



r-process nucleosynthesis



neutrinos and *r*-process nucleosynthesis

key quantity:

electron fraction

 $Y_e = \frac{1}{1 + (n/p)}$

Neutrinos can influence the:

- Initial electron fraction
- Entropy per baryon
- Free nucleons available for capture following seed formation

 $v_e + n \rightleftharpoons p + e^ \bar{v}_e + p \rightleftharpoons n + e^+$

$$Y_e = 0.02 Y_e = 0.14 Y_e = 0.18 Y_e = 0.24 Y_e = 0.12 Y_e = 0.16 Y_e = 0.21 Y_e = 0.28$$





neutrinos and *r*-process nucleosynthesis

Neutrinos can influence the:

- Initial electron fraction
- Entropy per baryon
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neutrino oscillations and *r*-process nucleosynthesis



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Facility for Rare Isotope Beams





First experiment started 9 May 2022





Interpreting observables of *r*-process nucleosynthesis

- What observables are currently limited by nuclear uncertainties that could be addressed in the FRIB era?
- Are there distinguishing observables that rise above nuclear uncertainties?
- What can we learn about nuclear physics far from stability from *r*-process observables?



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nuclear masses and the rare earth peak



nuclear masses and the rare earth peak



deducing *r*-process conditions from abundance pattern details: the rare earth peak

mass modification parameterization:

 $M(Z,N) = M_{DZ}(Z,N) + a_N e^{-(Z-C)^2/2f}$



Mumpower, McLaughlin, Surman, Steiner, 2016



deducing *r*-process conditions from abundance pattern details: the rare earth peak





reverse-engineering results for a hot wind *r*-process



Nicole Vassh, FIRE/ND postdoc now TRIUMF scientist

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reverse-engineering results for a hot wind *r*-process + new experimental masses



Nicole Vassh, FIRE/ND postdoc now TRIUMF scientist



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rare earth peak formation and *r*-process dynamics

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reverse-engineered nuclear mass predictions for three types of astrophysical environments compared to experimental data

Vassh, McLaughlin, Mumpower, Surman 2021





rare earth peak formation and fission products

Neodymium (*Z*=60) isotopes

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rare earth peak experimental prospects





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Korobkin, Hungerford, Fryer, Mumpower, Misch, Sprouse, Lippuner, Surman, Couture, Bloser, Shirazi, Evan, Vestrand, Miller 2020

also Hotokezaka+2016; Li 2019; Wu+2019; Ruiz-Lapuente, Korobkin 2020





Korobkin, Hungerford, Fryer, Mumpower, Misch, Sprouse, Lippuner, Surman, Couture, Bloser, Shirazi, Evan, Vestrand, Miller 2020

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Wang, Vassh, Sprouse, Mumpower, Vogt, Randrup, Surman, ApJL 2020







Wang, Vassh, Sprouse, Mumpower, Vogt, Randrup, Surman, ApJL 2020





Wang+ in preparation 2023









Wang, Clark, Ellis, Ertel, Fields, Fry, Liu, Miller, Surman, ApJ 2021; Wang, Clark, Ellis, Ertel, Fields, Fry, Liu, Miller, Surman, ApJ 2023





Wang, Clark, Ellis, Ertel, Fields, Fry, Liu, Miller, Surman, ApJ 2021; Wang, Clark, Ellis, Ertel, Fields, Fry, Liu, Miller, Surman, ApJ 2023





Wang, Clark, Ellis, Ertel, Fields, Fry, Liu, Miller, Surman, ApJ 2021; Wang, Clark, Ellis, Ertel, Fields, Fry, Liu, Miller, Surman, ApJ 2023

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Actinide observables: lunar regolith



Wang, Clark, Ellis, Ertel, Fields, Fry, Liu, Miller, Surman, ApJ 2023



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UNEDF1 masses

Sprouse, Navarro Perez, Surman, Mumpower, McLaughlin, Schunck 2020

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TABLE II: Optimized parameter set UNEDF1. Listed are bounds used in the optimization, final optimized parameter values, standard deviations, and 95% confidence intervals.

x	Bounds	$\hat{\mathbf{x}}^{(ext{fin.})}$	σ	95% CI
$ ho_{ m c}$	[0.15, 0.17]	0.15871	0.00042	[0.158, 0.159]
$E^{\rm NM}/A$	[-16.2, -15.8]	-15.800	_	—
$K^{\rm NM}$	[220, 260]	220.000	—	—
$a_{\mathrm{sym}}^{\mathrm{NM}}$	[28, 36]	28.987	0.604	[28.152, 29.822]
$L_{\rm sym}^{ m NM}$	[40, 100]	40.005	13.136	[21.841, 58.168]
$1/M_s^*$	[0.9, 1.5]	0.992	0.123	[0.823, 1.162]
$C_{0}^{\rho\Delta\rho}$	$[-\infty, +\infty]$	-45.135	5.361	[-52.548, -37.722]
$C_1^{\rho\Delta\rho}$	$[-\infty,+\infty]$	-145.382	52.169	[-217.515, -73.250]
V_0^n	$[-\infty, +\infty]$	-180.005	18.510	[-211.000, -100.404]
V_0^p	$[-\infty,+\infty]$	-206.580	13.049	[-224.622, -188.538]
$C_0^{\rho \nabla J}$	$[-\infty,+\infty]$	-74.026	5.048	[-81.006, -67.046] \approx
$C_1^{\rho \nabla J}$	$[-\infty,+\infty]$	-35.658	23.147	[-67.663, -3.654]

Sprouse, Navarro Perez, Surman, Mumpower, McLaughlin, Schunck 2020

UNEDF1 masses



weighted average A of the rare earth peak



Fission yield signatures





Fission yield signatures



Fission yield signatures



Roederer+ submitted 2023



summary

The origin of the heaviest elements in the *r*-process of nucleosynthesis has been one of the greatest mysteries in nuclear astrophysics for decades.

Despite considerable progress in the past several years, including the first direct detection of an *r*-process event, the *r*-process site(s) has not been definitively determined.

The neutrino and nuclear physics of candidate events remains poorly understood. FRIB has the potential to reduce key nuclear uncertainties, facilitating accurate interpretations of *r*-process observables such as abundance patterns and light curves.



Mumpower, Surman, McLaughlin, Aprahamian, JPPNP 2016

