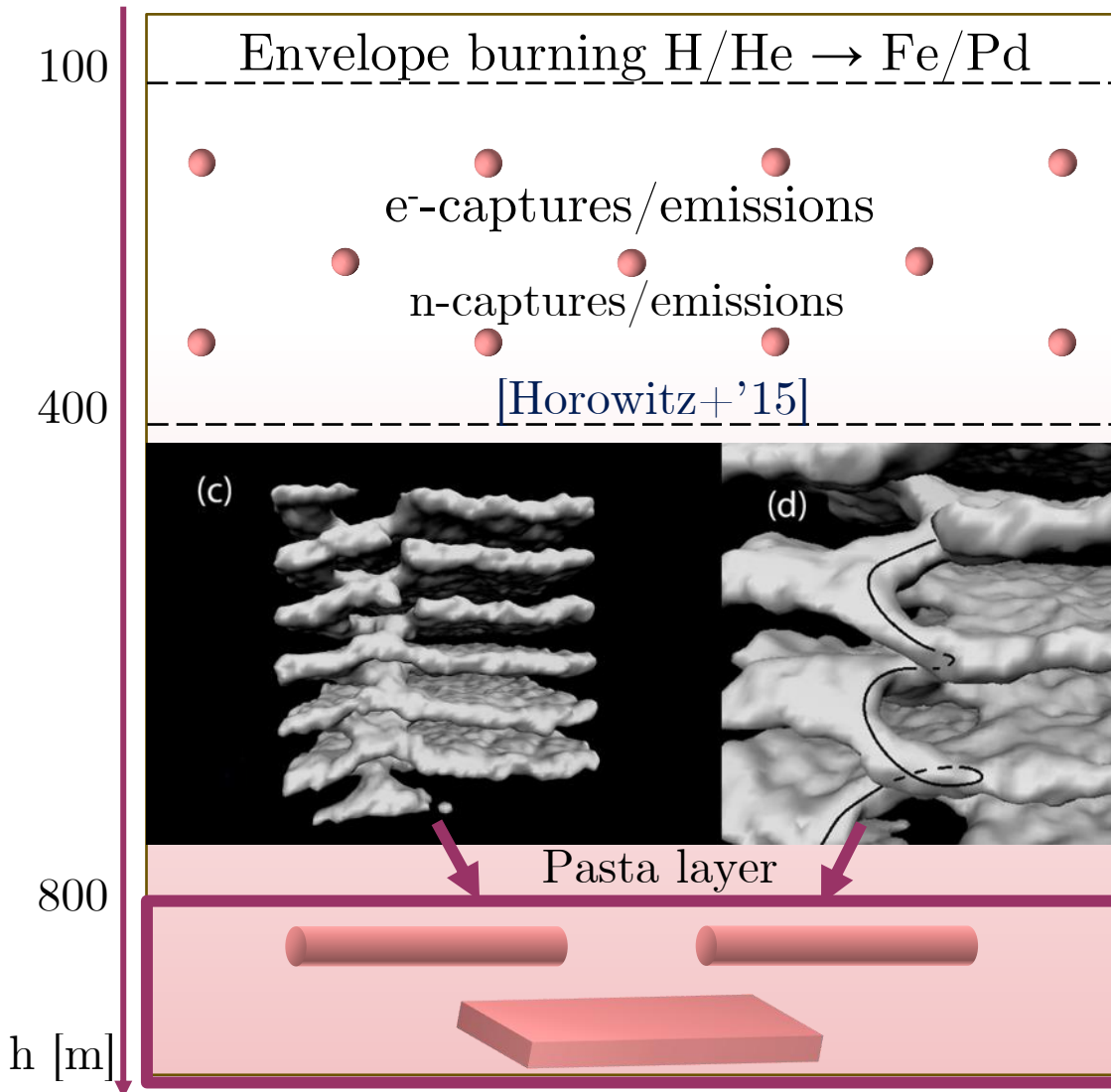
[NS, Gusakov, Chugunov '21,22,23]

One-component CLDM+Shell effects

[Gusakov & Chugunov '24]

Pasta imprints on the cooling curve

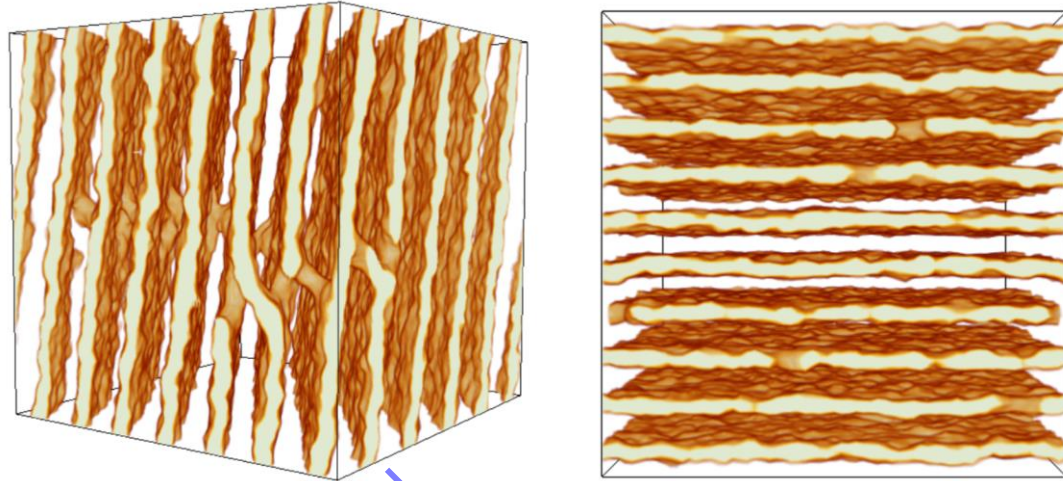
Topological defects in molecular dynamic simulations point towards low-conductivity layer in the crust bottom (typically modelled with $Q_{\text{imp}} \approx 40$)



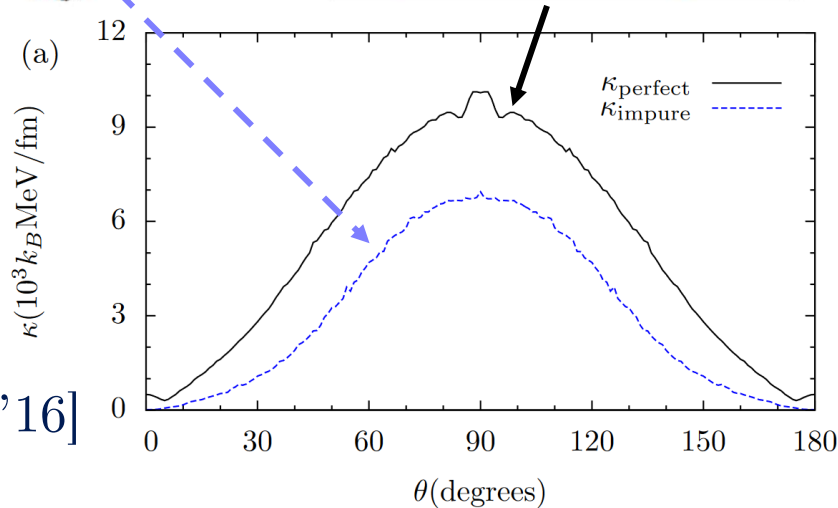
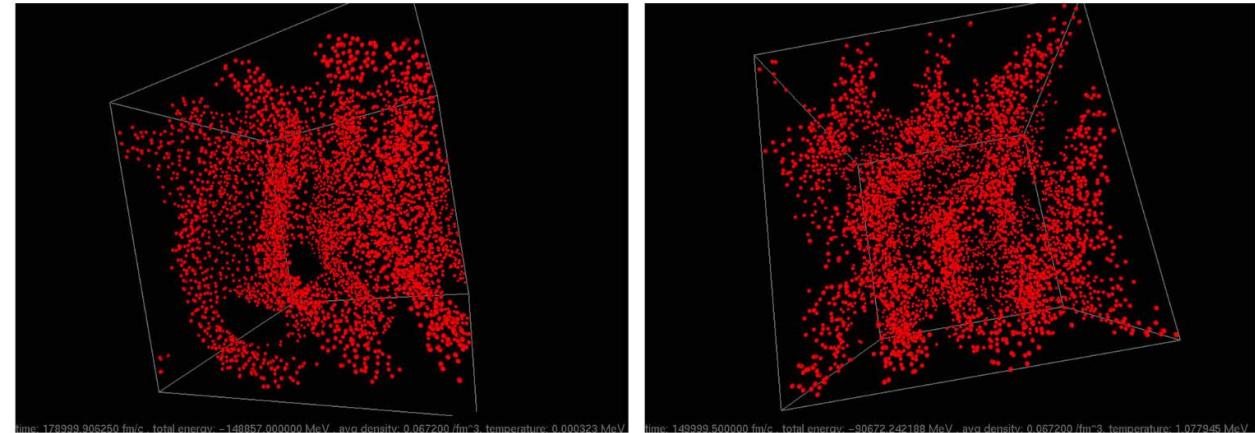
Talk of N. Chamel for explanation with gapless superfluidity

Do pasta have lower conductivity?

Classical molecular dynamics
(point-like particles)



VS **Quantum molecular dynamics**
(Gaussian wave-packets)



[Schneider+'16]

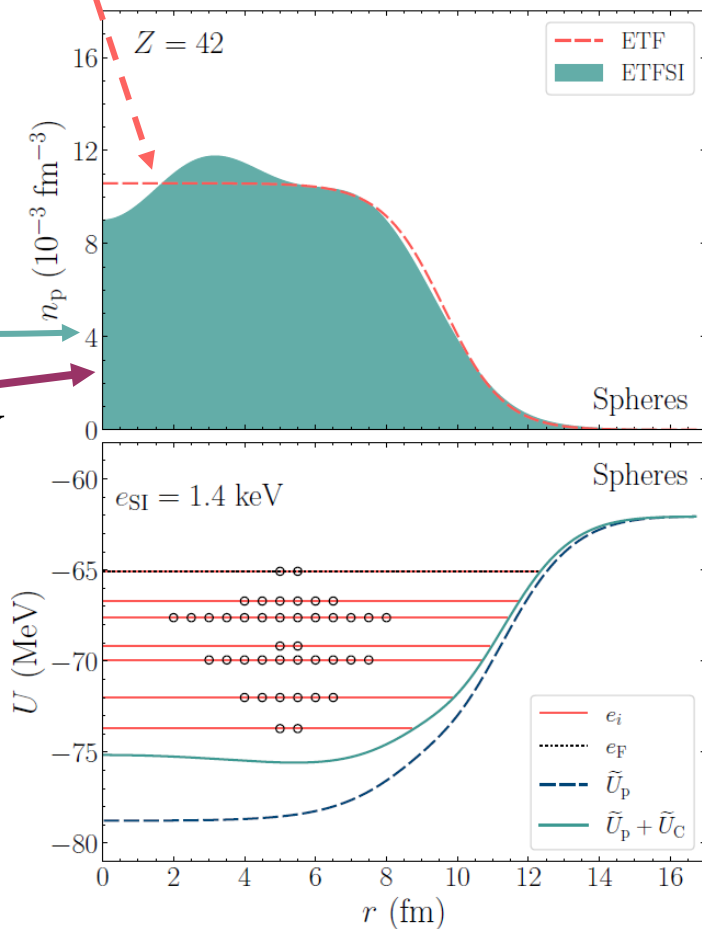
Suppressed conductivity for impure lasagna,
 $n_b = 0.05 \text{ fm}^{-3}$, $T = 1 \text{ MeV}$, $Y_p = 0.4$, $N \approx 4 \cdot 10^5$

$n_b = 0.064 \text{ fm}^{-3}$, $T = 0,1 \text{ MeV}$, $Y_p = 0.3$, $N \approx 1.6 \cdot 10^4$

[Nandi & Schramm '18] the effect of pasta is not dramatic for the transport properties

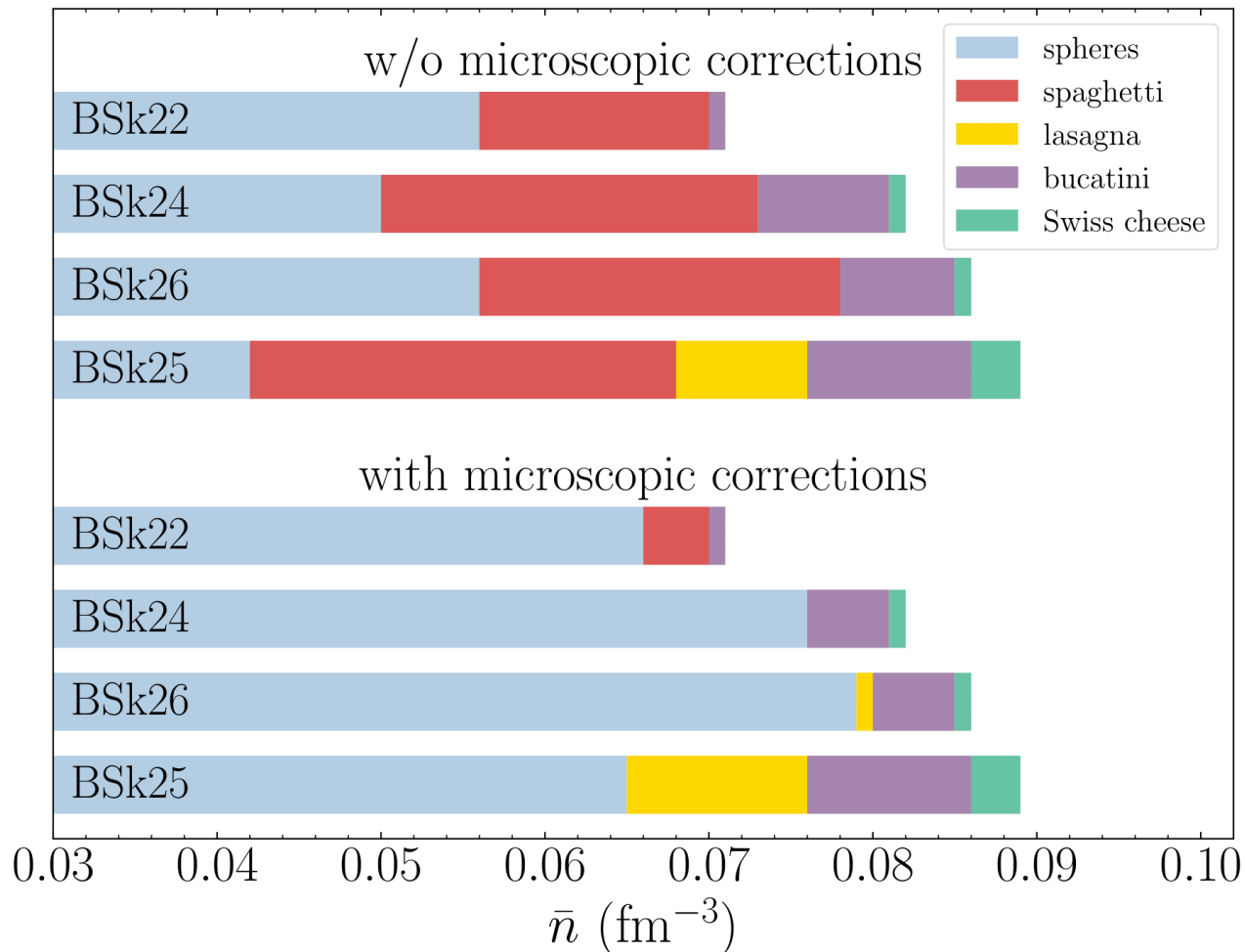
Do we have sizable amount of pasta at all?

With perturbative method for adding microscopic (shell + pairing) corrections on top of semi-classical calculations



Proton density distribution

Shrinking of pasta layer with microscopic corrections



Pasta sequence in the mantle

3D HF+BCS in the mantle

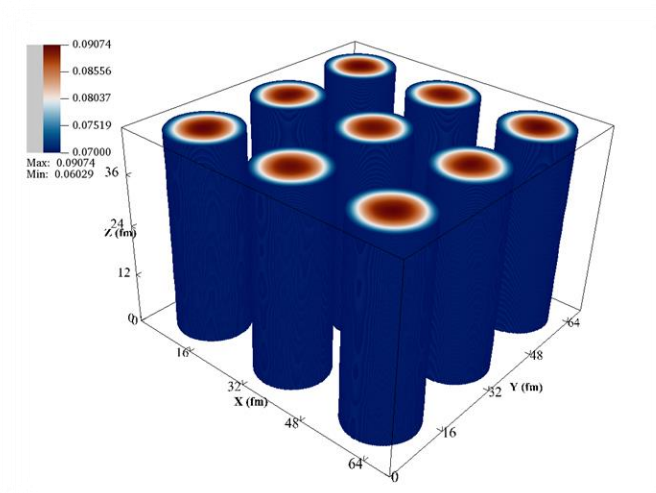
In attempt to **systematically study pasta phases in large domains within fully microscopic framework** we generalize 3D HFB code MOCCa [Ryssens, Heenen, Bender '15]

Proof of principle example:

Iterative procedure for calculating spaghetti shape using ETFSI solution as a guide

$$n_b = 0.07 \text{ fm}^{-3}, N = 16164, Y_p = 0.033$$

Preliminary



[NS, Ryssens, Chamel]

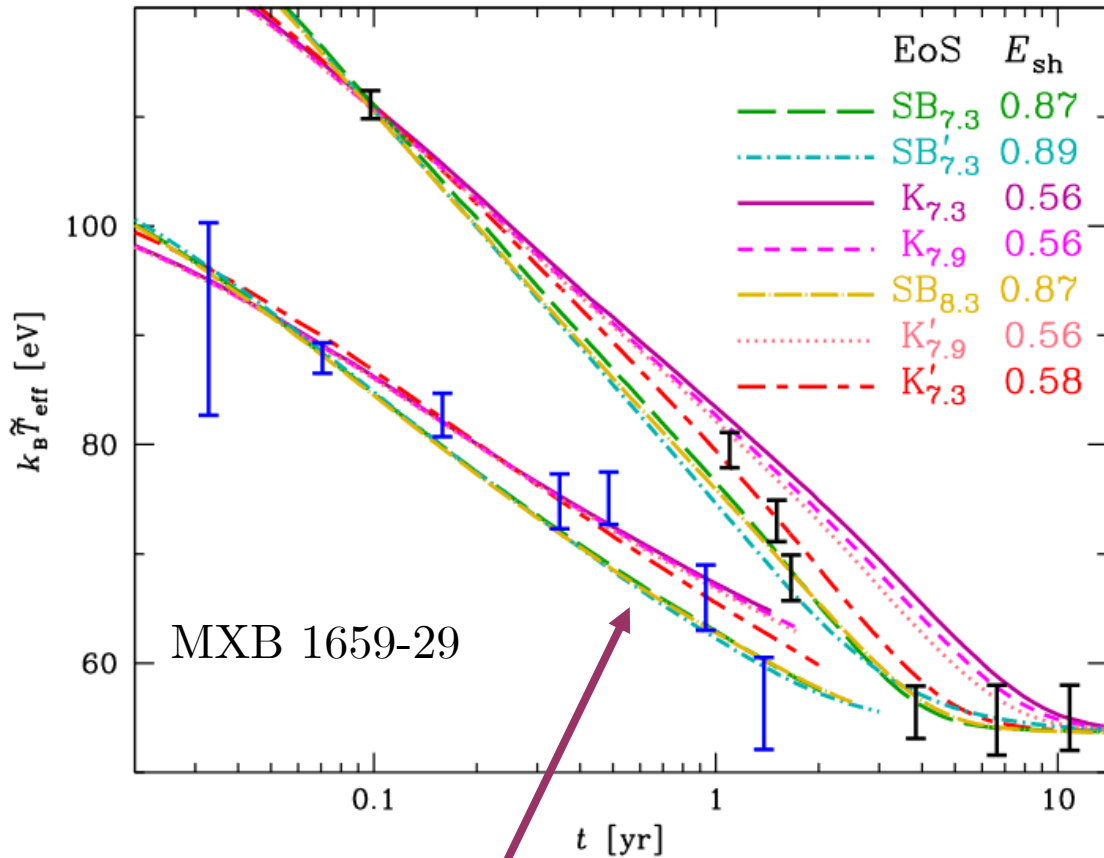
Using spatial symmetries

Conclusions

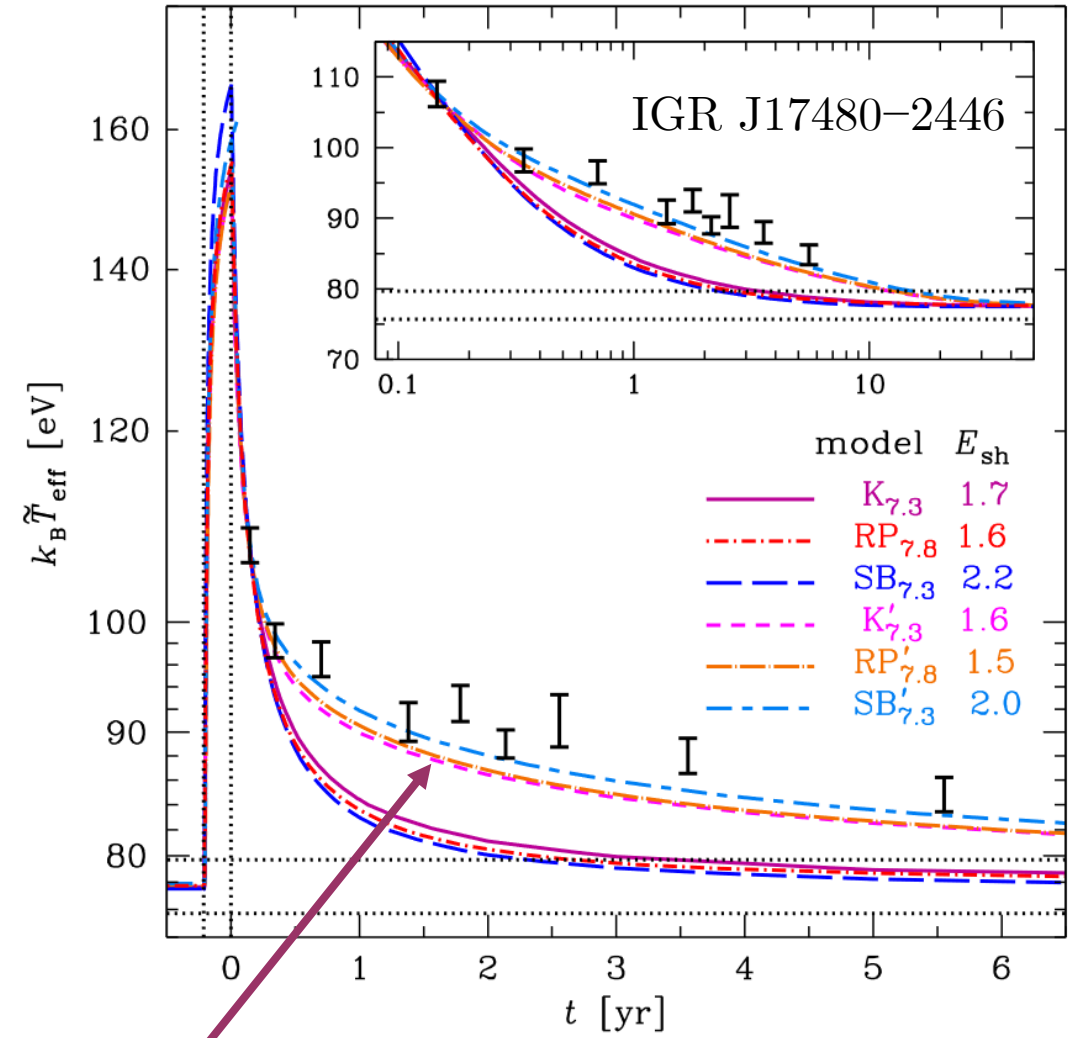
- We developed multicomponent simplified reaction network for the accreted crust considering **neutron hydrostatic and diffusion equilibrium** condition.
In contrast to the traditional approach:
 - Total deep crustal heating amounts only to **0.2-0.5 MeV** per baryon
 - Outer-inner crust transition and equation of state is close to the catalyzed crust
 - Pycnonuclear reactions purify the composition only if initial ashes were abundant in light elements ($A \approx 30$)
- It is **uncertain whether the pasta** layer possesses **reduced thermal conductivity** and whether it **exists at all**. To provide more definitive answer within **fully microscopic framework** at large scales we generalized 3D HFB code MOCCa and made **proof-of-principle calculations**. More results are coming...

Cooling with realistic crust composition and heating

[Potekhin, Chugunov, NS, Gusakov'24]



- Explained with pure ashes
- superburst?
 - partially accreted crust?



Disordered layer is required with $Q_{imp} = \langle Z \rangle^2$