

A new class of three-nucleon forces In Chiral EFT

Based on: [arXiv:2411.00097](https://arxiv.org/abs/2411.00097)

with

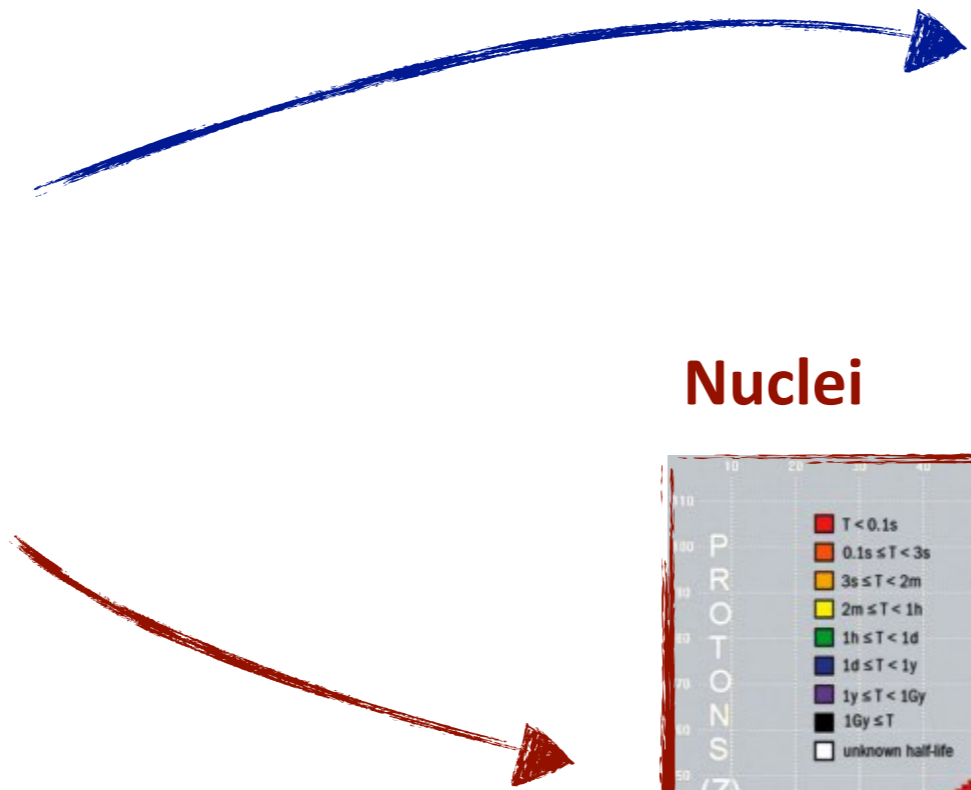
Maria Dawid, V. Cirigliano, S. Reddy



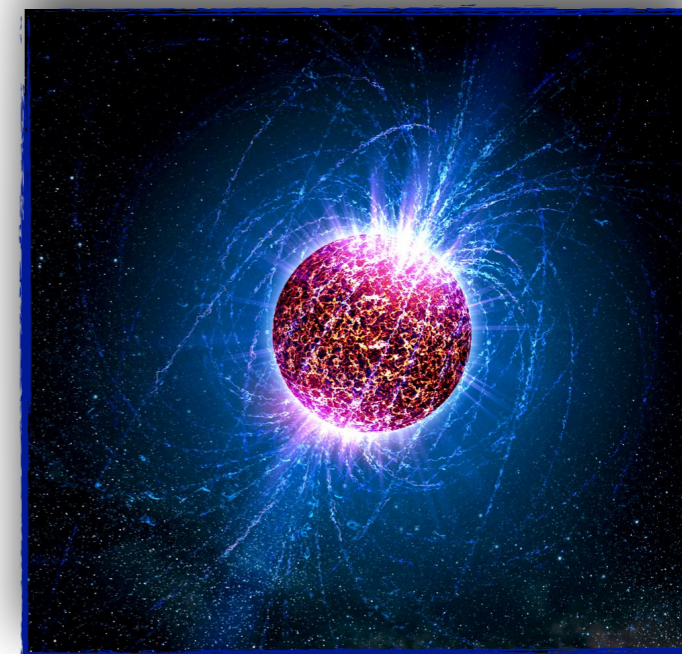
UNIVERSITY *of* WASHINGTON



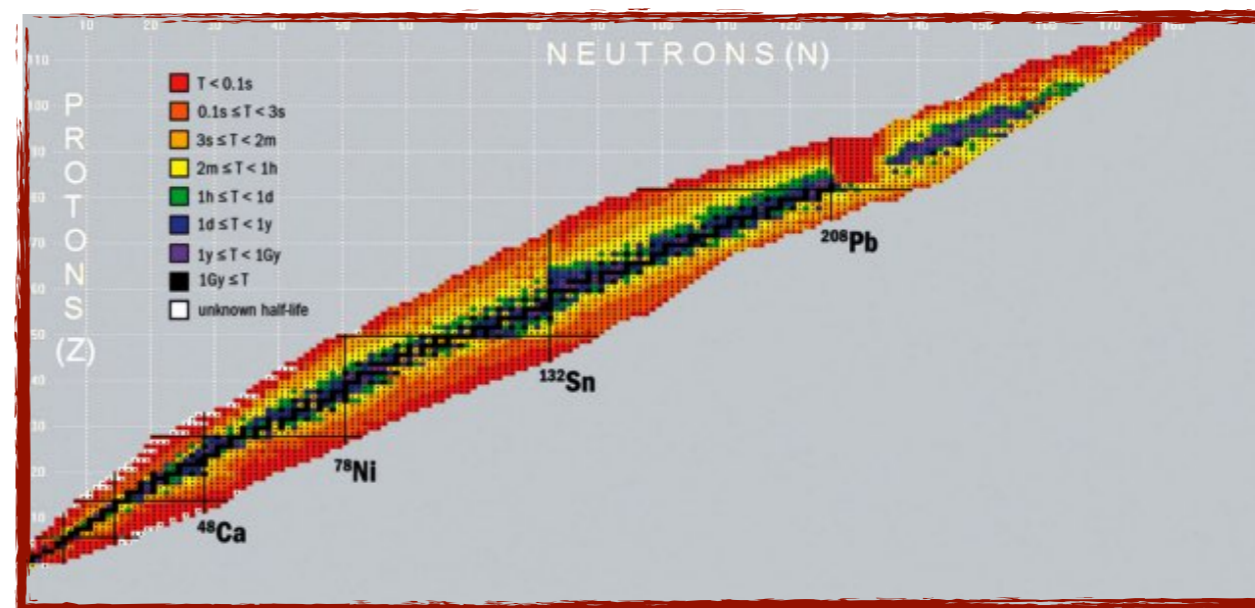
INSTITUTE *for*
NUCLEAR THEORY



Neutron Stars



Nuclei

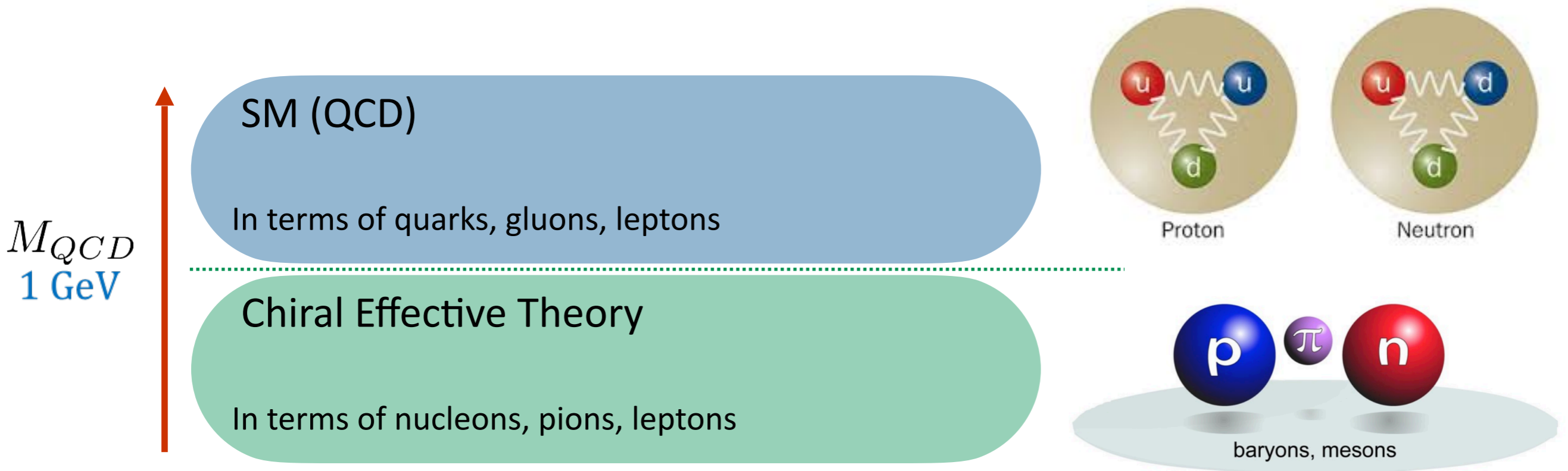


Chiral 3 nucleon forces

- Play an important role in nuclei & dense matter
 - Contribute to binding energies
 - Needed to describe saturation

From quarks to nucleons

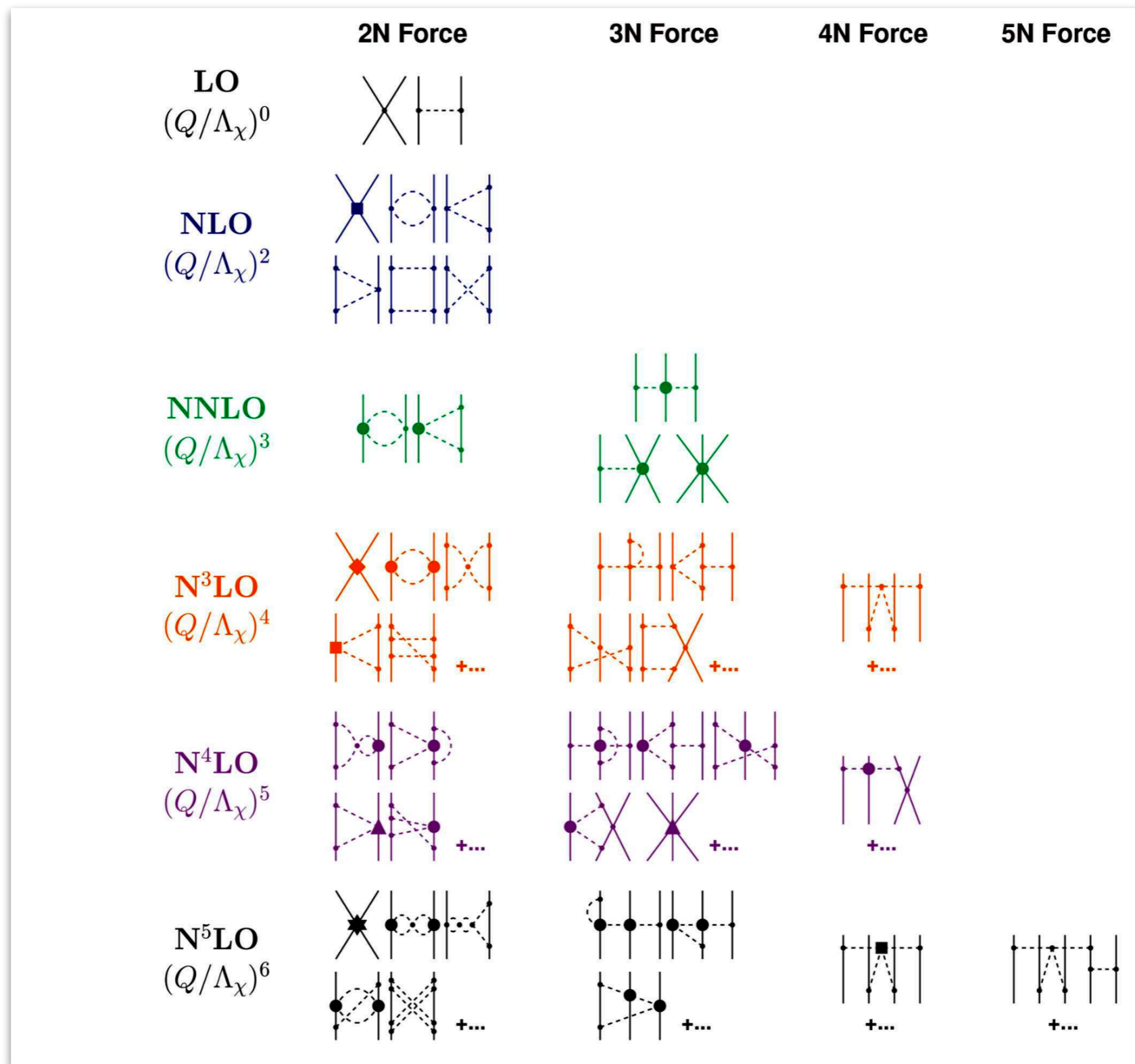
Modern view of nucleon forces



- Possible hadronic interactions determined by (chiral) symmetries
- Expansion in $Q/\Lambda_\chi \sim m_\pi/\Lambda_\chi$, where $\Lambda_\chi \sim 1 \text{ GeV}$
- Interactions come with (unknown) coupling constants
 - *Power counting needed to determine their importance*

Chiral Effective Theory

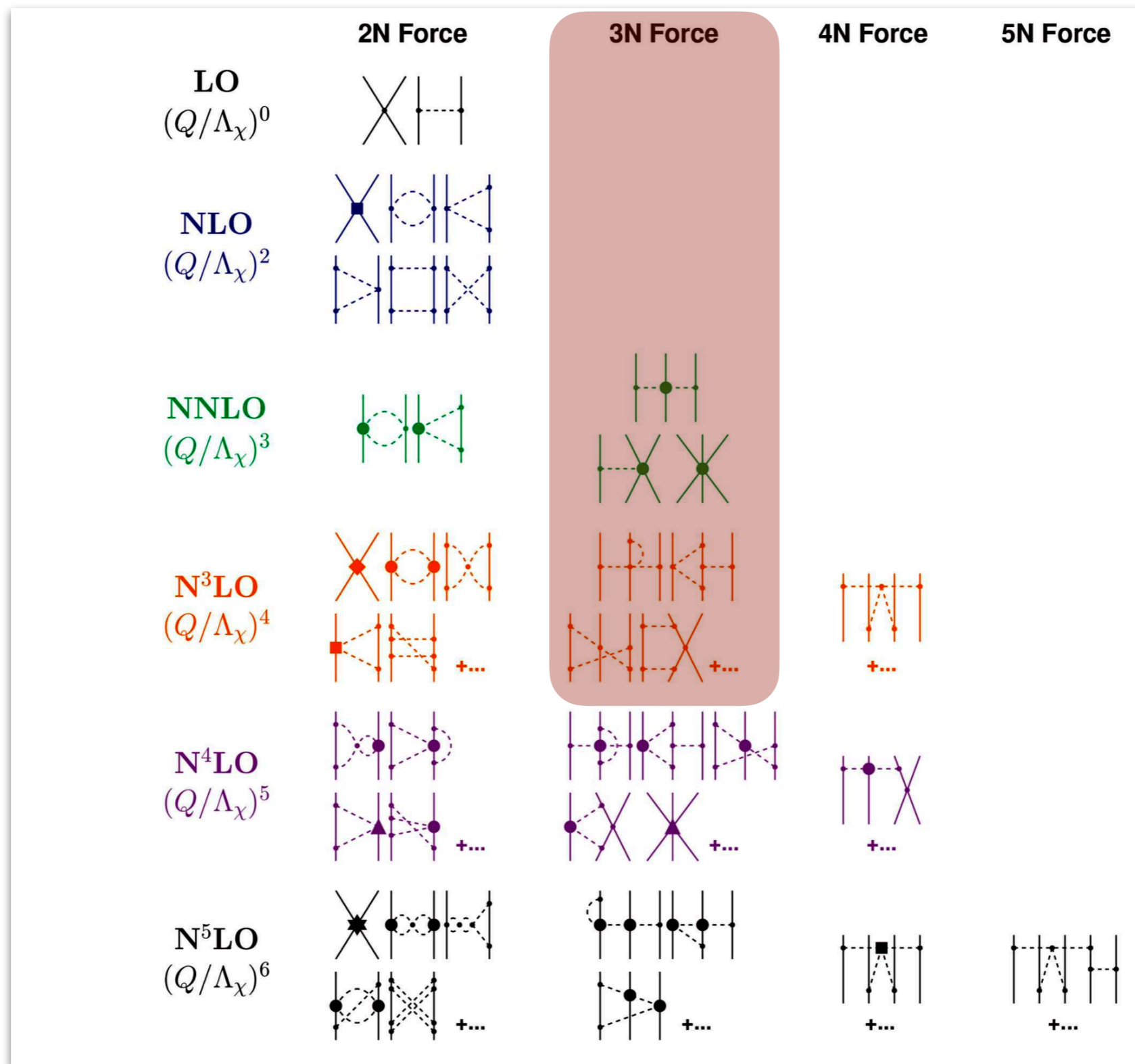
Nucleon forces in Weinberg counting



Chiral Effective Theory

Nucleon forces in Weinberg counting

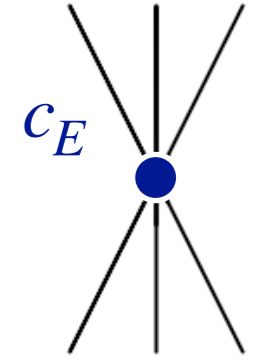
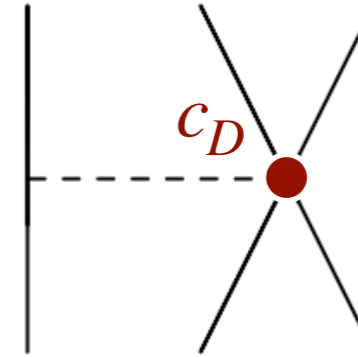
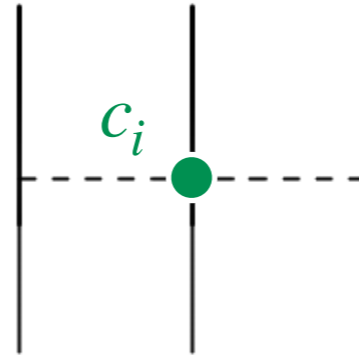
This talk



Conventional three-body force

N2LO contributions

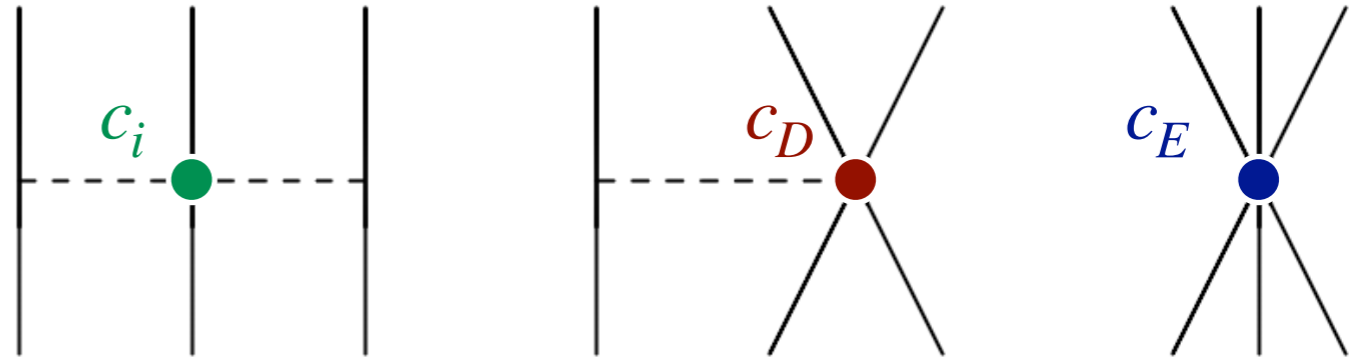
- Consists of
 - **One-** and **two-**pion exchange
 - **Short-distance** contributions



Conventional three-body force

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- LECs determined by

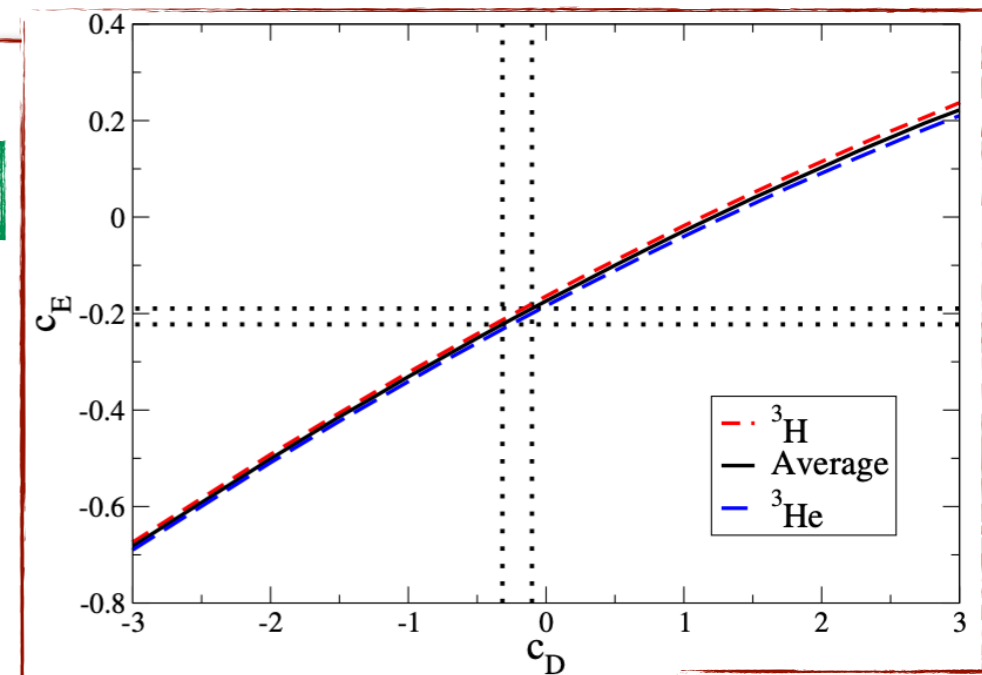
- c_i : πN scattering

Hoferichter et al Phys. Rept. '15; PRL '15; Phys. Lett. B '16

See Martin's talk

- $c_{D,E}$:

- Nd scattering
- light nuclei
- tritium β decay

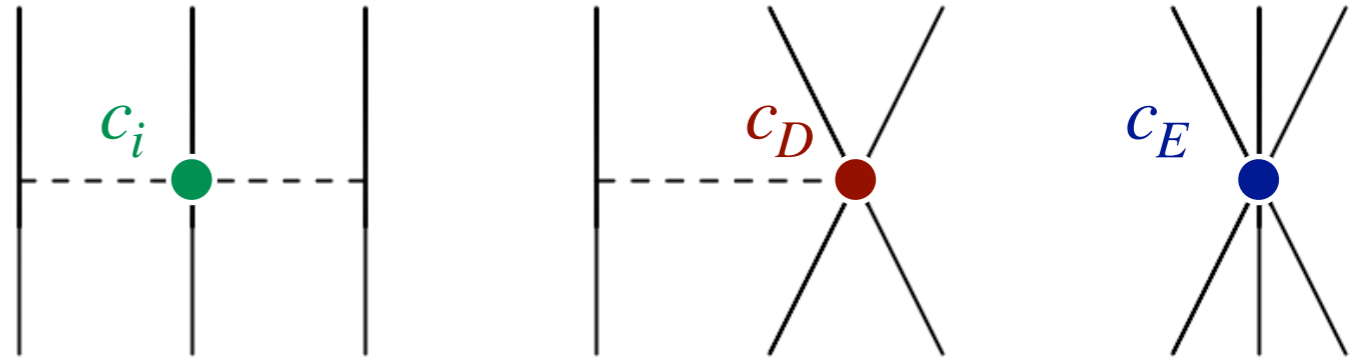


Gazit et al '09

Conventional three-body force

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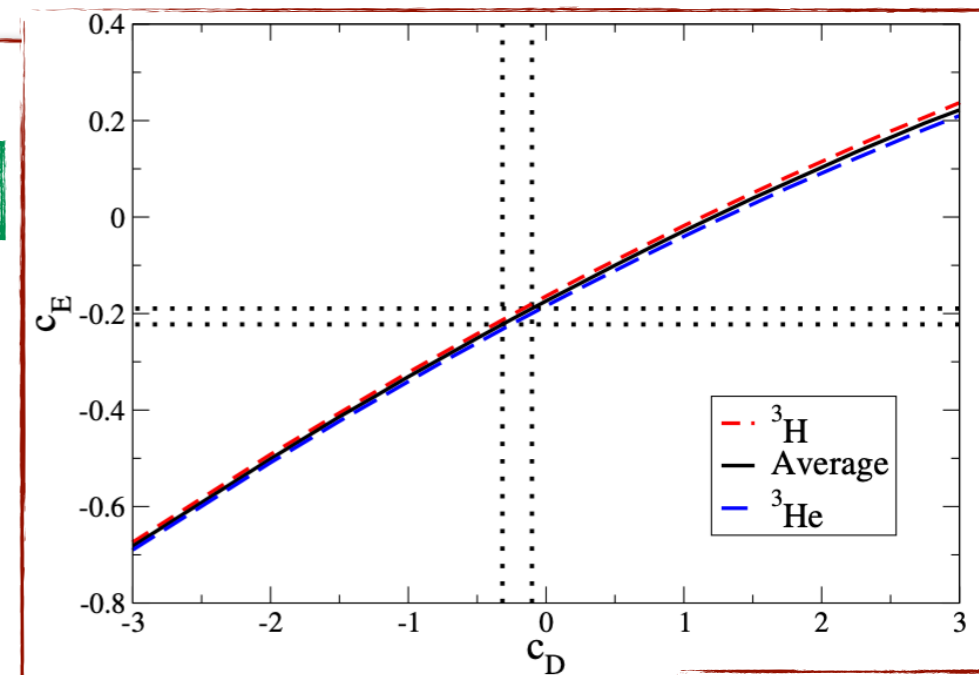
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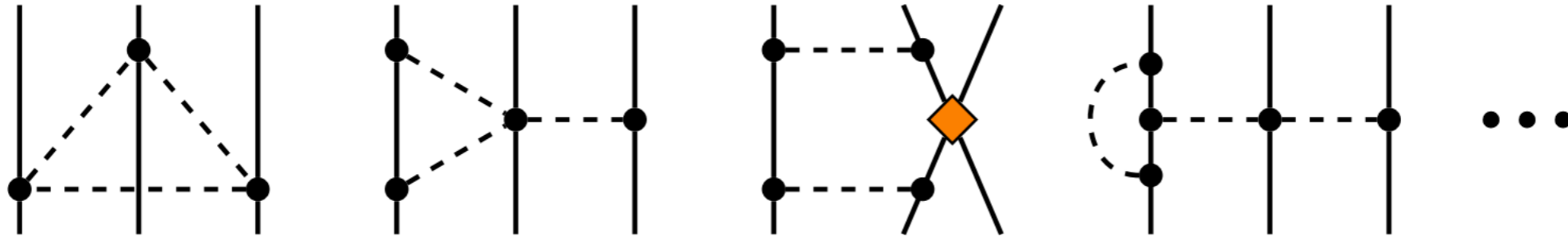


Gazit et al '09

- **Only c_1 and c_3 contribute in neutron matter**

Conventional three-body force

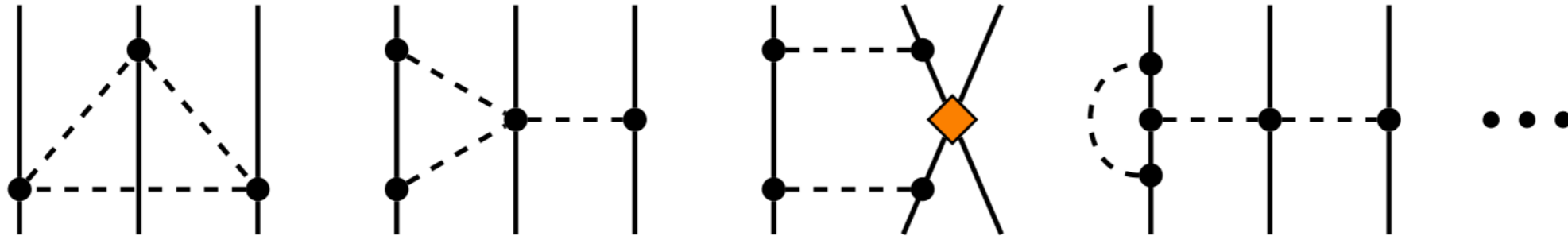
N³LO contributions



- Consists of
 - Loop diagrams with LO vertices
 - Tree graphs involving relativistic corrections

Conventional three-body force

N3LO contributions



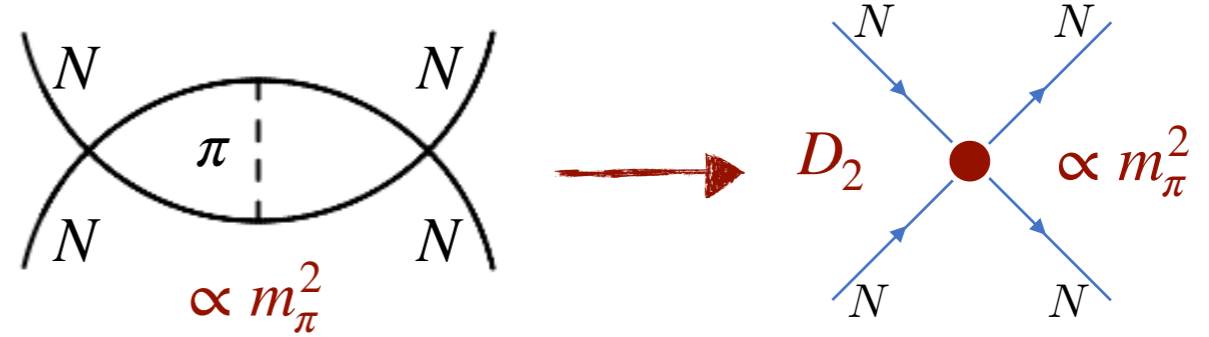
- Consists of
 - Loop diagrams with LO vertices
 - Tree graphs involving relativistic corrections

- **No new LECs**

Beyond Weinberg counting

Beyond Weinberg counting

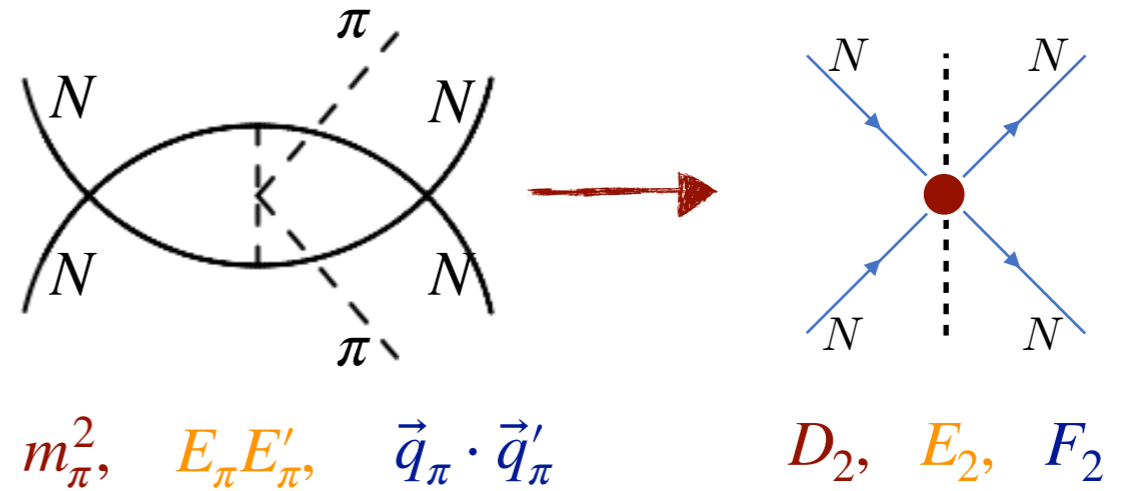
- Bubble diagrams in 1S_0 lead to divergences
- Requires short-distance interactions
- Can be absorbed in $C'_0 = C_0 + m_\pi^2 D_2$



Beyond Weinberg counting

Borasoy, Griesshammer, '01, '03

- Similar diagrams induce $\pi^2 NN$ couplings
- **Cannot** be absorbed in C_0
- Depend on m_π or momenta:

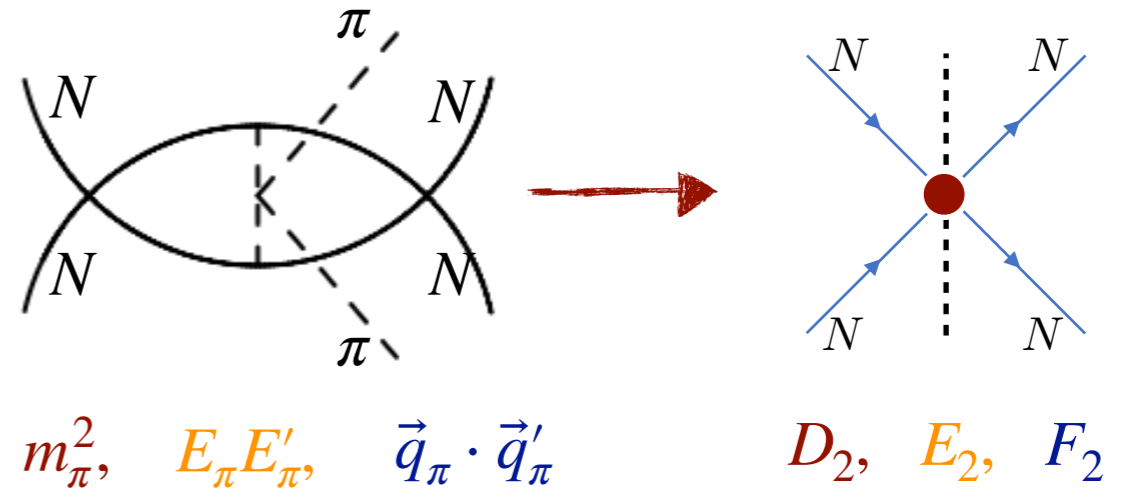


$$\mathcal{L} = -\frac{1}{4} D_2 (N^T P_i N)^\dagger (N^T P_i N) \langle \chi_+ \rangle + \frac{1}{4} [E_2 \langle (v \cdot u)^2 \rangle + F_2 \langle u \cdot u - (v \cdot u)^2 \rangle] (N^T P_i N)^\dagger (N^T P_i N)$$

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D_2, E_2, F_2 currently poorly known

- Estimate from renormalization group:

$$\frac{d}{d \ln \mu} \left[\frac{X}{\tilde{C}_0^2} \right] = \gamma_X \left(\frac{m_N}{4\pi f_\pi} \right)^2, \quad X \in \{D_2, E_2, F_2\}$$

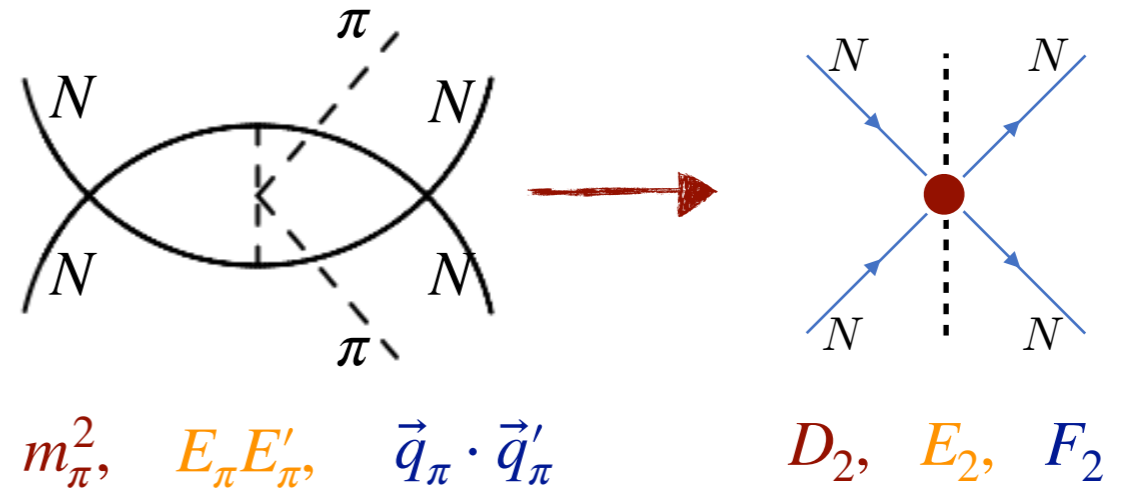
$$\begin{aligned} \gamma_{D_2} &= g_A^2/4 \\ \gamma_{E_2} &= -(1 + g_A^2)/3, \\ \gamma_{F_2} &= -g_A^2/3 \end{aligned}$$

- Suggests LO size $X \simeq C_0^2/5 \sim 1/(5 F_\pi^4)$
 - Compared to N2LO Weinberg $X \sim 1/(F_\pi^2 \Lambda_\chi^2)$

Beyond Weinberg counting

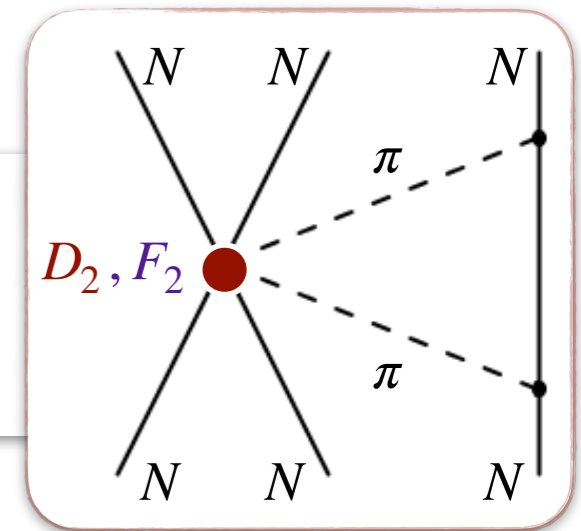
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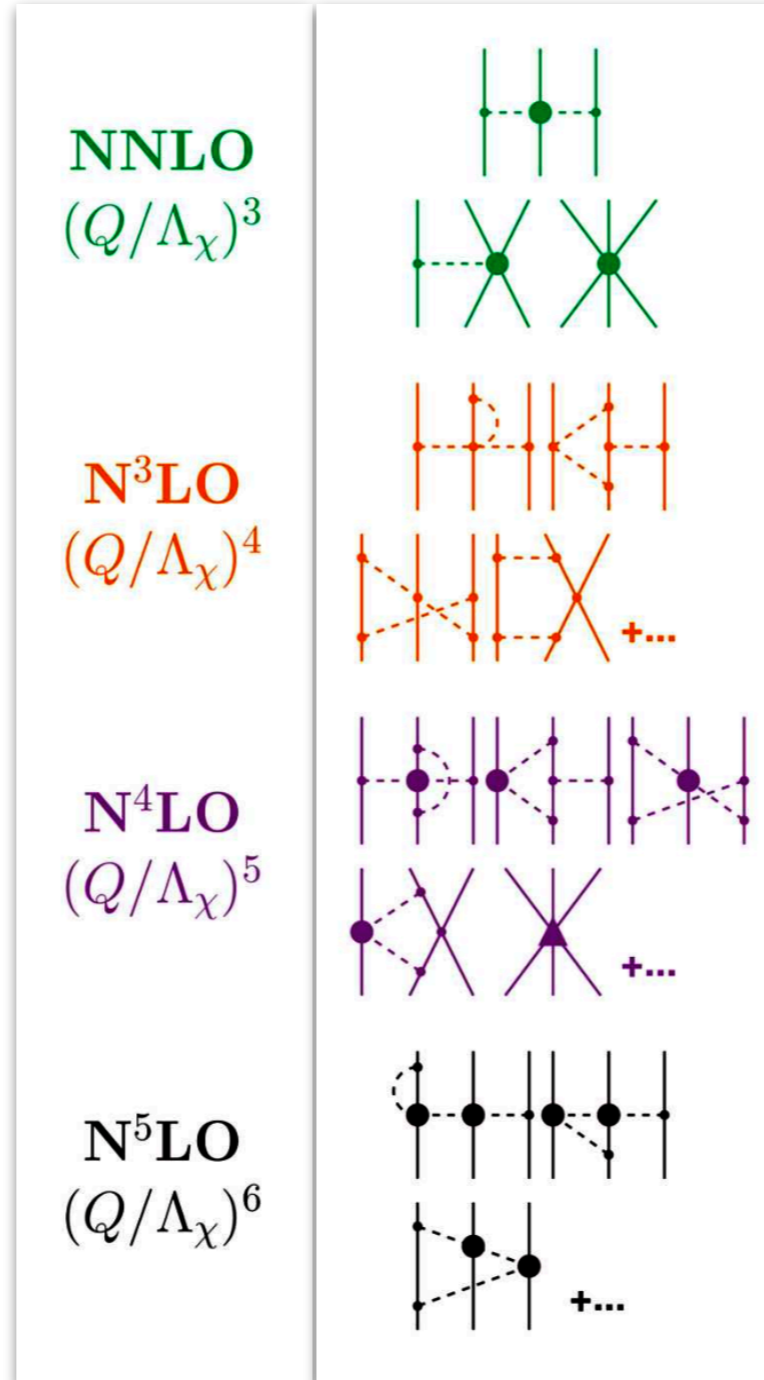
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- $\pi^2 NN$ interactions do not affect NN
- Induce 3-nucleon forces through loops



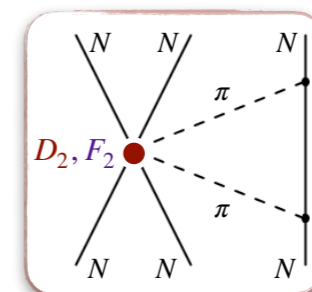
Beyond Weinberg counting

Three- nucleon forces



Hierarchy of nuclear forces up to N5 LO in ChiPT. Solid lines represent nucleons and dashed lines pions. Entem, Machleidt, Y. Nosyk, (arXiv:1703.05454)

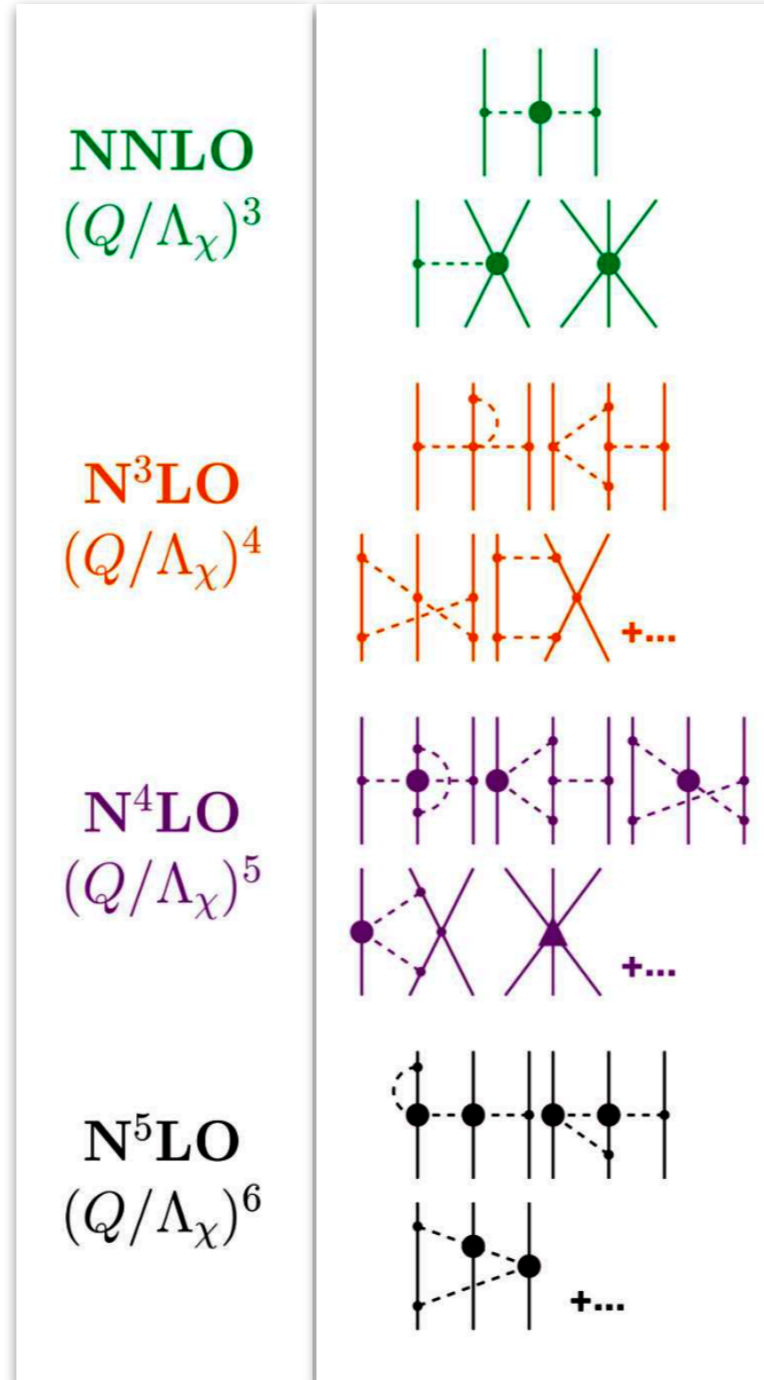
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Weinberg's counting

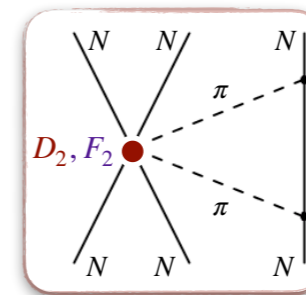
Beyond Weinberg counting

Three- nucleon forces

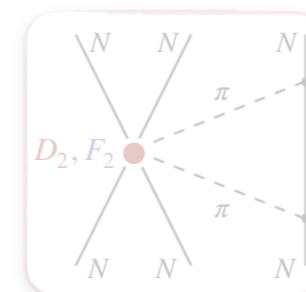


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Consistency requirement



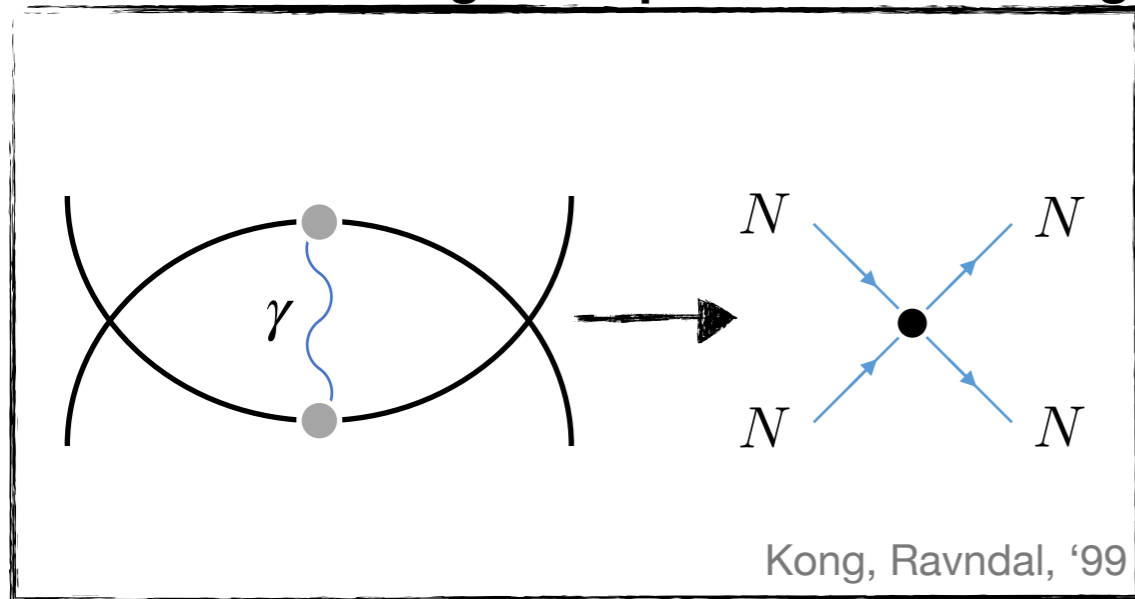
Weinberg's counting

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Electroweak/BSM interactions

Electroweak and BSM interactions affected by the same RG arguments:

Charge-independence Breaking



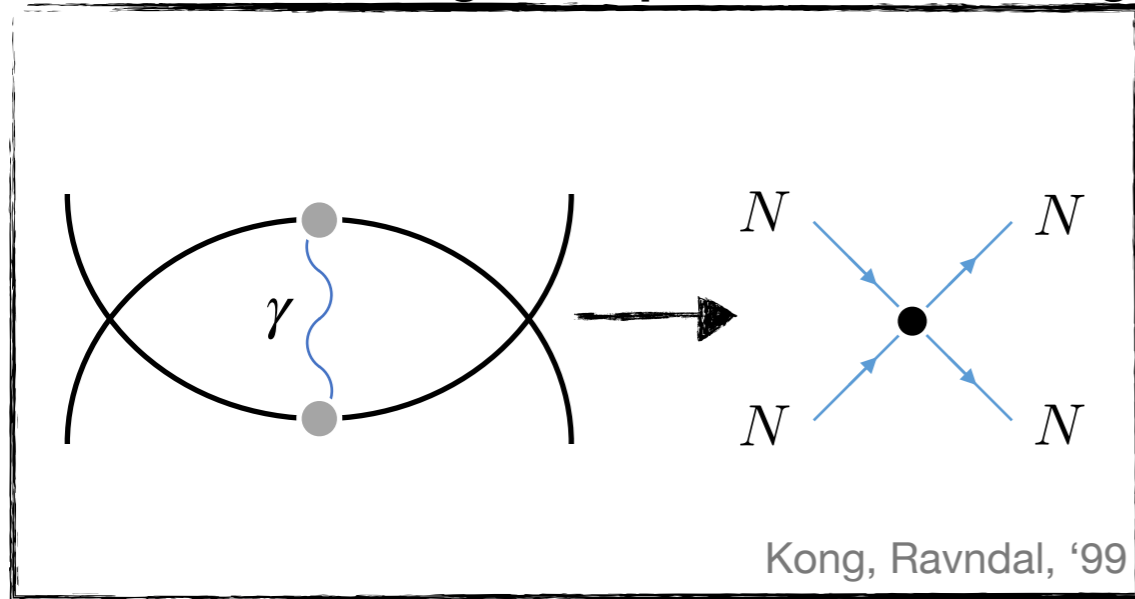
Beyond Weinberg counting

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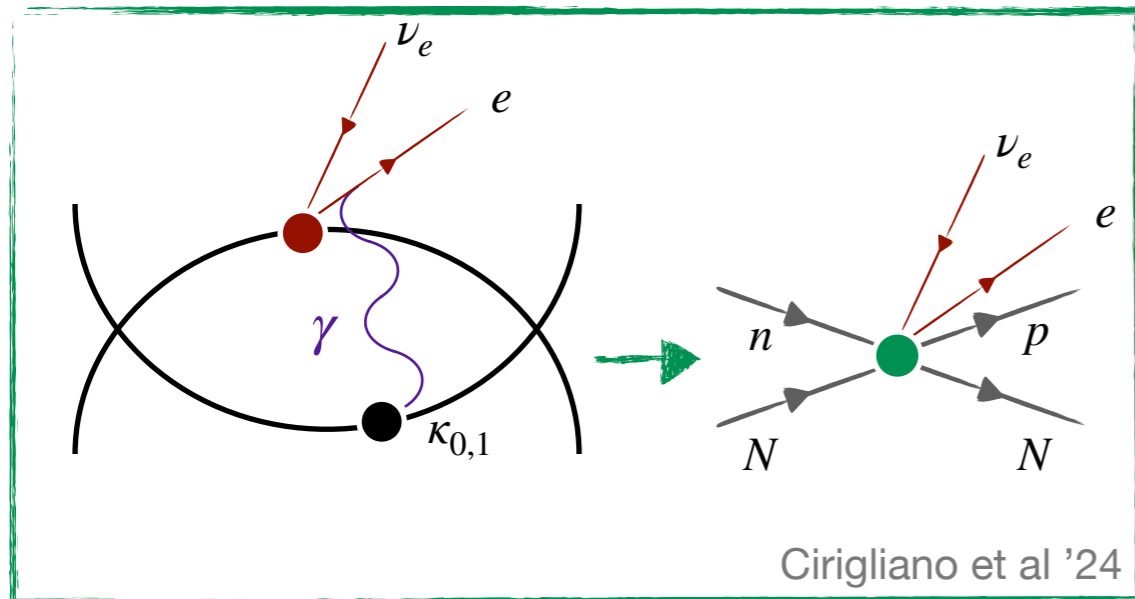
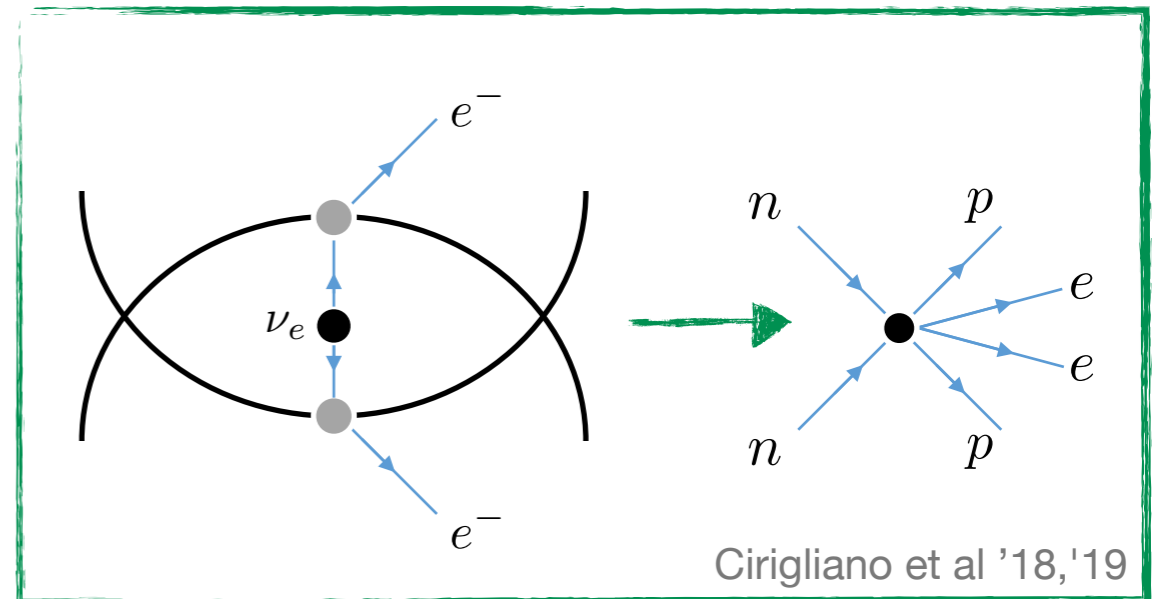
Electroweak and BSM interactions affected by the same RG arguments:

See Emanuele/Jordy's Talks on Thursday

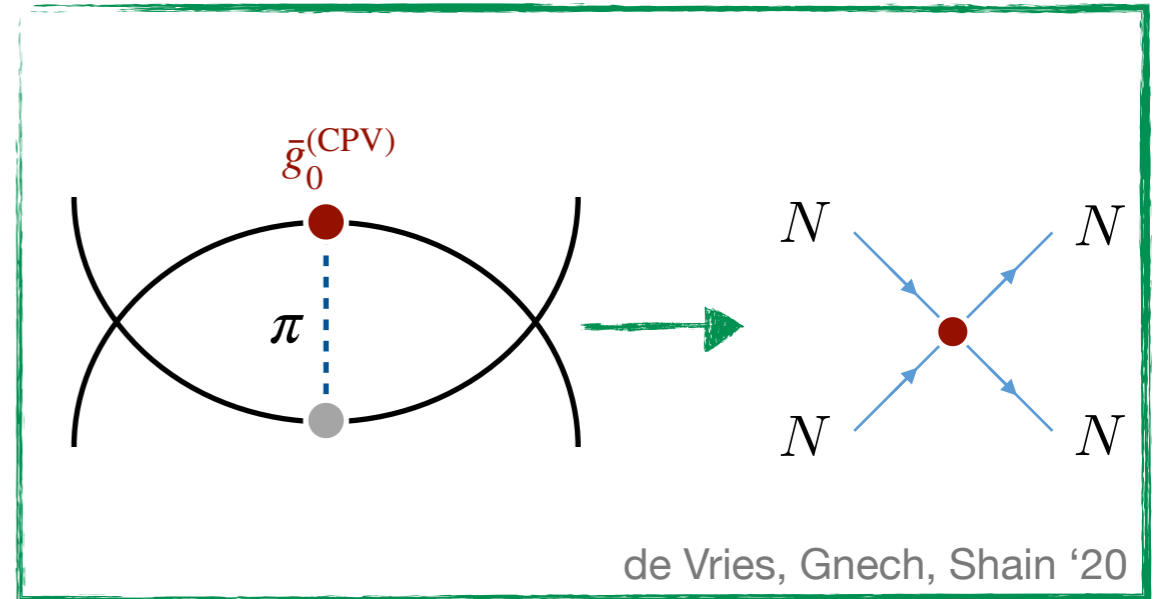
Charge-independence Breaking



$0\nu\beta\beta$



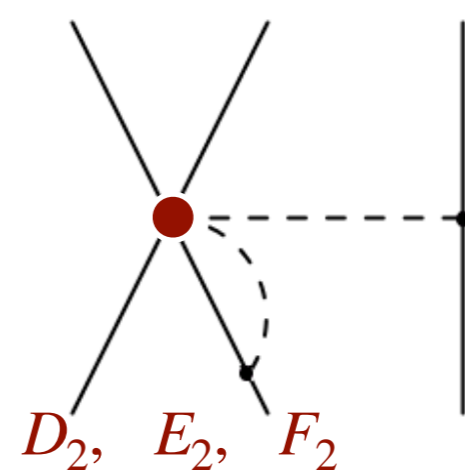
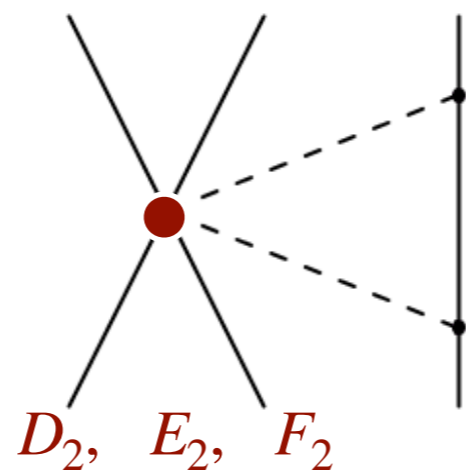
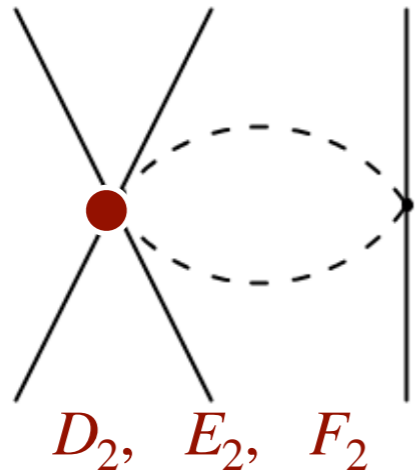
Radiative corrections in β decay



Electric Dipole Moments

Beyond Weinberg counting

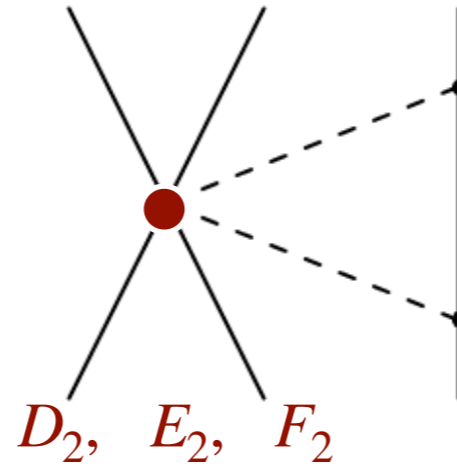
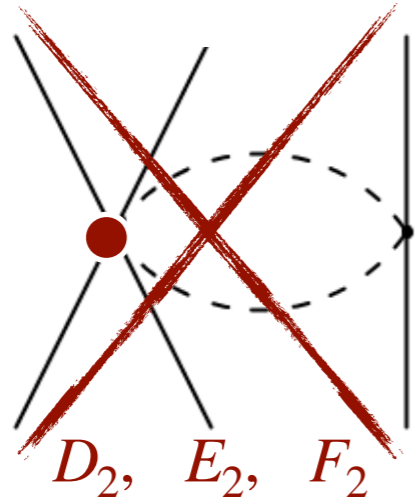
Induced potential



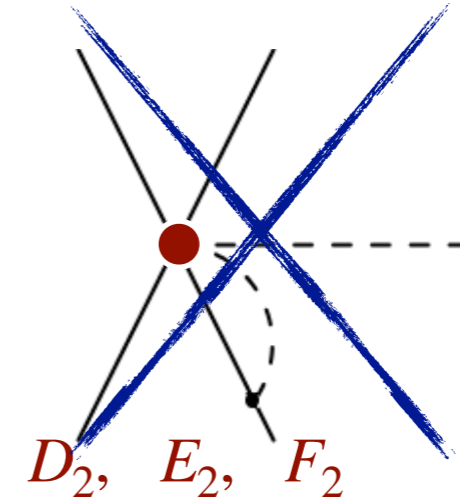
Beyond Weinberg counting

Induced potential

Vanishes

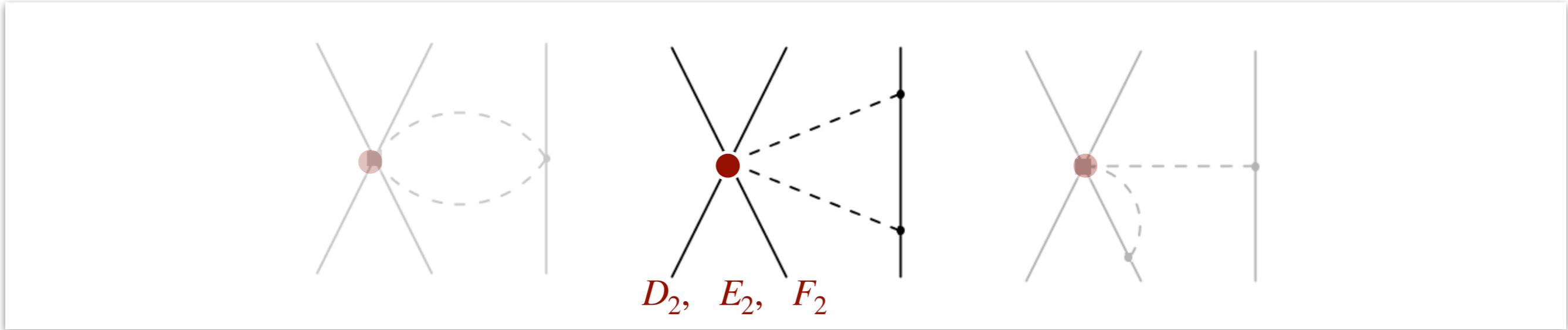


Induces a shift in c_D



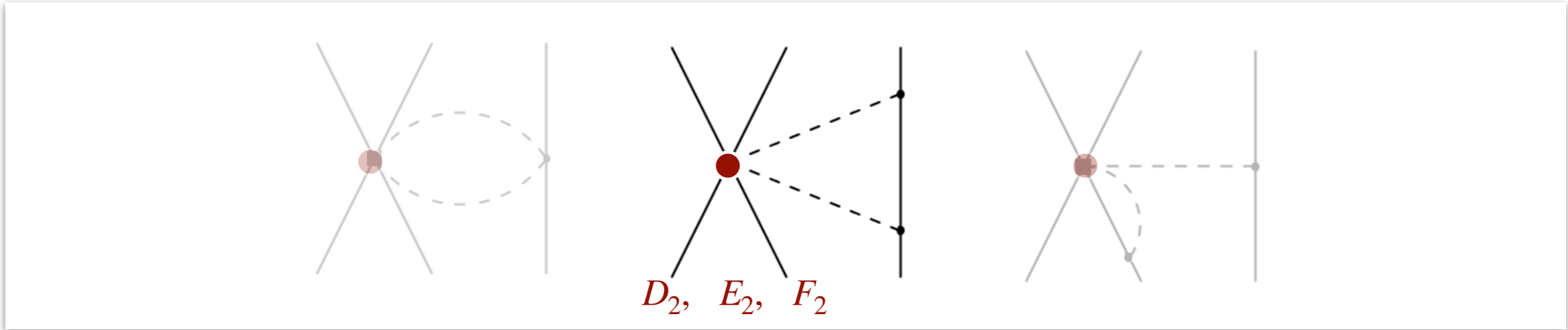
Beyond Weinberg counting

Induced potential



Beyond Weinberg counting

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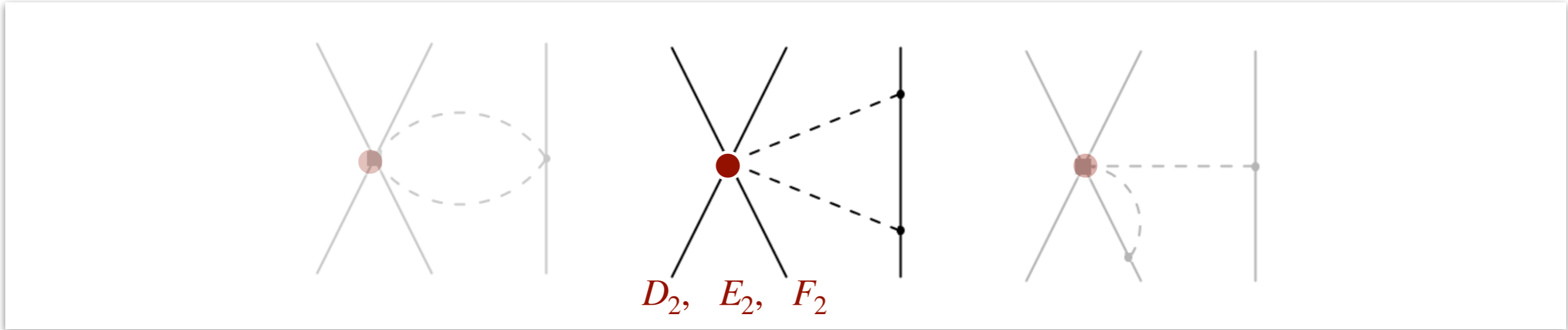


Three-nucleon potential from $m_\pi^2 D_2$

$$V(\vec{q}_1, \vec{q}_2, \vec{q}_3) = \frac{9g_A^2 D_2 m_\pi^3}{512\pi f_\pi^4} F\left(\frac{\vec{q}_3^2}{4m_\pi^2}\right) (1^{(i)}1^{(j)} - \vec{\sigma}^{(i)} \cdot \vec{\sigma}^{(j)}) \quad , \quad F(b) = \frac{2}{3} \left(1 + \left(\frac{1}{2\sqrt{b}} + \sqrt{b} \right) \tan^{-1}(\sqrt{b}) \right) .$$

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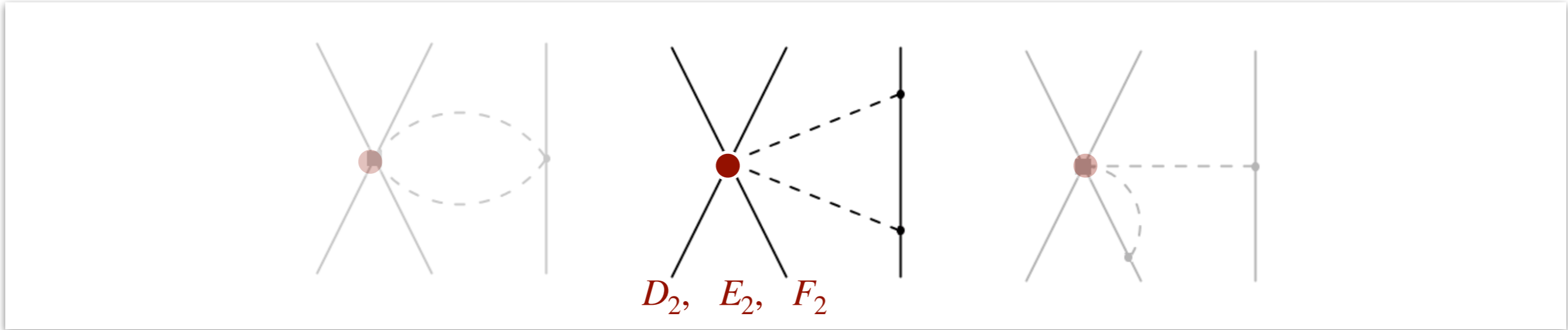
Contributions from $\sim E_\pi^2 E_2$

- E_2 induces same structure

- With $m_\pi^2 \rightarrow (\vec{q}^2/m_N)^2$

Beyond Weinberg counting

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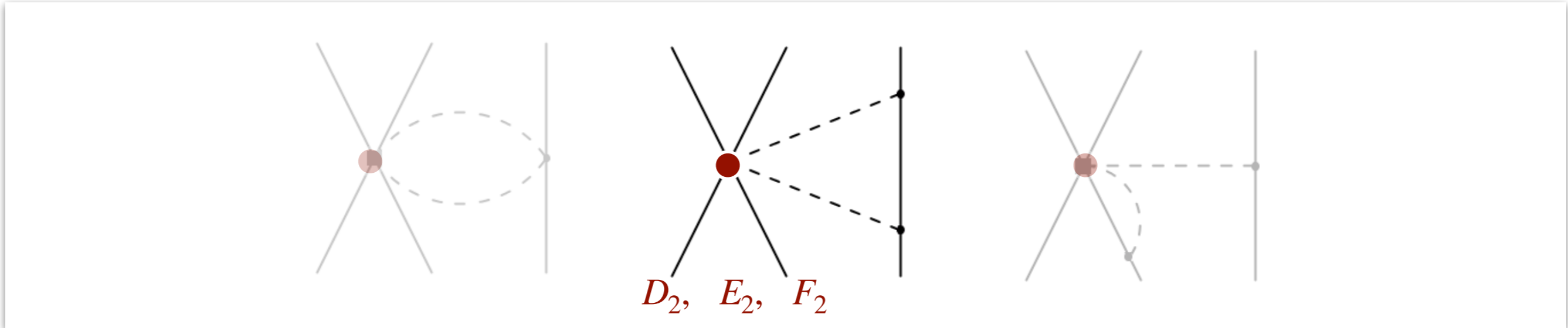
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Contributions from $\sim \vec{q}^2 F_2$

- F_2 induces same structure as D_2
- With additional factors of \vec{q}^2

Beyond Weinberg counting

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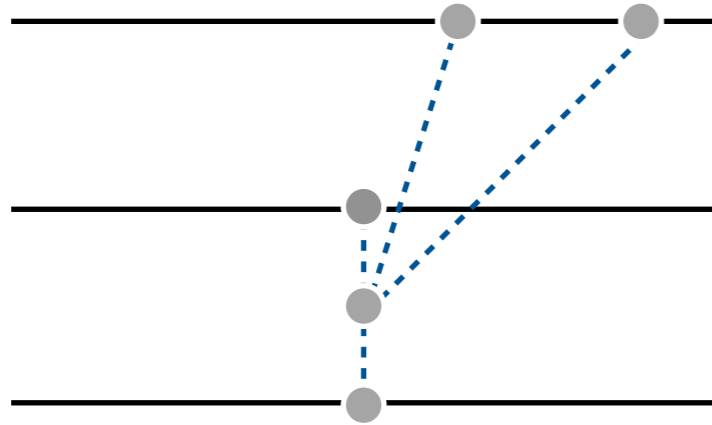
Contributions from $\sim E_\pi^2 E_2$

- E_2 Negligible for $n \lesssim 2n_{\text{sat}}$
- With $m_\pi^2 \rightarrow (\vec{q}^2/m_N)^2$

Contributions from $\sim \vec{q}^2 F_2$

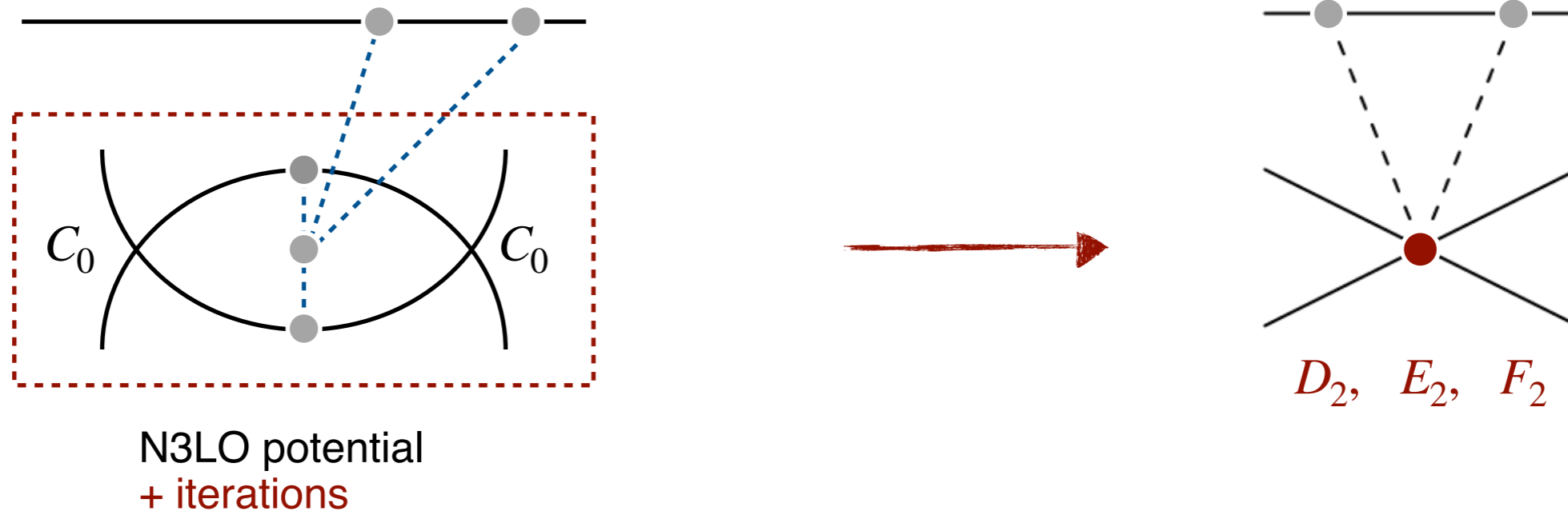
- F_2 induces same structure as D_2
- With additional factors of \vec{q}^2

Connection to usual N3LO graphs



N3LO potential

Connection to usual N3LO graphs



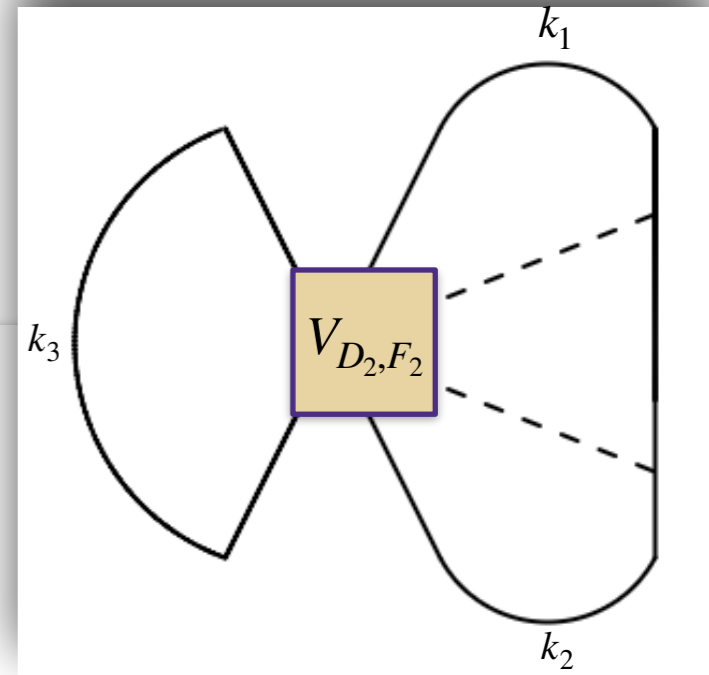
- Part of the 'conventional' N3LO potential is connected to D_2, E_2, F_2
- Generates the divergent diagrams that induce D_2, E_2, F_2
- Need to be considered simultaneously for a consistent calculation

Impact in dense matter

Effect in dense matter

Hartree-Fock estimate

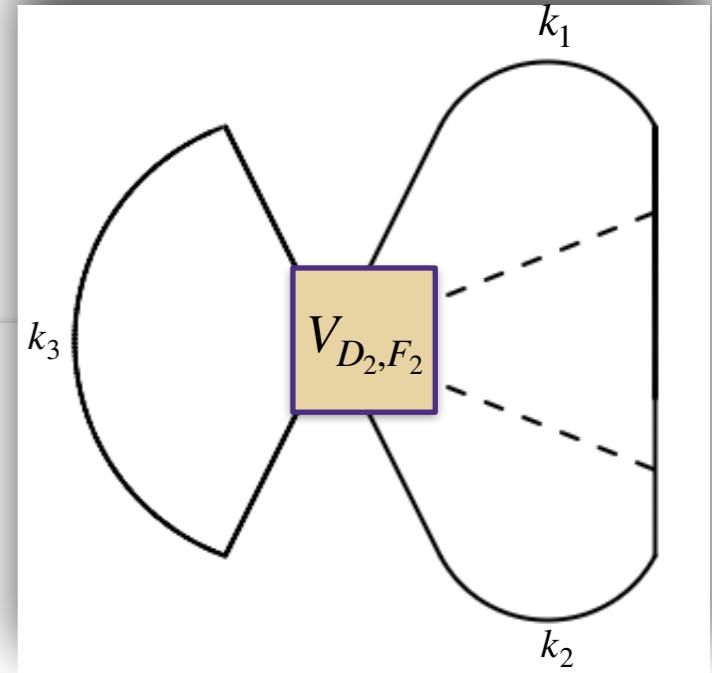
- Expectation value of V_{D_2, F_2} in Fermi gas state
 - Diagrammatically; contract nucleons in all possible ways
 - Finite, no need for additional regulators



$$\langle \mathcal{H}(0) \rangle = \int_{\vec{p}_1, \vec{p}_2, \vec{p}_3 \leq k_f} \left[V_{ijk}^{ijk}(0,0,0) - V_{ijk}^{ikj}(0, \vec{p}_{32}, \vec{p}_{23}) + V_{ijk}^{jki}(\vec{p}_{21}, \vec{p}_{32}, \vec{p}_{13}) + V_{ijk}^{kij}(\vec{p}_{31}, \vec{p}_{12}, \vec{p}_{21}) - V_{ijk}^{kji}(\vec{p}_{31}, 0, \vec{p}_{13}) - V_{ijk}^{jik}(\vec{p}_{21}, \vec{p}_{12}, 0) \right]$$

Effect in dense matter

Hartree-Fock estimate



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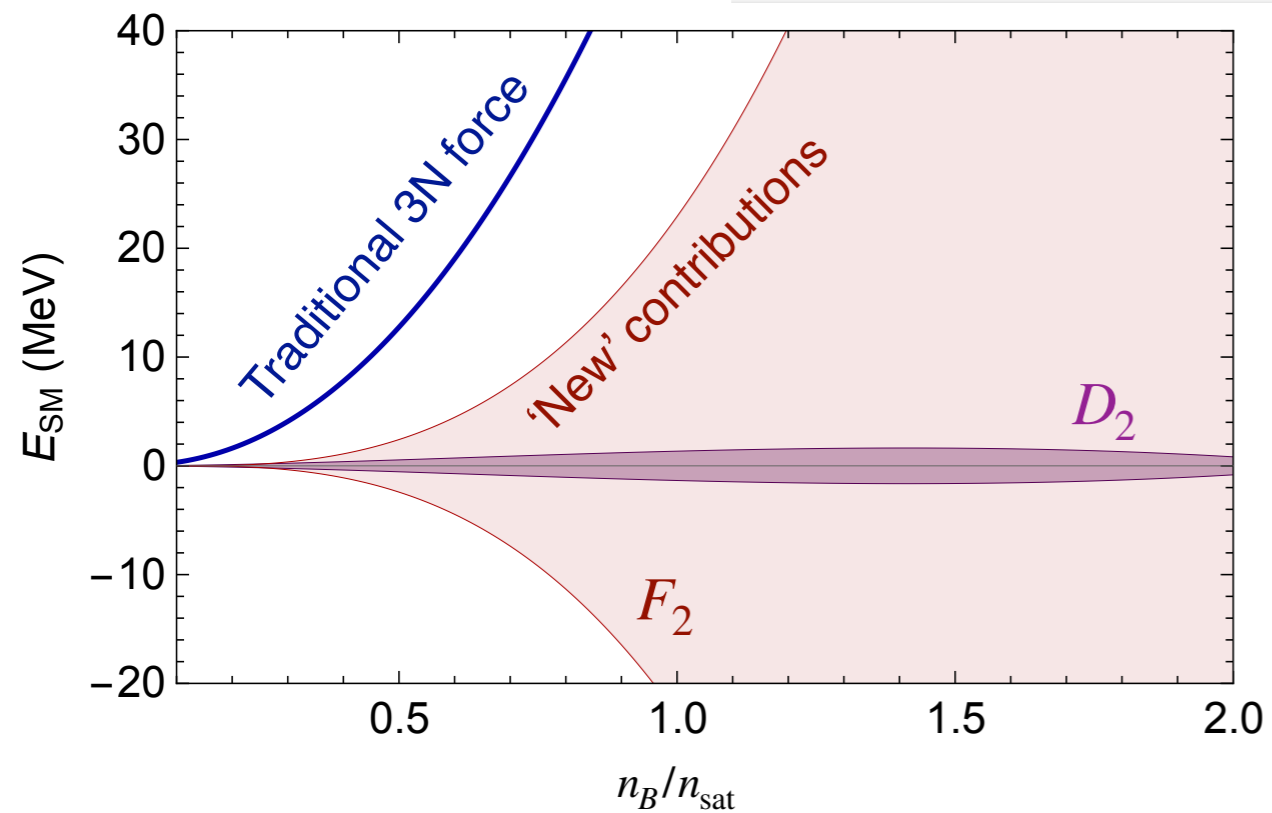
- Need the size of the LECs
 - Use expectation from RGE for a first estimate:

$$D_2 \leq 1/(5 F_\pi^4), \quad F_2 \leq 1/(5 F_\pi^4)$$

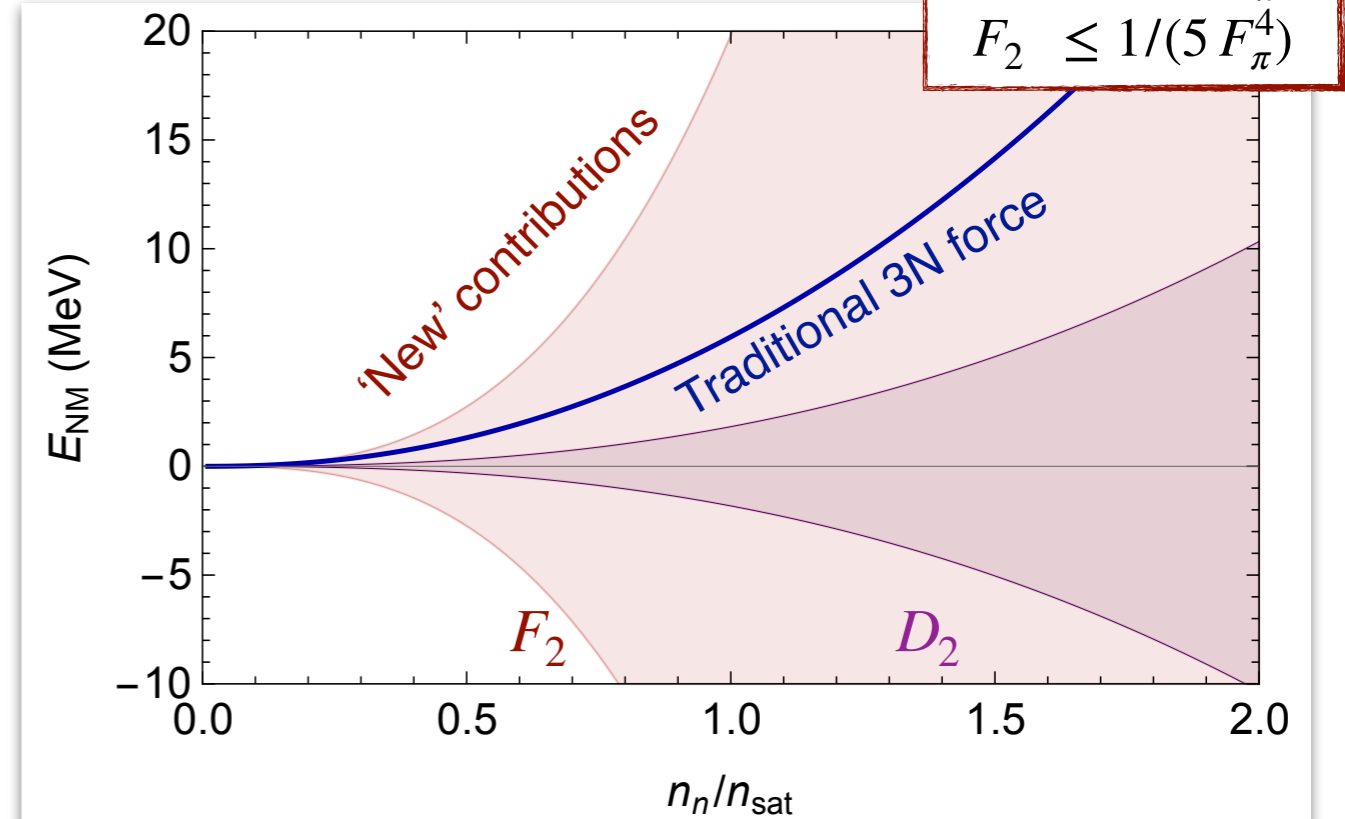
New class of three-nucleon forces

Effects in dense matter

Symmetric Matter



Neutron Matter

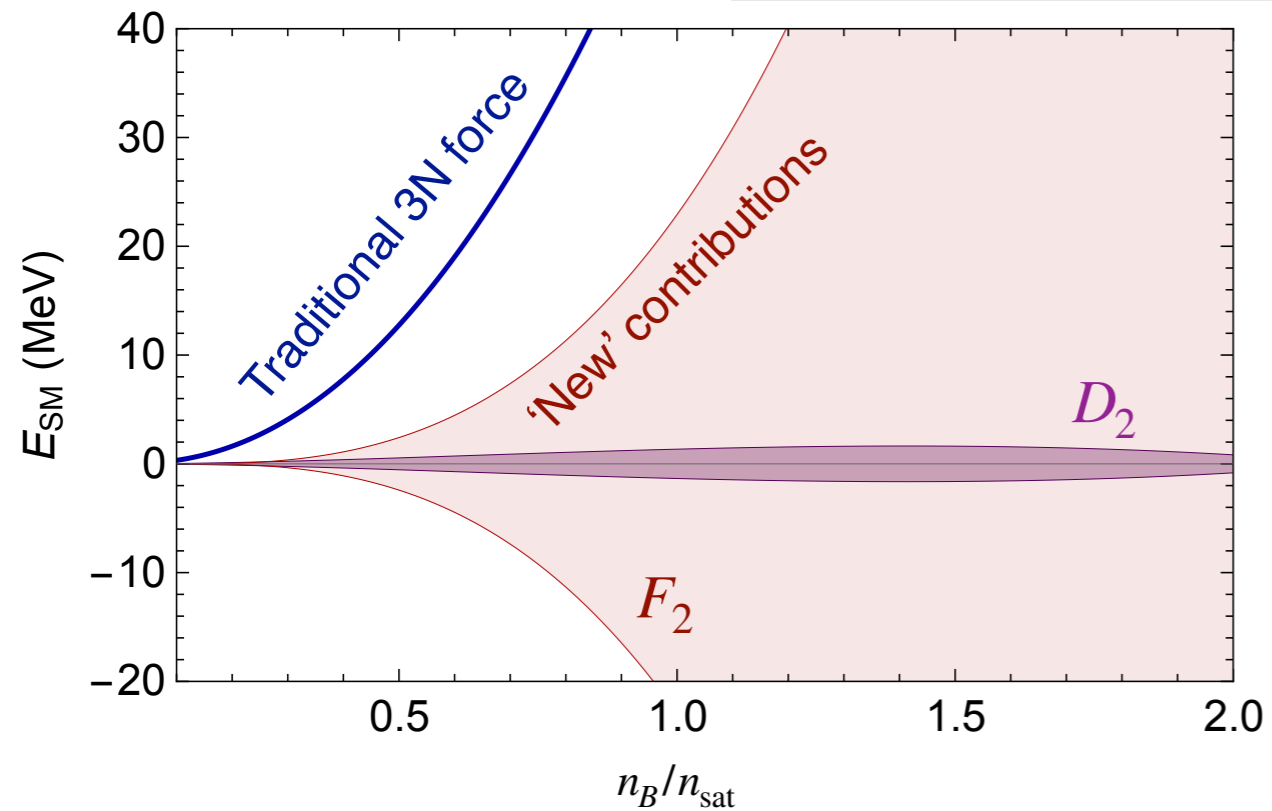


- Sizable contributions
- D_2, F_2 induce a stronger density dependence

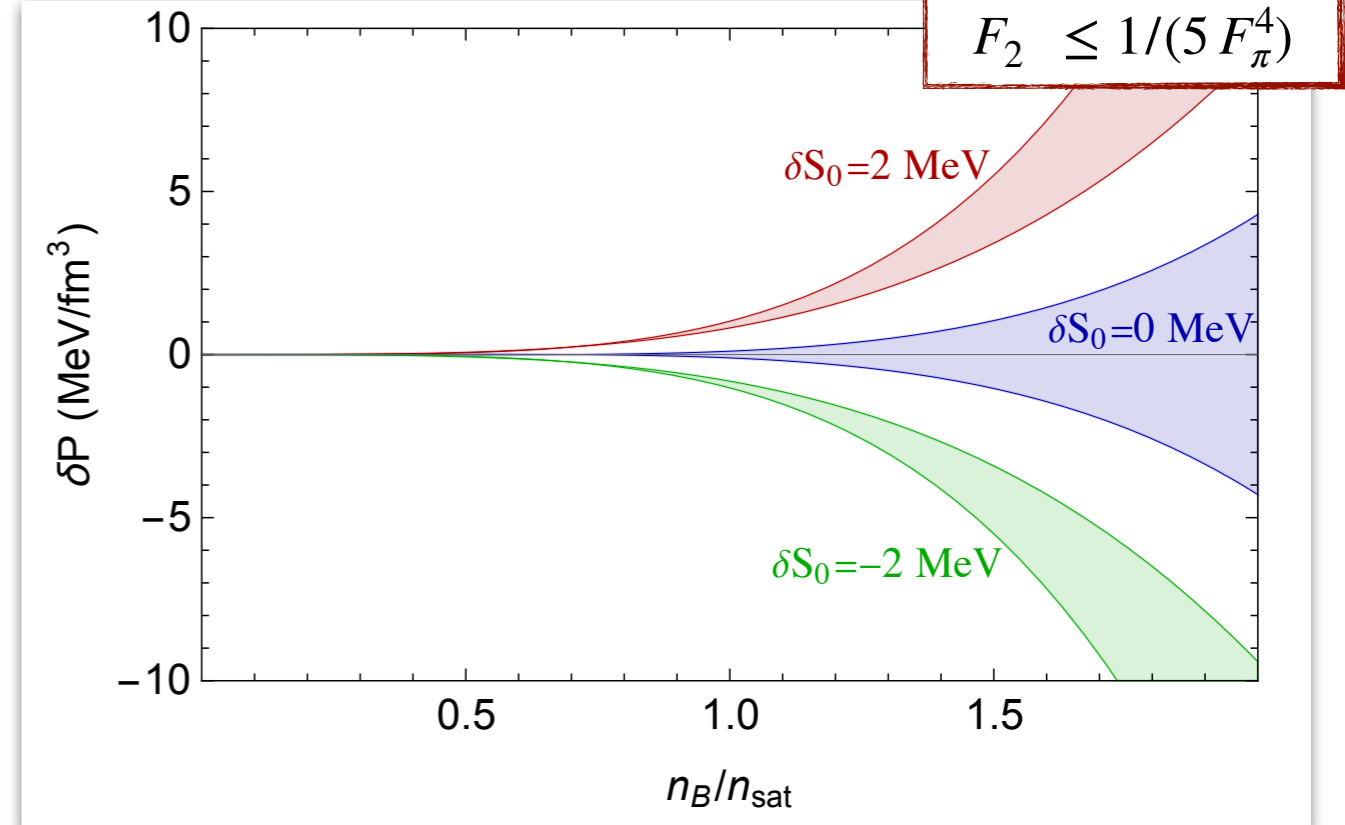
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Pressure in Symmetric Matter

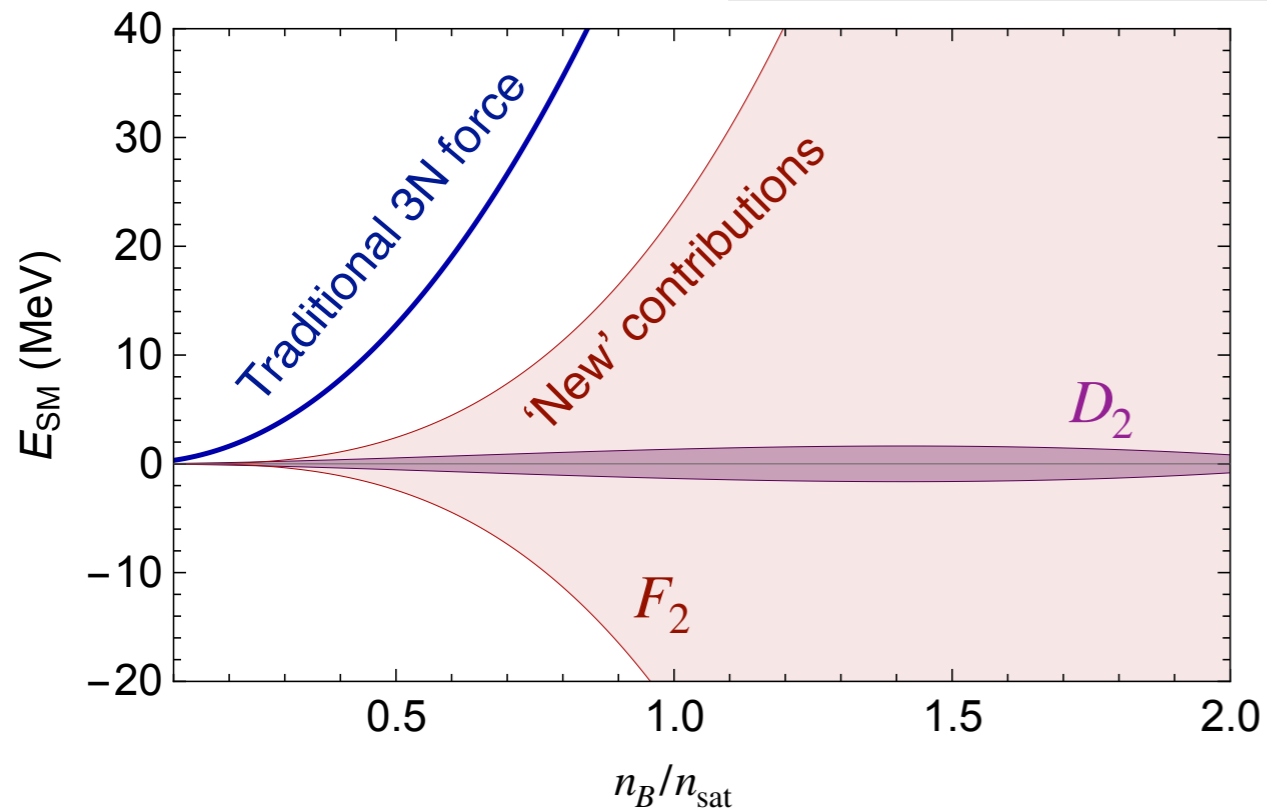


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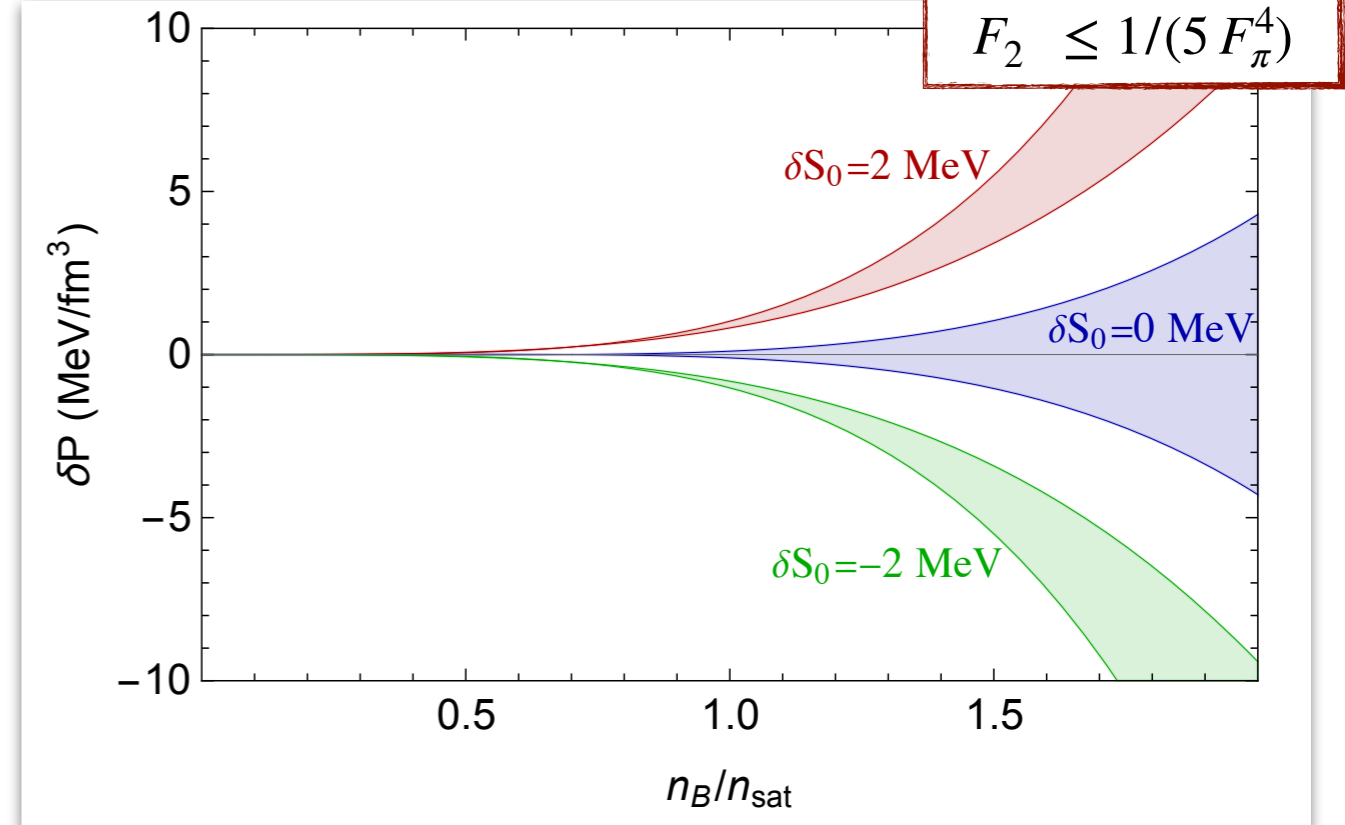
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• Scheme and regulator dependent

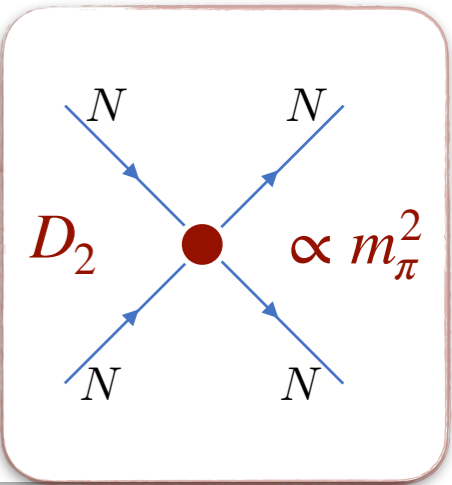
- Roughly $\sim 1/3$ smaller using dispersive regulator scheme [for $\Lambda = 500$ MeV at n_{sat}]
- Requires consistent combination with 'usual' N3LO three-nucleon force

- Crucially depend on the size of D_2, F_2

Impact Beyond the Standard Model

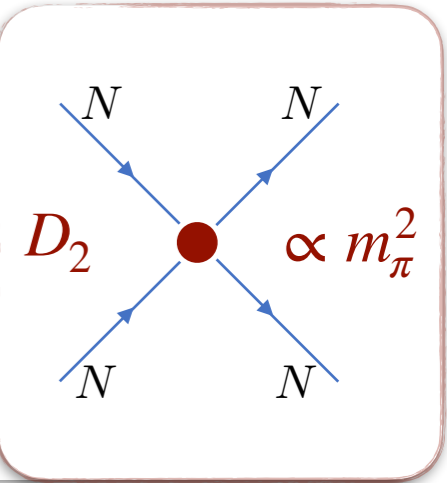
Effects on BSM scenarios

- D_2 induces m_π dependence of NN interactions



Effects on BSM scenarios

- D_2 induces m_π dependence of NN interactions

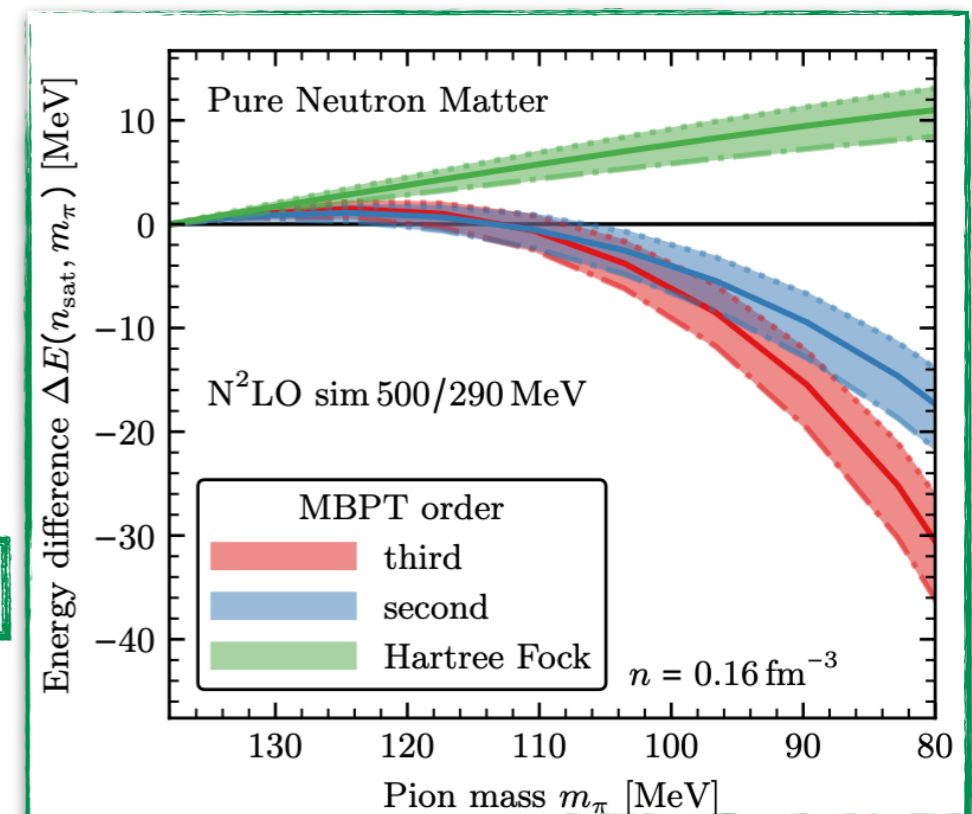


BSM scenarios can affect the quark masses

- Variations of fundamental constants
 - Lead to time dependent $m_q(t)$
- **Axion scenarios**
 - Axion could condense in dense matter like neutron stars
 - Would change $m_\pi(\theta = 0) \rightarrow m_\pi(\theta = \pi) \simeq 80 \text{ MeV}$

See Sanjay's talk

- Can be probed through their effect on the nuclear force
 - Requires m_π dependence of the nuclear force and D_2



Determining the LECs

How to determine D_2, F_2

From theory:

- First principles determination using Lattice QCD
 - Currently only calculations at unphysical m_π

e.g. Beane, Bedaque, Orginos, Savage, '06; Beane et al '15;

How to determine D_2, F_2

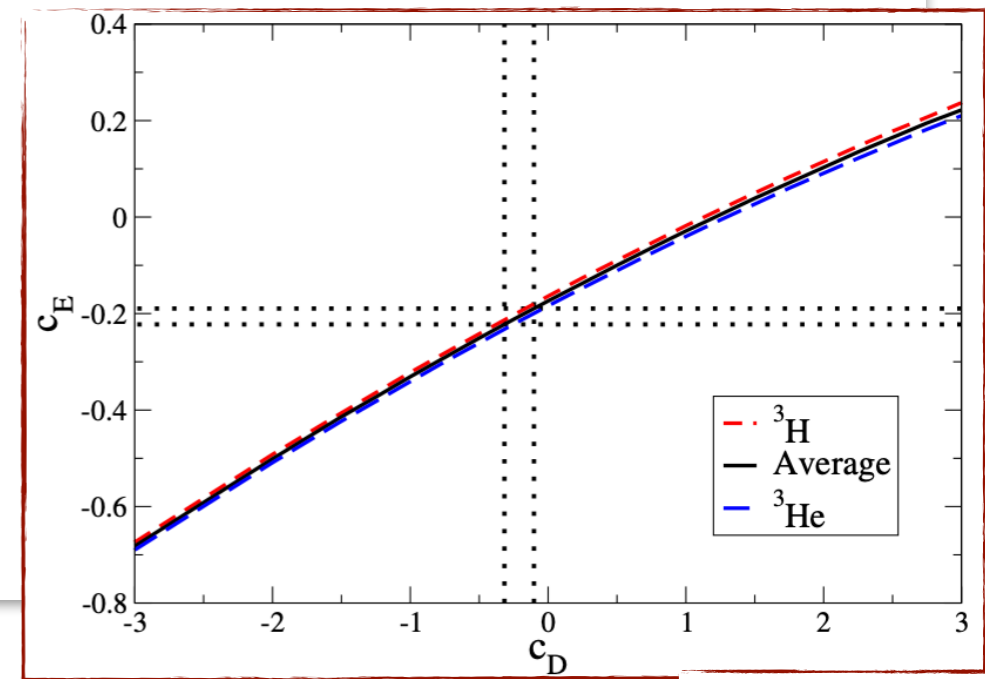
From theory:

- First principles determination using Lattice QCD
 - Currently only calculations at unphysical m_π

e.g. Beane, Bedaque, Orginos, Savage, '06; Beane et al '15;

From experiment:

- Determine D_2, F_2 together with $c_{D,E}$ from
 - Light systems:
 - Nd scattering
 - Binding energies
 - tritium β decay
 - Properties of dense matter
 - Properties of neutron stars
 - π -nucleus scattering

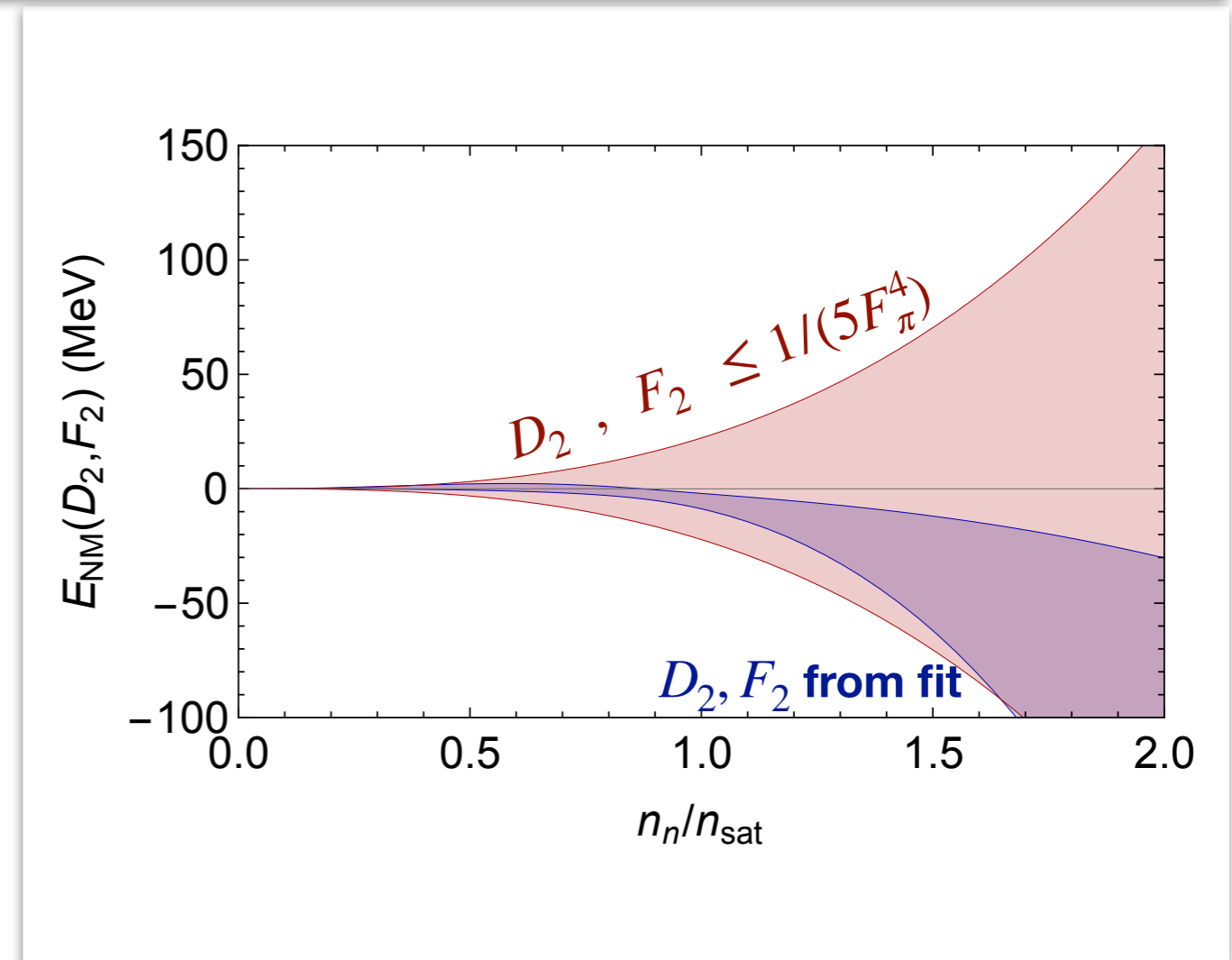
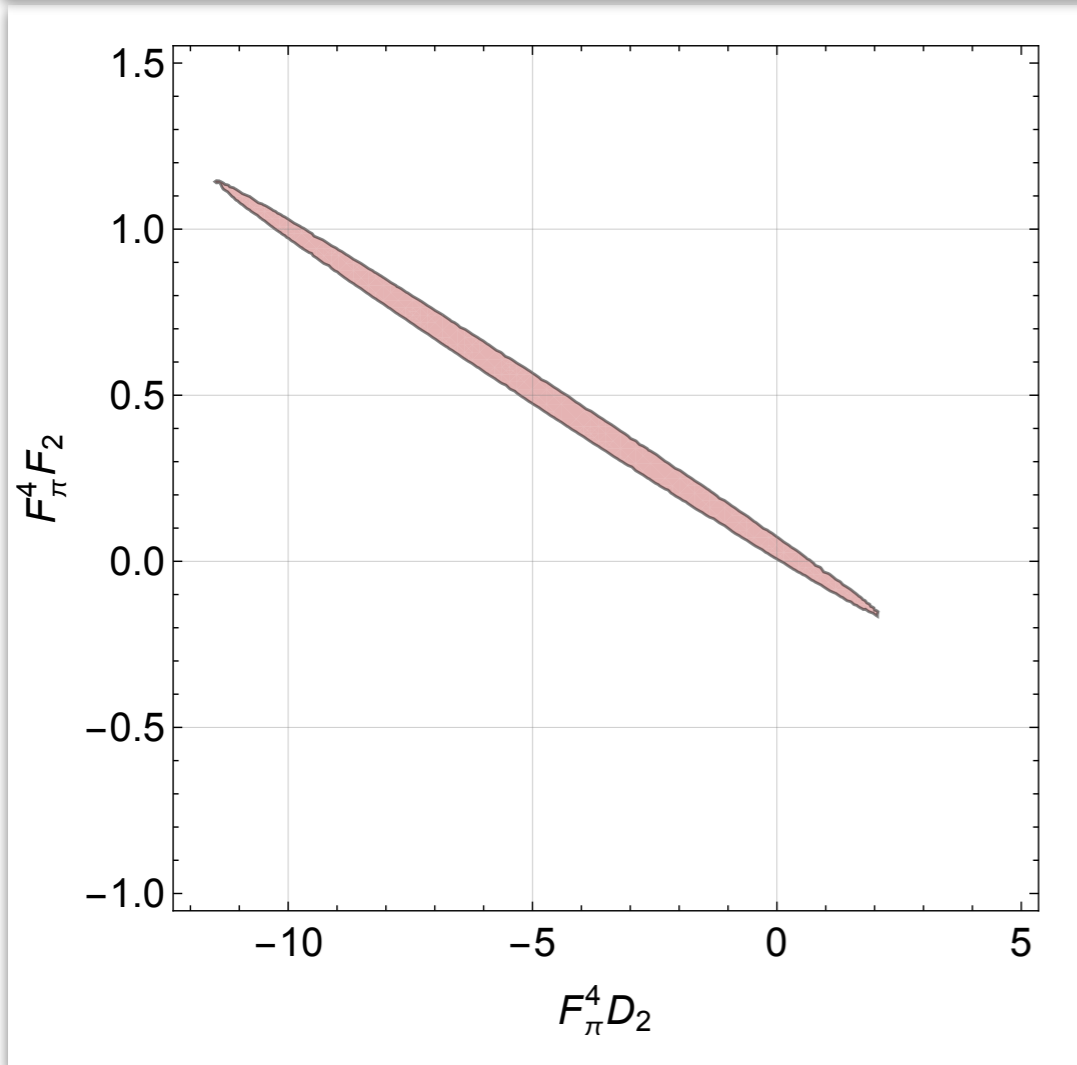


Gazit et al '09

How to determine D_2, F_2

Properties of dense matter

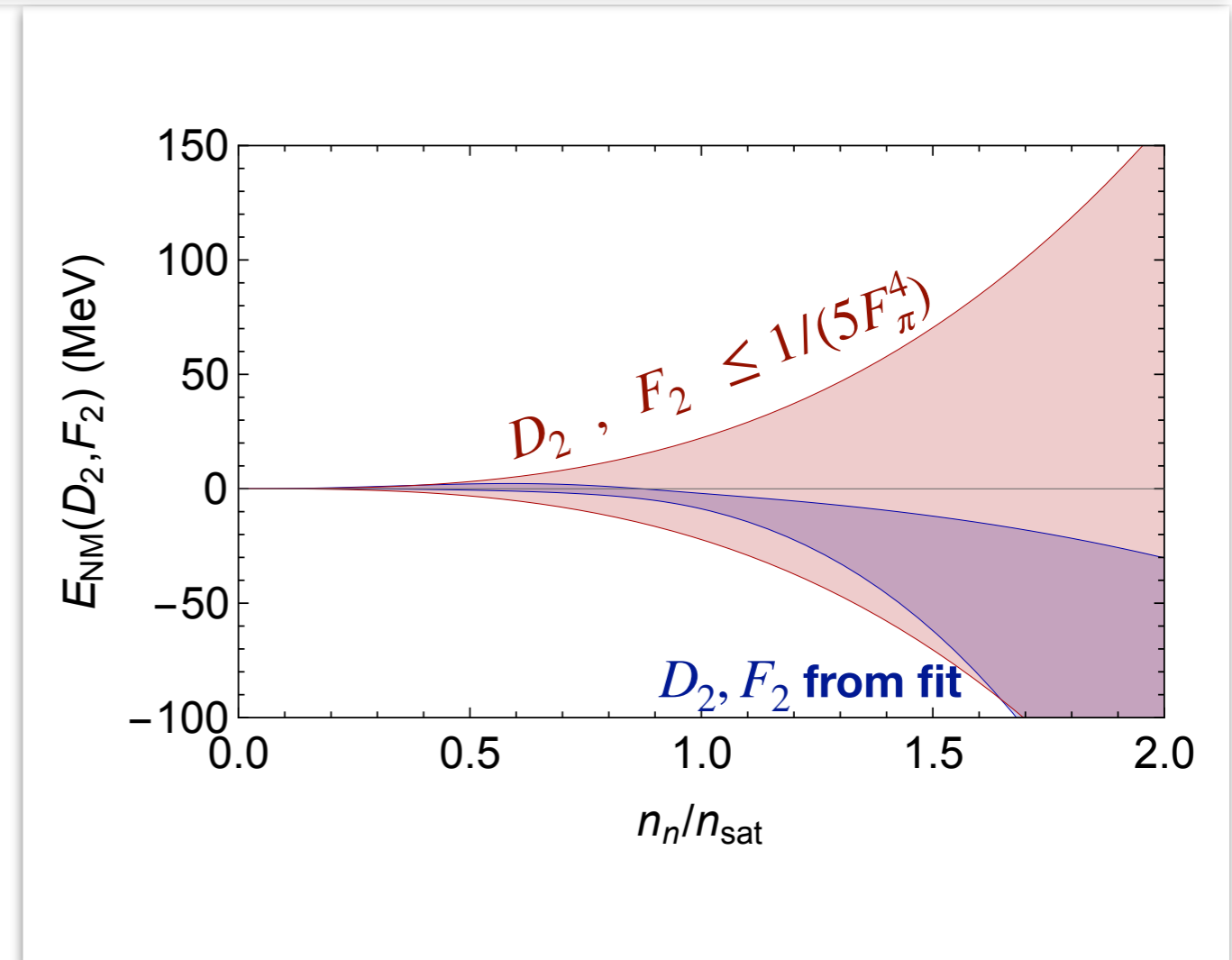
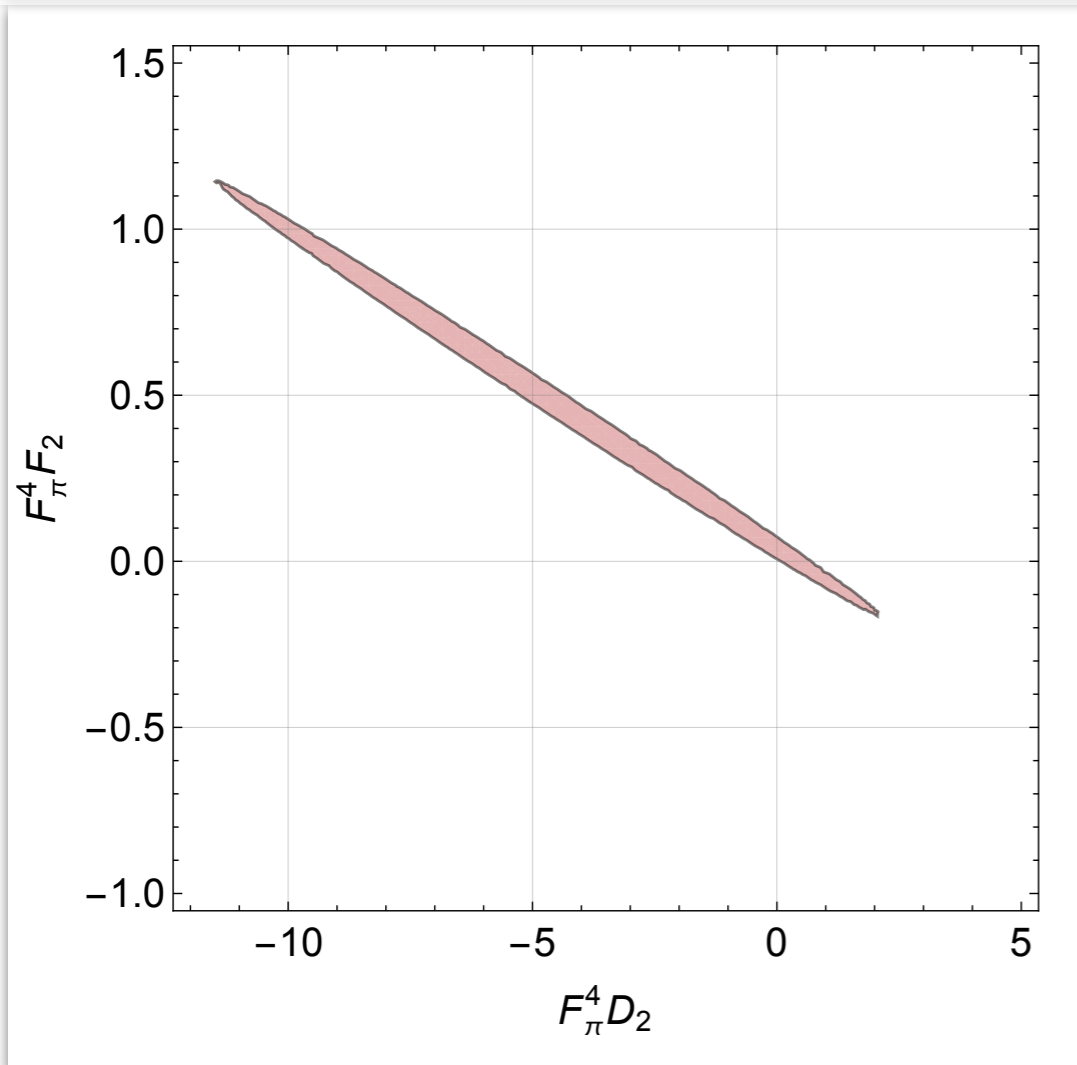
- **Naive** implementation
 - Combines HF estimates of 3-nucleon force with 2nucleon contributions
 - Fits to properties of dense matter near saturation



How to determine D_2, F_2

Properties of dense matter

Naive (inconsistent) implementation

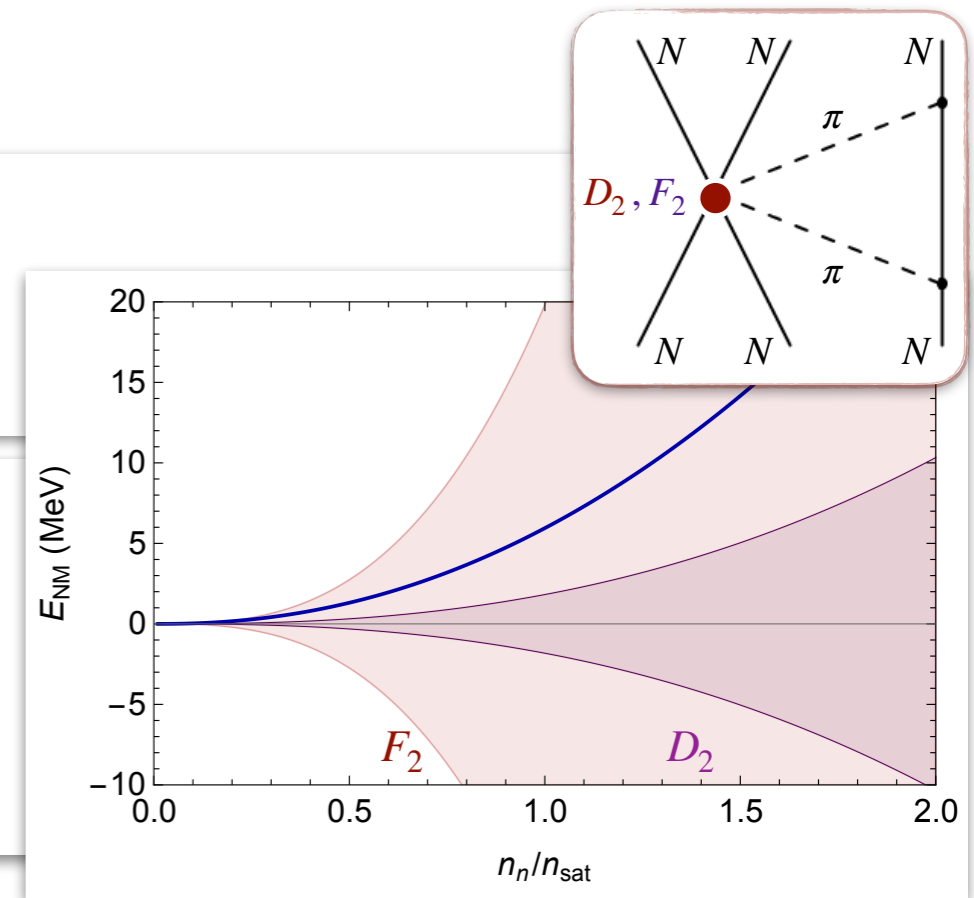


- **Inconsistent:** 2- and 3-nucleon forces using different regulators, Hartree-Fock estimates...
- Does suggest dense matter can help pin down D_2, F_2

- A **consistent** determination is work in progress (w/ C. Drischler, M. Kumamoto, M. Dawid, S. Reddy)

Summary

- Renormalization requires $\pi^2 NN$ interactions at LO
- Induce 3-nucleon forces through loops
- Significant contributions in dense matter
 - Important for neutron stars (equation of state)
 - Nuclei
- Crucially depends on the value of the LECs
 - Need to be fit to data



Questions/outlook

- Are these the only 'new' terms?
 - Additional contributions other channels?

de Vries, Gnech, Shain, '20

- Determining the LECs
 - From light nuclei or properties of dense matter
- Needed to assess the impact on
 - Dense systems
 - BSM scenarios & m_π dependence on nuclear force

Energy/particle at $m_\pi \neq m_\pi^{(phys)}$

